

New low temperature and superconductivity technology challenges

Angelo Nucciotti

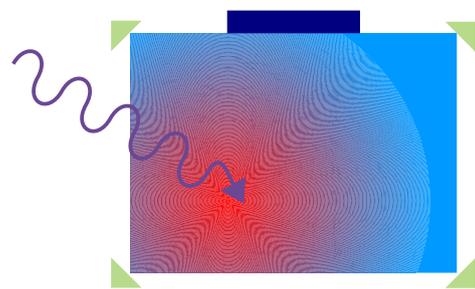
Dip. di Fisica "G. Occhialini", Università di Milano-Bicocca, Italia
INFN - Sezione di Milano-Bicocca, Italia



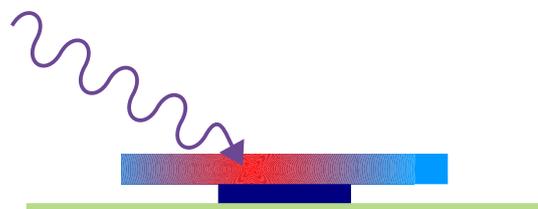
Sensor technologies for LTD

Dark Matter
 $\beta\beta 0\nu$
 γ spect. / n spect.
 Scint./Cherenkov
 light detectors
 Coherent ν scatt.
 Low E γ spect.
 neutrino mass
 X-ray spect.
 α spect.
 single photon
 mm \rightarrow THz
 bolometry/
 radiometry
 ...

quasi-equilibrium LTD
 thermal phonon sensors



calorimeter



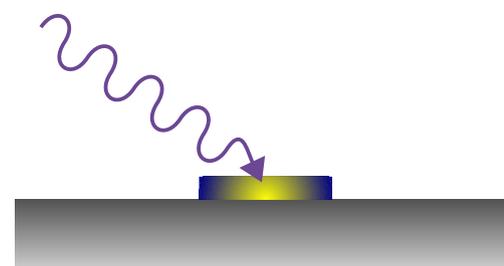
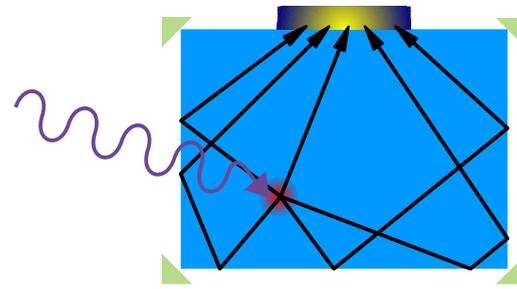
micro-calorimeter

$$E \rightarrow \Delta T \approx E/C \rightarrow \Delta X(T)$$

e.g: $R=R(T)$, $M=M(T)$

thermodynamic limit $\sigma_E^2 \approx k_B T^2 C$

out-of-equilibrium LTD
 athermal phonon sensors



kg



μg

Cooper pair breaking in S/C

$$E \rightarrow \Delta n_{qp} \approx E/\Delta \rightarrow \Delta X(n_{qp})$$

statistical limit $\sigma_E^2 \approx FE\Delta/\eta$

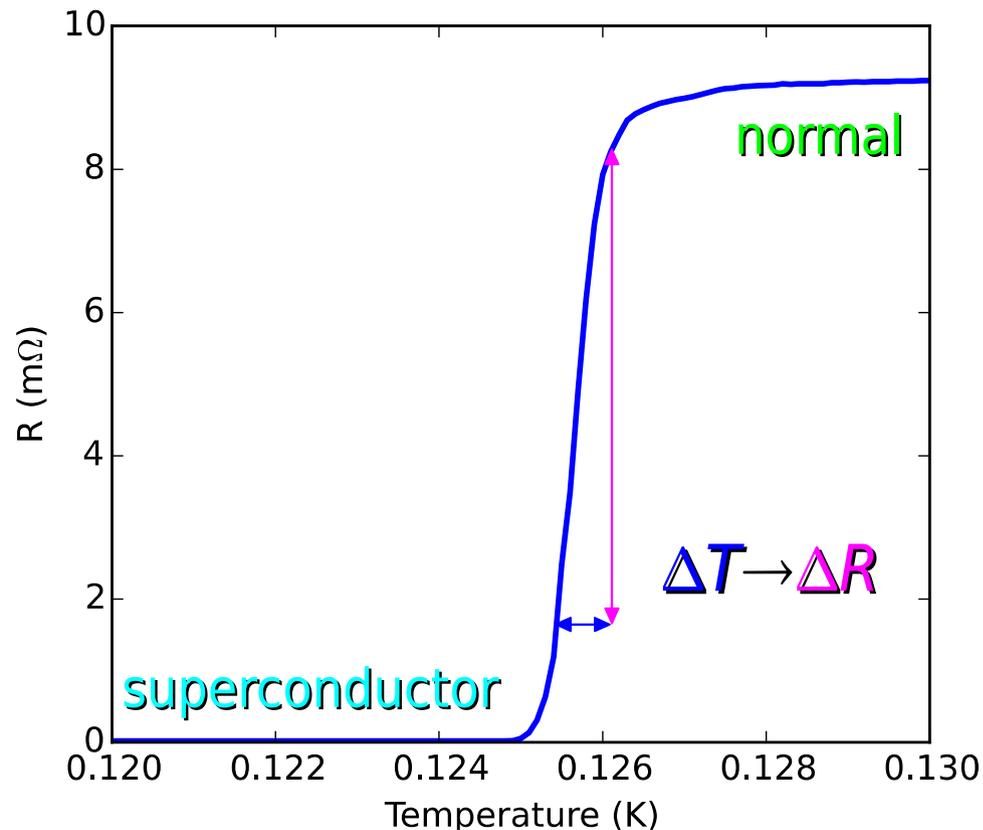
- high resolution ΔE / low E_{thres}
- fast τ_{rise} (\rightarrow BW)
- multiplexability \rightarrow large number of channels
- material choice flexibility

Low Temperature Detector (LTD) sensor technologies

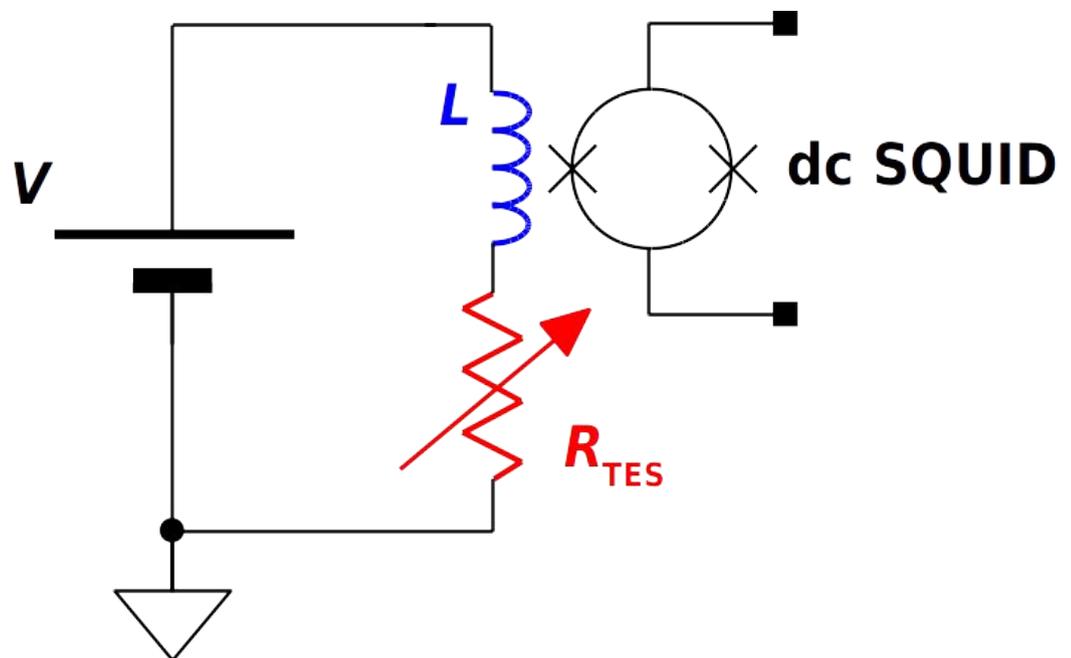
- Transition Edge Sensors (**TES**)
 - ▶ TES microcalorimeters
 - ▶ HOLMES TES
- Magnetic Metallic Calorimeters (**MMC**)
 - ▶ ECHo MMC
- Kinetic Inductance Detectors (**KID**)
 - ▶ KID for athermal phonon detection
 - ▶ CALDER KID
 - ▶ thermal KID
- KID microwave **multiplexing**
 - ▶ detector microwave multiplexing by upconversion

Superconducting transition edge sensors (TES)

- superconductor thin films used inside the phase transition at T_c
 - ▶ pure superconductors: Ir ($T_c = 112$ mK), W ($T_c = 15$ mK), ...
 - ▶ metal-superconductor bilayers \Rightarrow tunable T_c (20 \div 200 mK) : Mo/Cu, Ti/Au, Ir/Au, ...
- high sensitivity ($A \approx 100$) \Rightarrow high energy resolution
 - ▶ as **thermal sensors** $\rightarrow \sigma_E^2 \approx \xi^2 k_B T^2 C$
 - ▶ also as **athermal sensors**
- high electron-phonon coupling \Rightarrow high intrinsic speed
- low impedance \Rightarrow SQUID read-out \Rightarrow multiplexing for large arrays

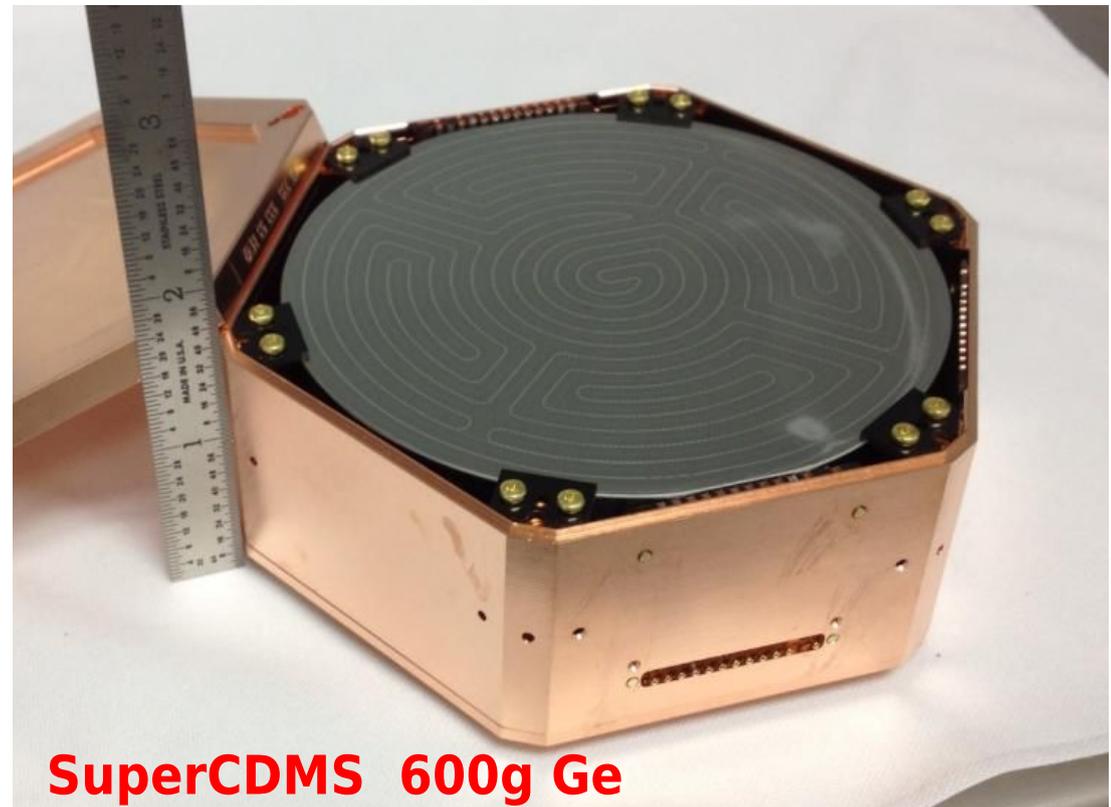
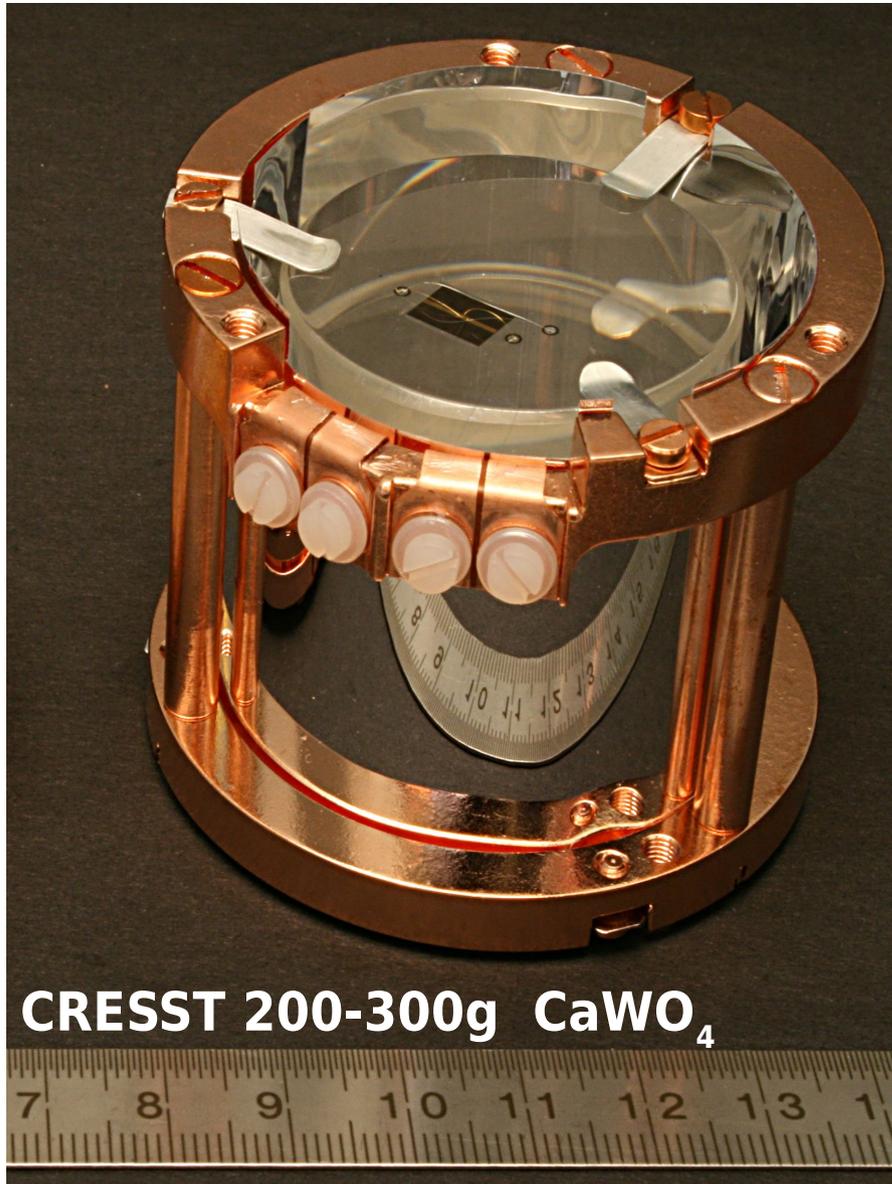


TES read-out: constant voltage bias



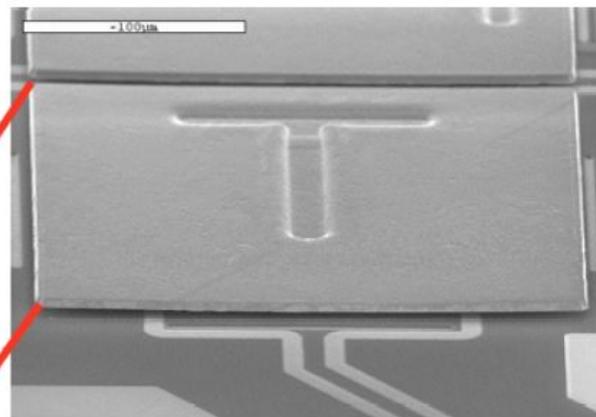
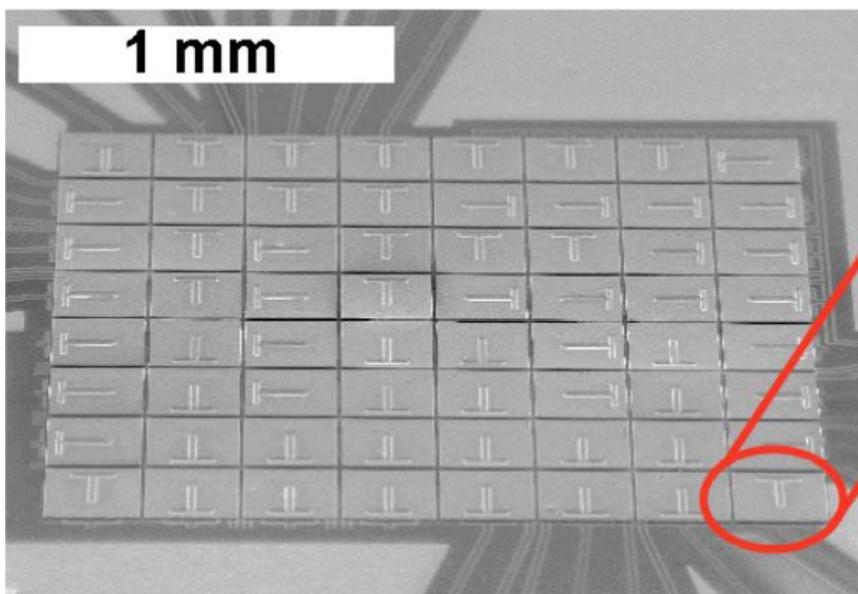
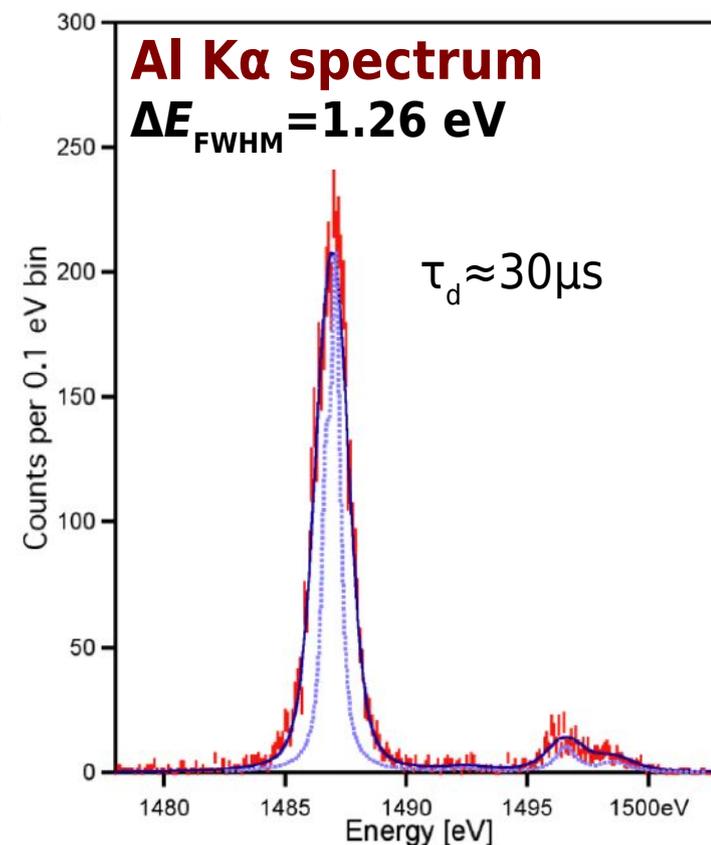
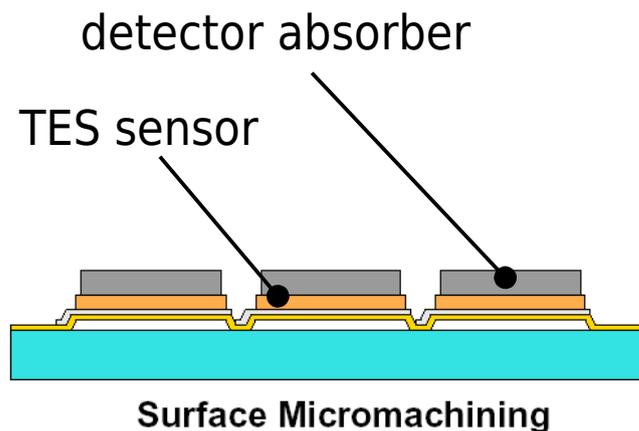
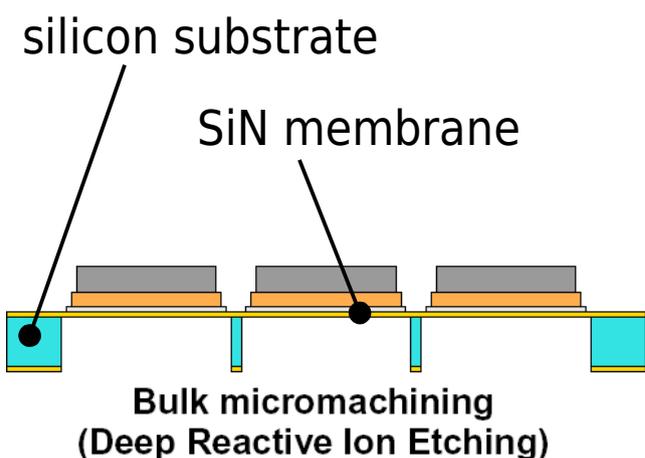
TES for massive LTDs: direct Dark Matter searches

- thin film deposition
- photolithographic patterning



TES microcalorimeters / 1

- detectors for α , β , low E γ , and X-ray science
 - astrophysics, material science, nuclear physics, ...



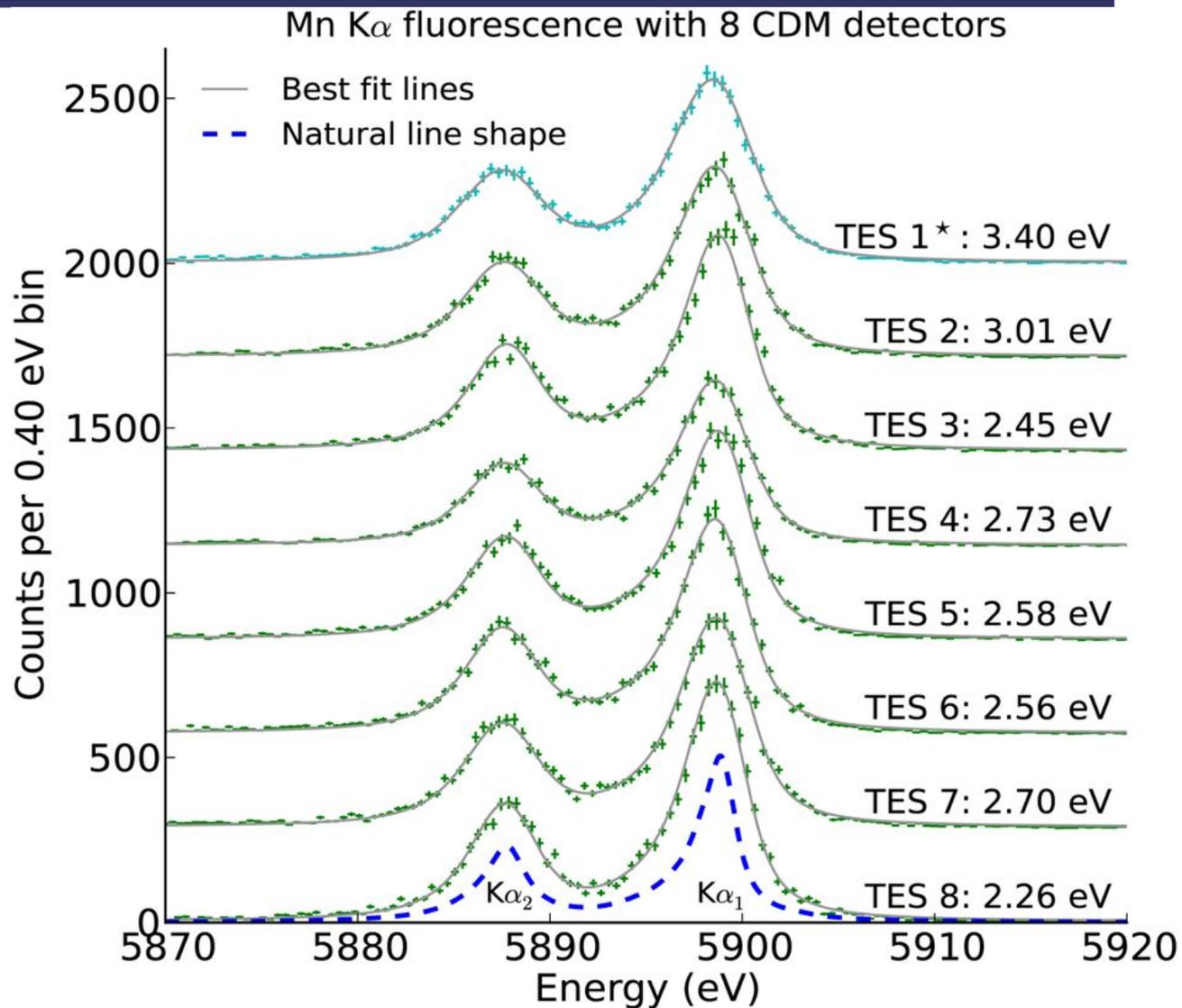
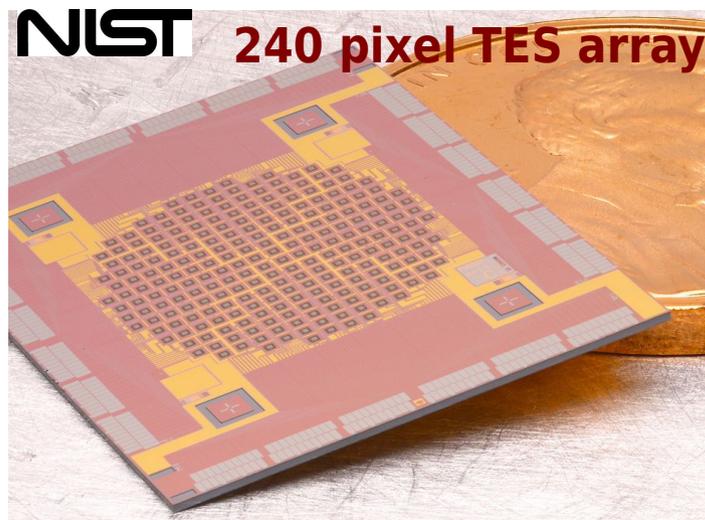
← 250 μm →



TES microcalorimeters / 2

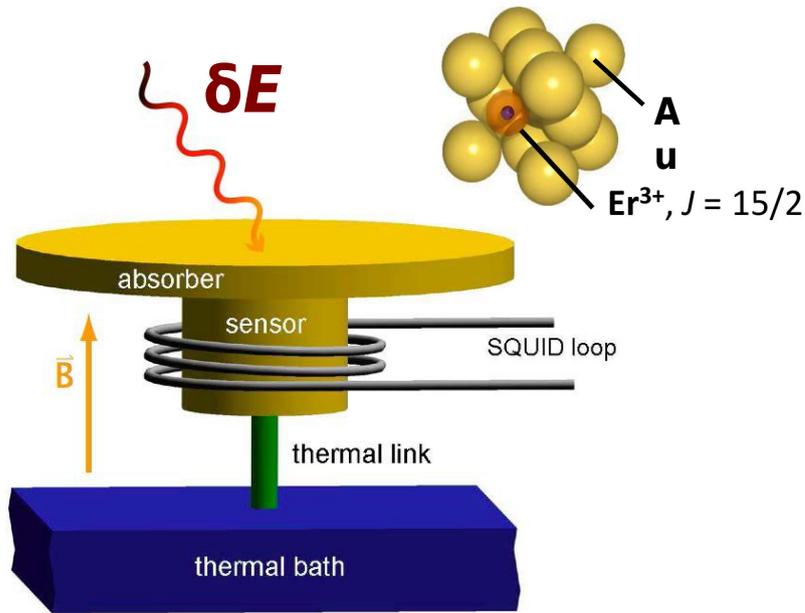
X-ray spectroscopy

XES, EXAFS, XANES, RSXS ...
time resolved X-ray spectroscopy

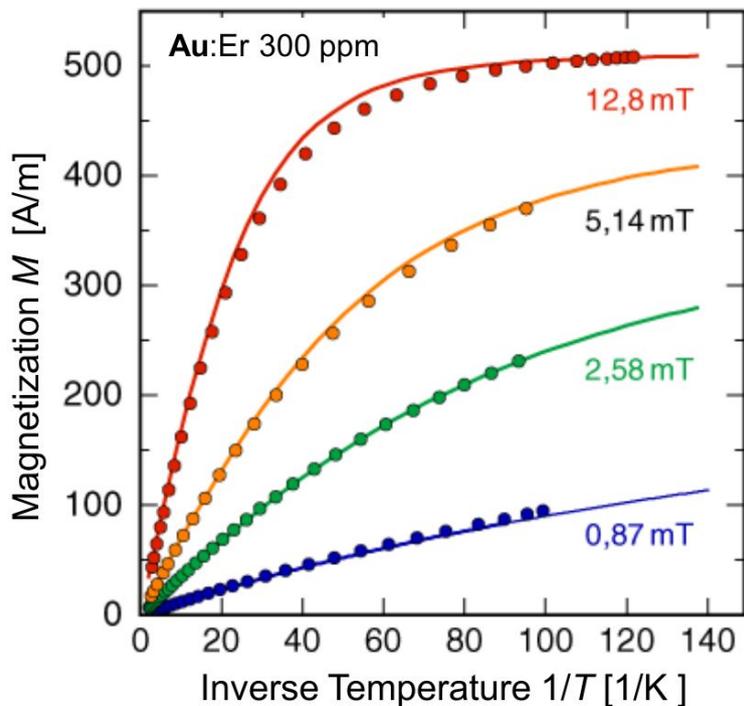


- 240 pixels, 23.4 mm² active area, 30% fill factor
- Code Division SQUID multiplexing
- $\Delta E \approx 2.5$ eV at 6 keV
- 80% Quantum Efficiency at 6 keV

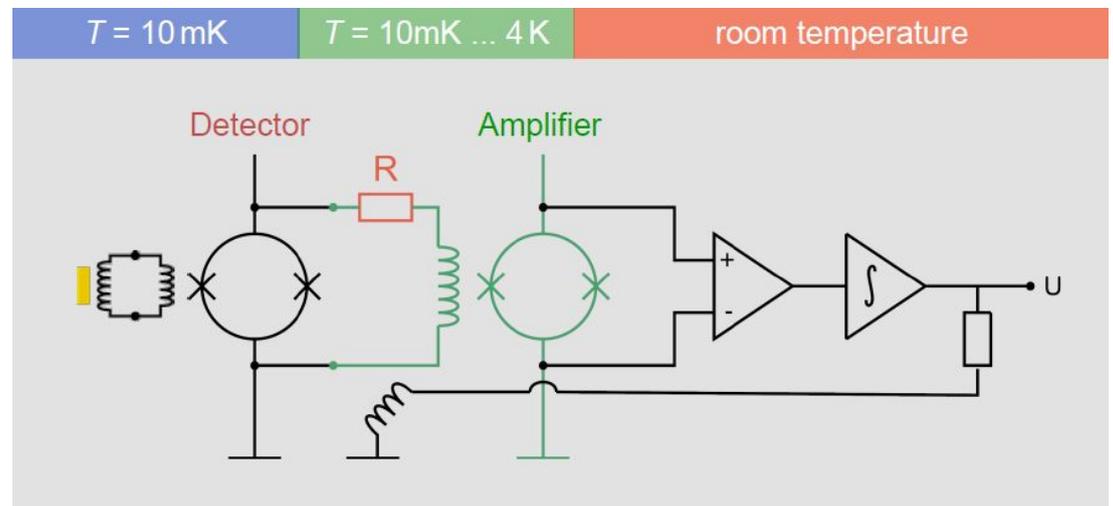
Metallic Magnetic Calorimeters (MMC)



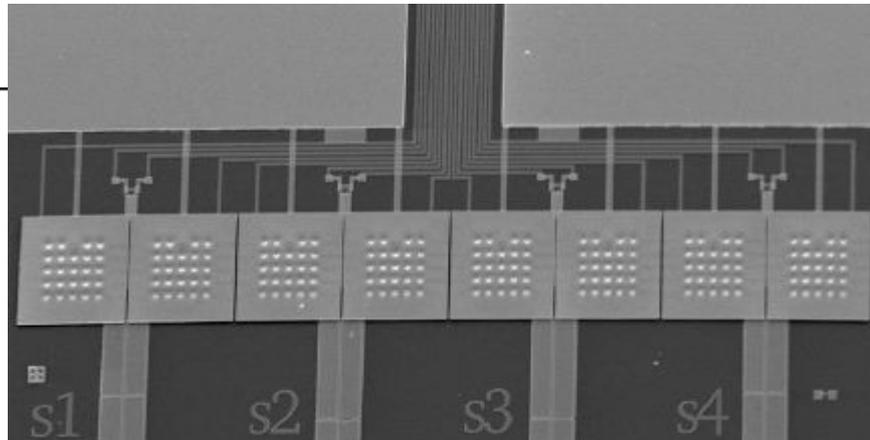
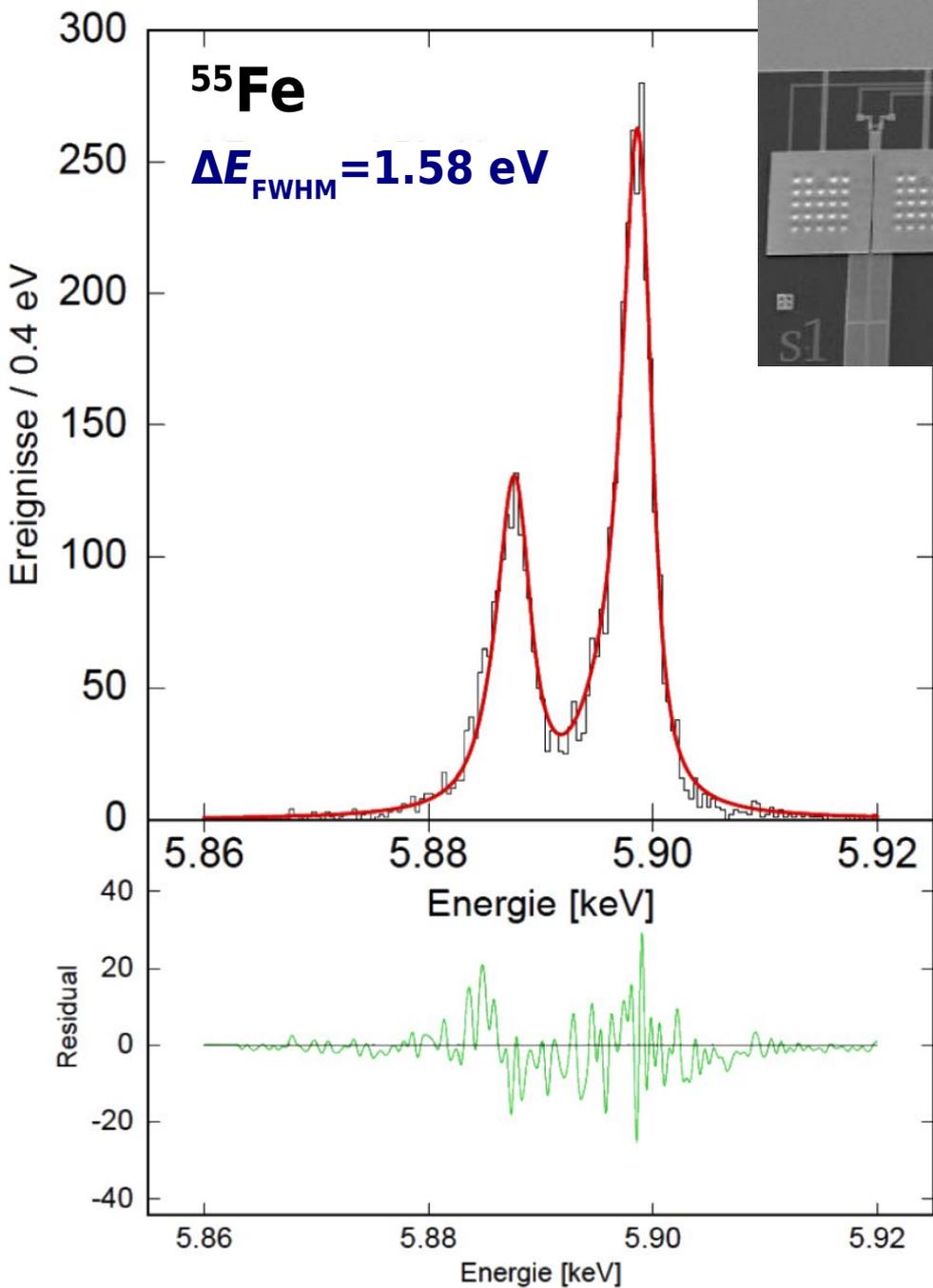
$$\delta E \rightarrow \delta M \rightarrow \delta \Phi$$



- paramagnetic temperature sensor
 - ▶ **Au:Er**, Ag:Er, PbTe:Er, Dy:W, W:Fe ...
- **dc-SQUID read-out**
 - ▶ high energy resolution with metallic absorbers
 - ▶ **fast rise time ($\approx 100\text{ns}$)**
- **high C**
 - ▶ **massive absorbers**
 - ▶ high linearity
- no power dissipation in the sensor
- multiplexing by frequency up-conversion



MMC development at Heidelberg



Electron capture end-point experiments

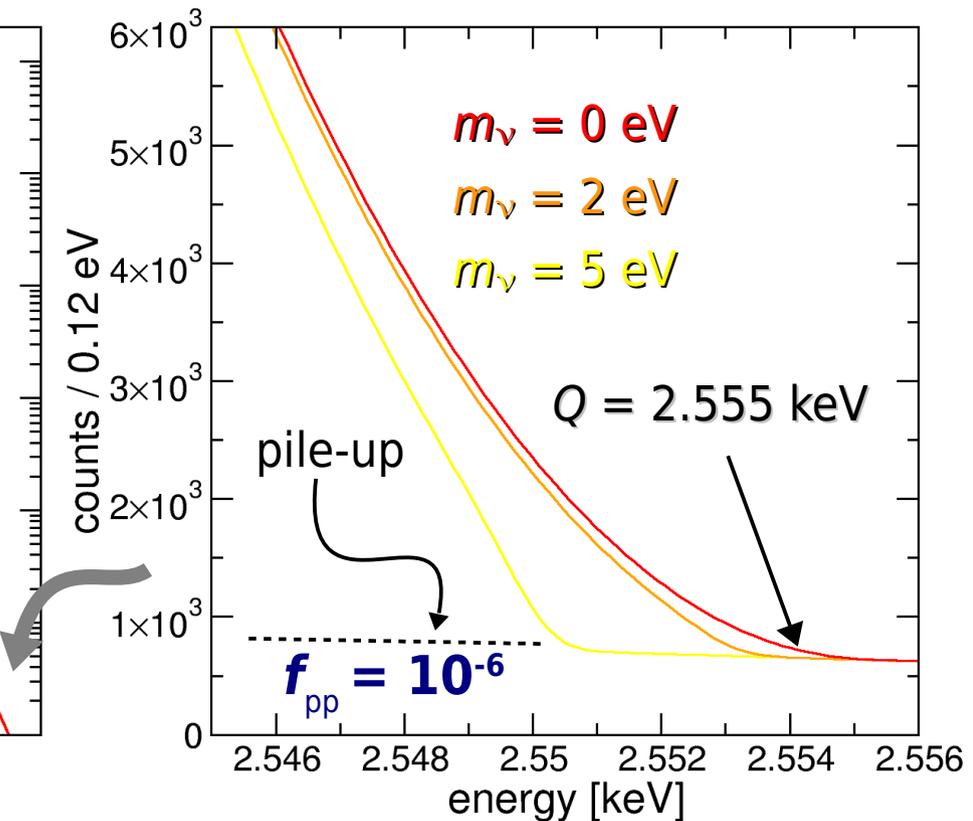
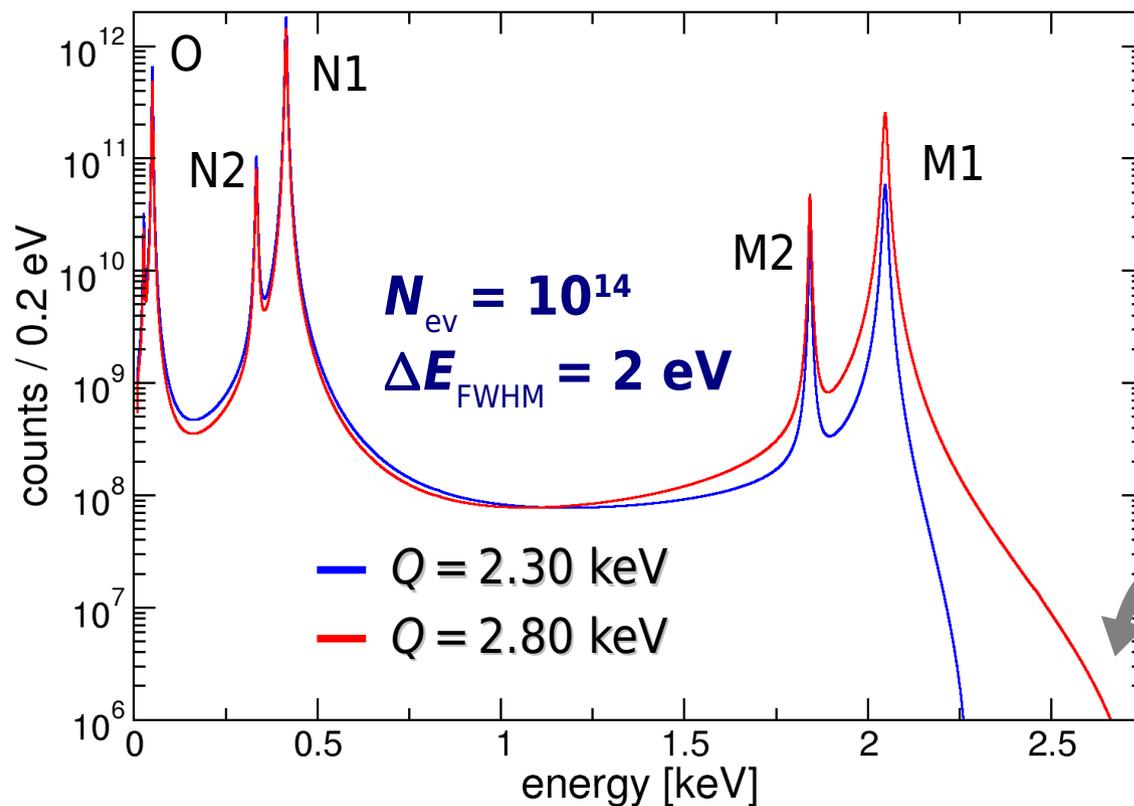


electron capture from shell \geq M1

A. De Rujula and M. Lusignoli, Phys. Lett. B 118 (1982) 429

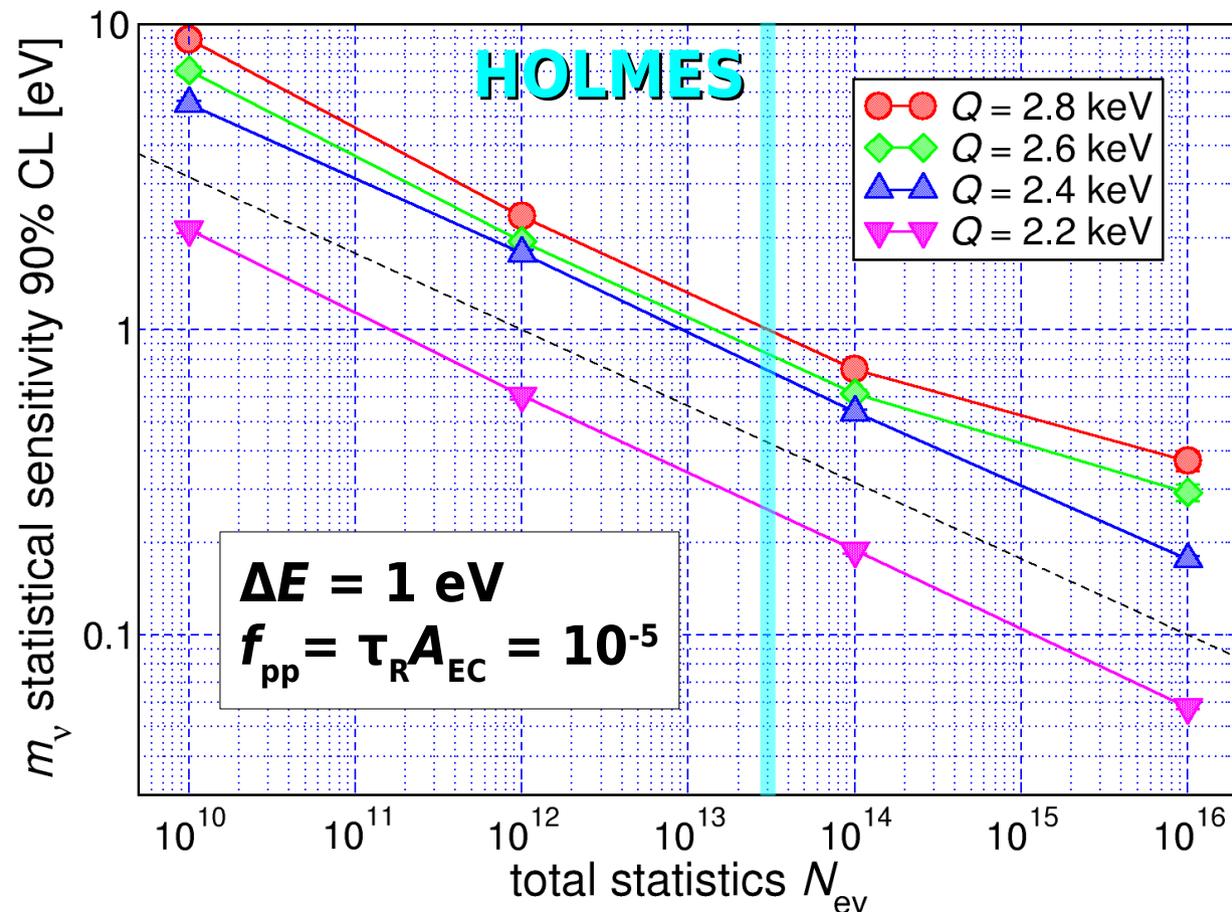
- calorimetric measurement of Dy atomic de-excitations (mostly non-radiative)
- rate at end-point and ν mass sensitivity depend on Q
 - ▶ Past measurements: $Q = 2.3 \div 2.8$ keV. Recently measured $Q = 2.83 \pm 0.04$ keV
- $\tau_{1/2} \approx 4570$ years \rightarrow few active nuclei are needed

$$\frac{d\lambda_{EC}}{dE_c} = \frac{G_\beta^2}{4\pi^2} (Q - E_c) \sqrt{(Q - E_c)^2 - m_\nu^2} \times \sum_i n_i C_i \beta_i^2 B_i \frac{\Gamma_i}{2\pi} \frac{1}{(E_c - E_i)^2 + \Gamma_i^2/4}$$



m_ν statistical sensitivity: Montecarlo simulations

- 2×10^{11} ^{163}Ho nuclei \rightarrow 1 decay/s
- ^{163}Ho production: p.e. neutron irradiation of ^{162}Er enriched Er
- embed ^{163}Ho in thermal detectors for low energy X-rays spectroscopy

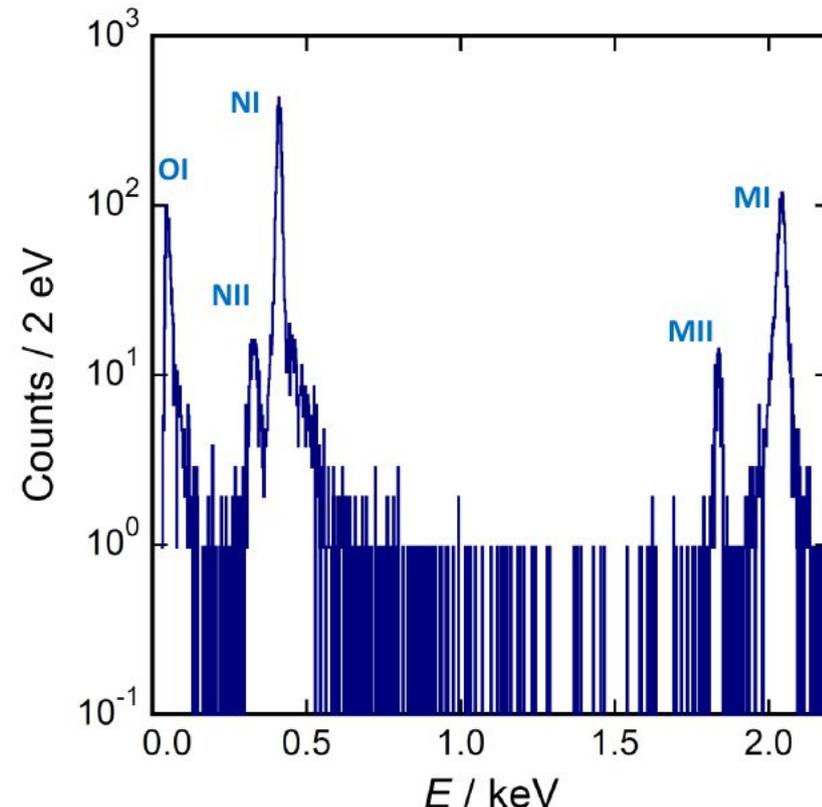
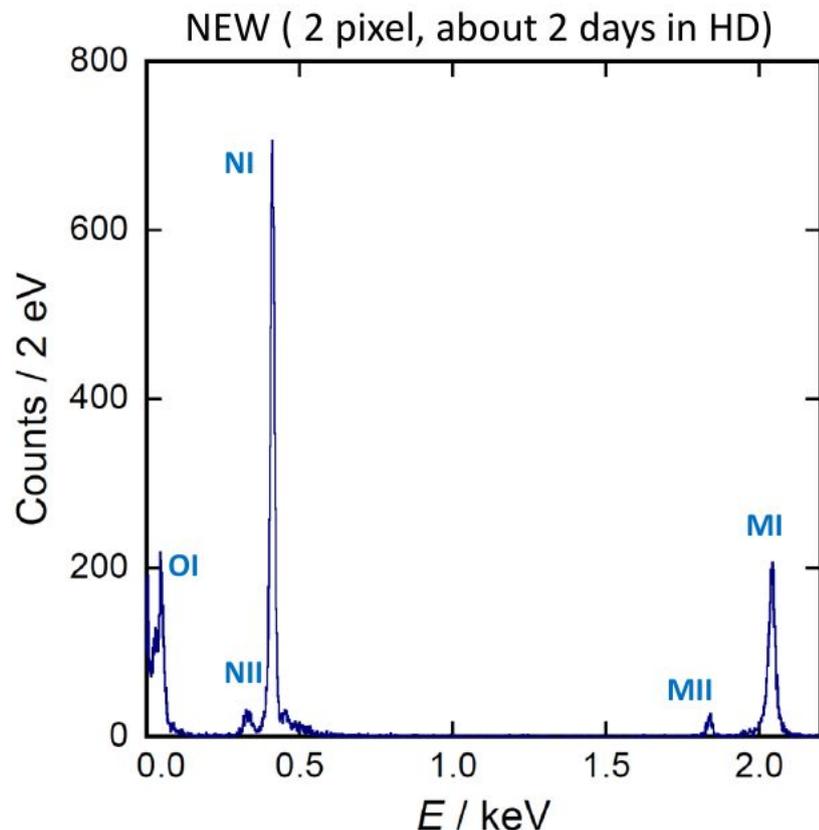
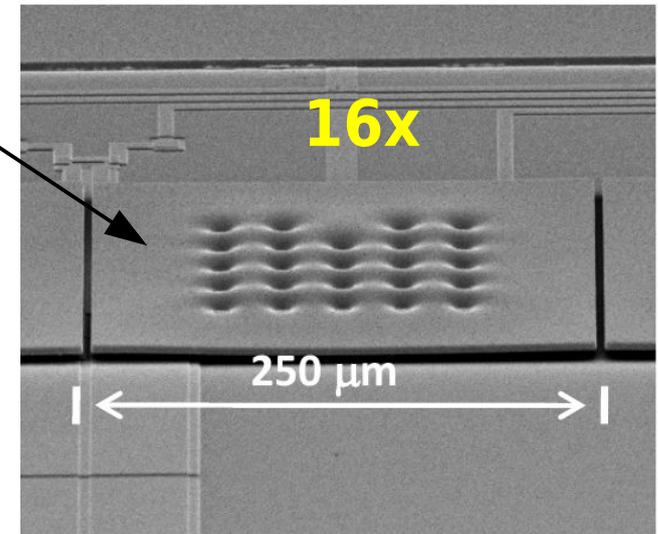


- ▶ high energy resolution $\approx 1\text{eV}$
- ▶ fast response $\approx 1\mu\text{s}$
- ▶ large multiplexable array ≈ 1000

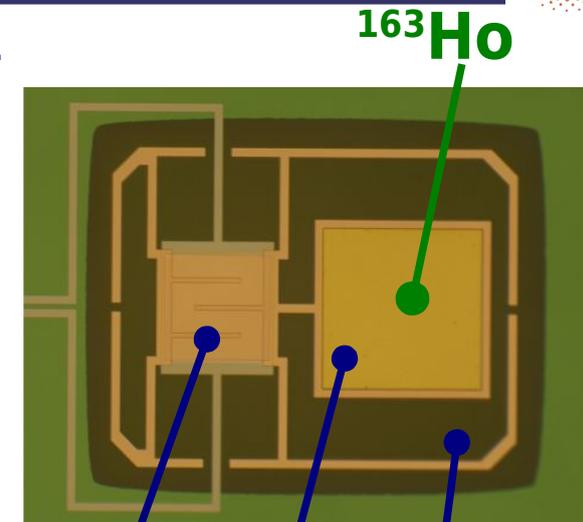
$$\propto \sqrt[4]{1/N_{ev}}$$

ECHo MMC microcalorimeters

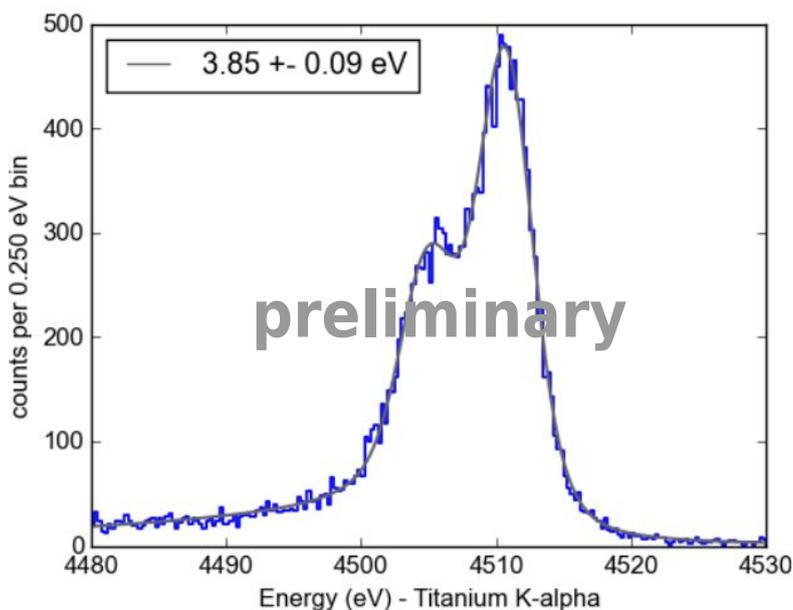
- 16x **MMC** with sandwiched **Au** absorber
- chemically purified ^{163}Ho
- offline ^{163}Ho implantation at ISOLDE (CERN)
- $A(^{163}\text{Ho}) = 0.1 \text{ Bq/pixel}$
- $\Delta E \approx 5 \text{ eV}$
- $\tau_{\text{rise}} \approx 130 \text{ ns}$
- **ECHo-1k**: 100 pixels, $A(^{163}\text{Ho}) = 10 \text{ Bq/pixel}$



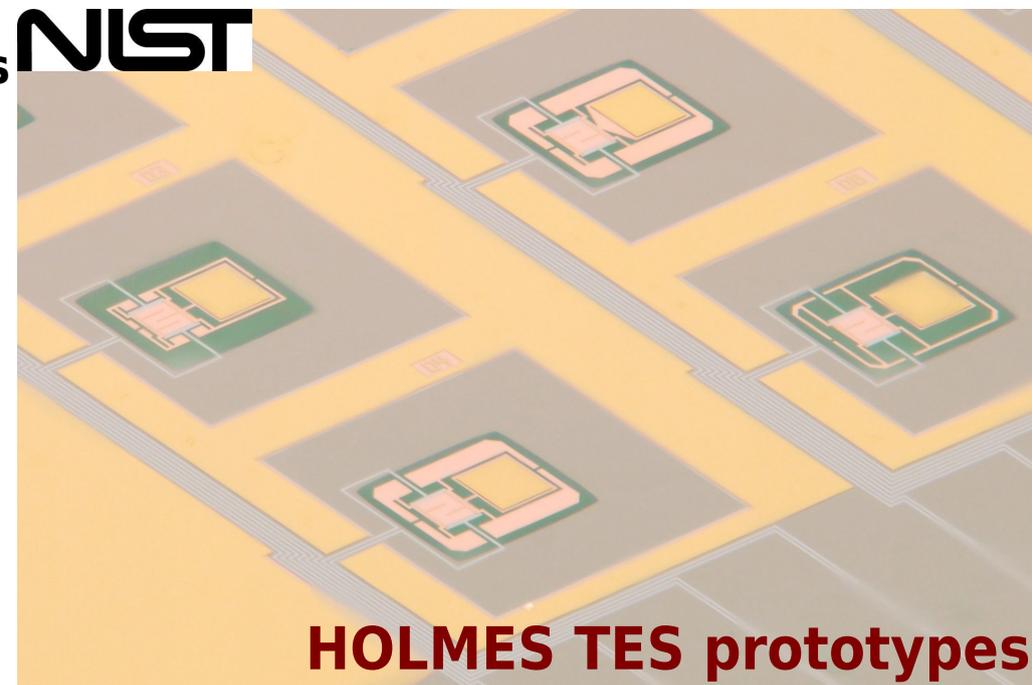
- Transition Edge Sensors (TES) with Au (or Bi) absorber
 - ▷ **2 μm Au** total thickness for *full* electron and X absorption
- MoCu proximity TES $\rightarrow T_c \approx 100\text{mK}$
- on Si_2N_3 membrane
- optimize design for speed and resolution
 - ▷ **specs @2.5keV :**
 - $\Delta E_{\text{FWHM}} \approx 1\text{eV}$
 - $\tau_{\text{rise}} \lesssim 5\mu\text{s}, \tau_{\text{decay}} \approx 100\mu\text{s}$
- from preliminary X-ray measurements:
 - ▷ $\Delta E_{\text{FWHM}} \approx 3\text{eV}, \tau_{\text{rise}} \approx 5\mu\text{s}, \tau_{\text{decay}} \approx 150\mu\text{s}$



TES with Au absorber membrane Si_2N_3

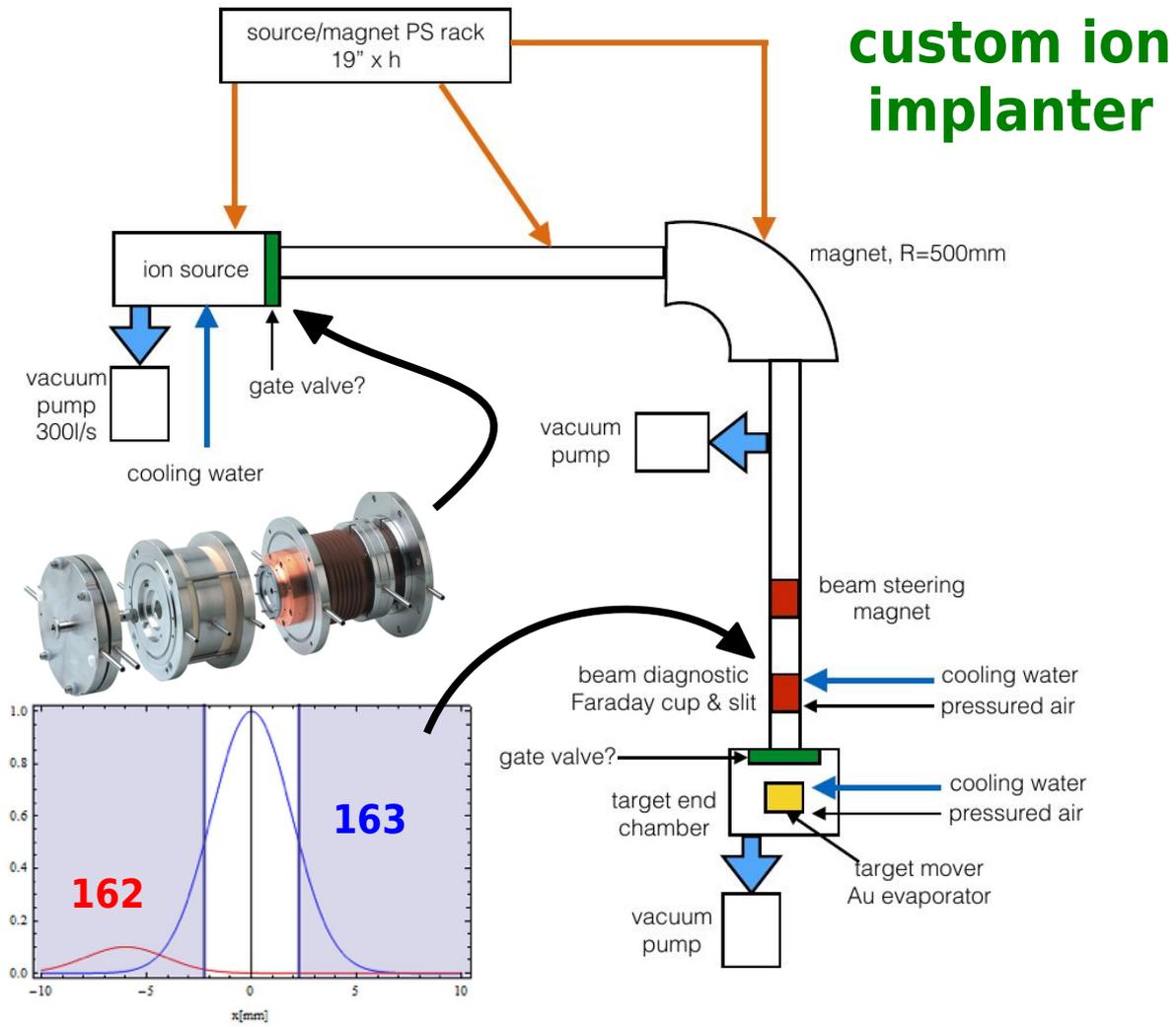


NLST



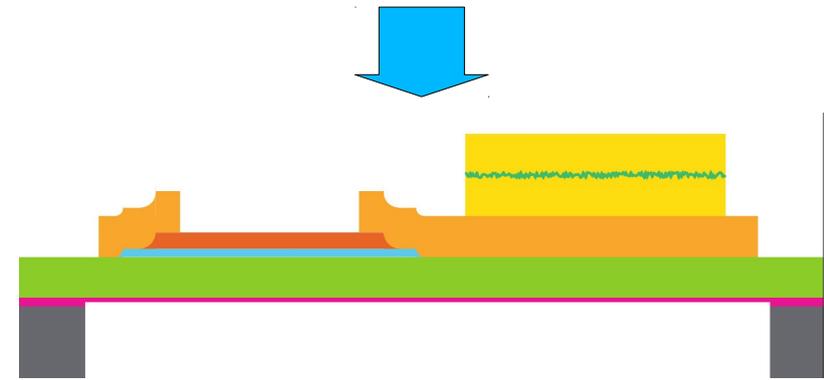
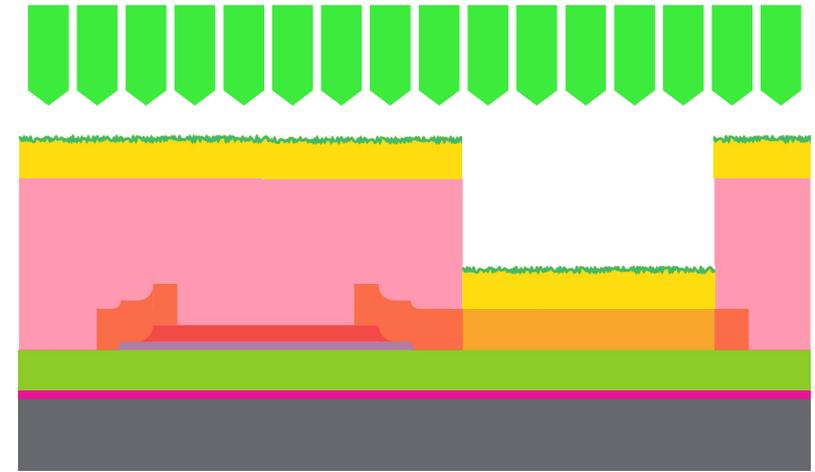
HOLMES TES prototypes

HOLMES detector array fabrication



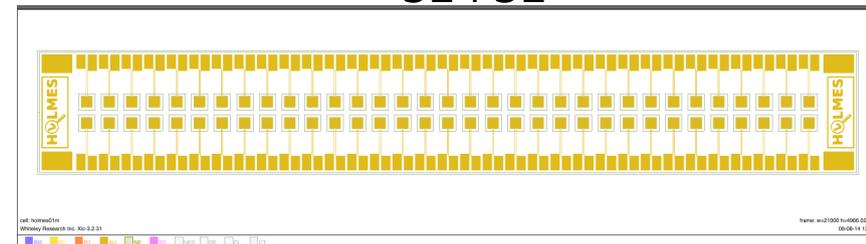
custom ion implanter

^{163}Ho



Si SiO₂ Si₃N₄ Mo Cu Cu Bi Au Ho

32+32



- 2 μm thick **Au** encapsulating implanted ^{163}Ho
- TES fabricated at NIST, Boulder, CO, USA
- ^{163}Ho implantation at INFN Genova
- **HOLMES**: 1000 pixels, $A(^{163}\text{Ho}) = 300 \text{ Bq/pixel}$

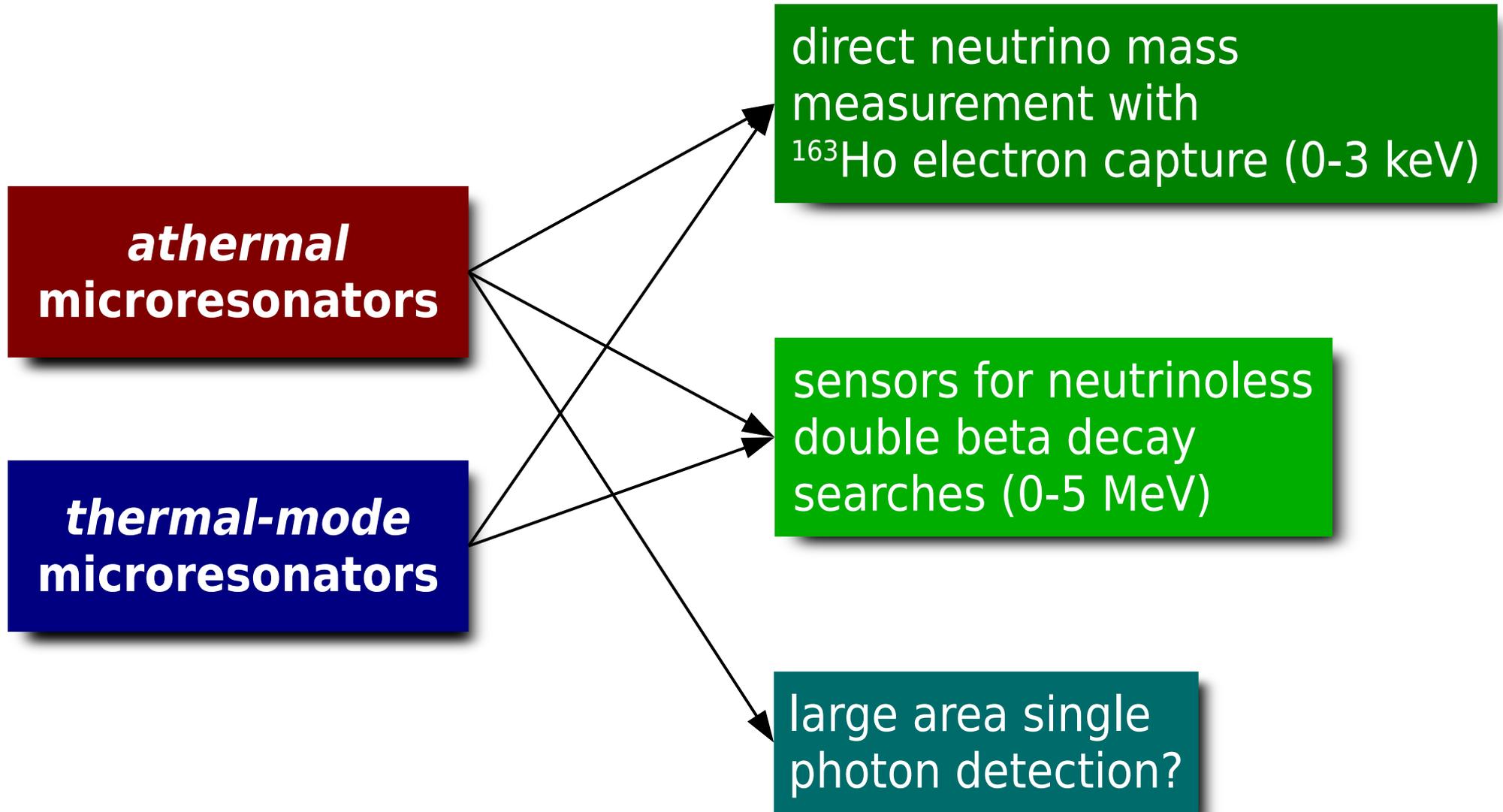
KID: an alternative to TES?

Development of microresonator detectors for neutrino physics

2011 - 2015 @ UniMiB



fondazione
c a r i p l o



KID principles: athermal mode

first developed as radiometers for mm astronomy

$$E = h\nu > 2\Delta (\approx \text{meV})$$

pair breaking detector



- Cooper pair breaking
- quasi-particle creation ($N_{qp} \approx \eta h\nu / \Delta$)
- increase qp density n_{qp}
- change in complex surface impedance $Z_s = R_s + i\omega L_s$

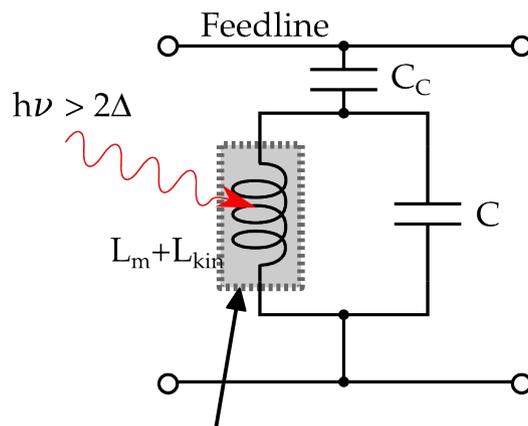
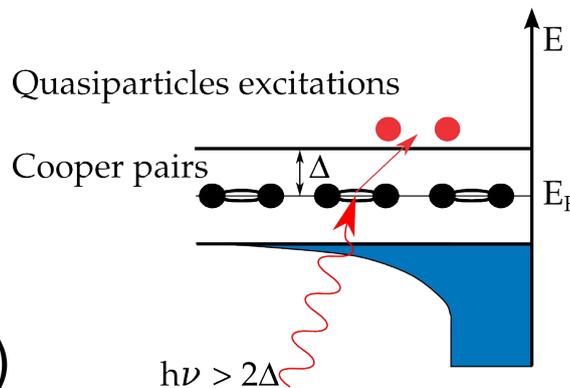


$$\frac{\delta f}{f_0} = -\frac{\alpha}{2} \frac{\delta L_s}{L_s} \quad \delta Q^{-1} = \alpha \frac{\delta R_s}{\omega L_s}$$

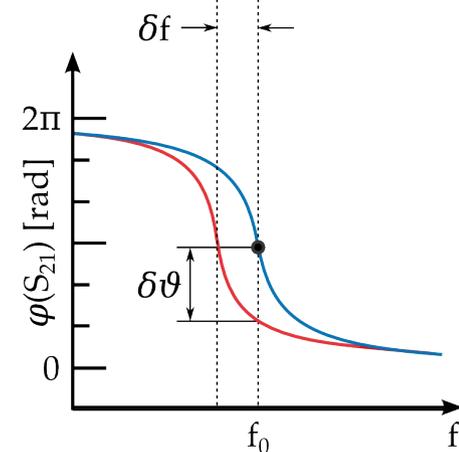
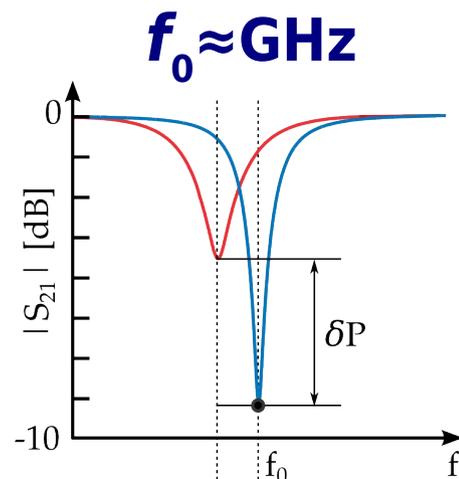
α surface inductance L_s fraction in circuit inductance



relaxation time: qp recombination τ_{qp}



detector



P. Day et al., Nature, 425 (2003) 817

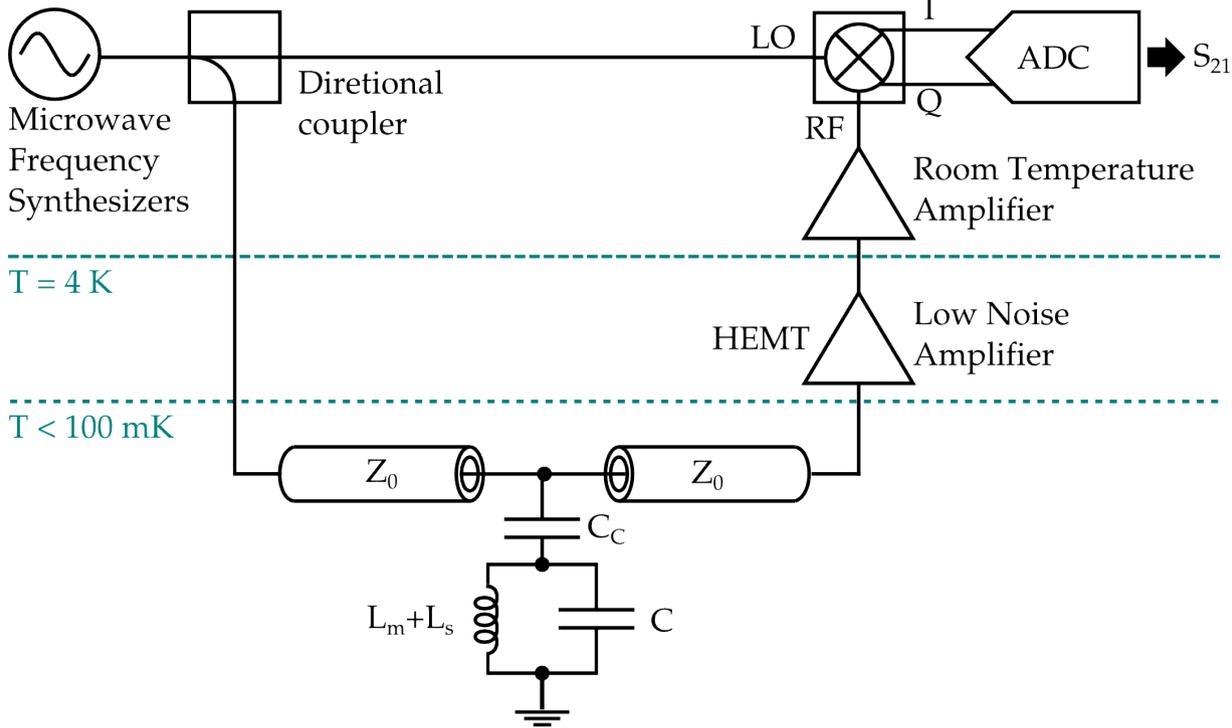
resolution limited by

qp statistics $\sigma_E^2 \approx FE\Delta/\eta$

(order of 1eV for $E = 6 \text{ keV}$)

KID homodyne read-out

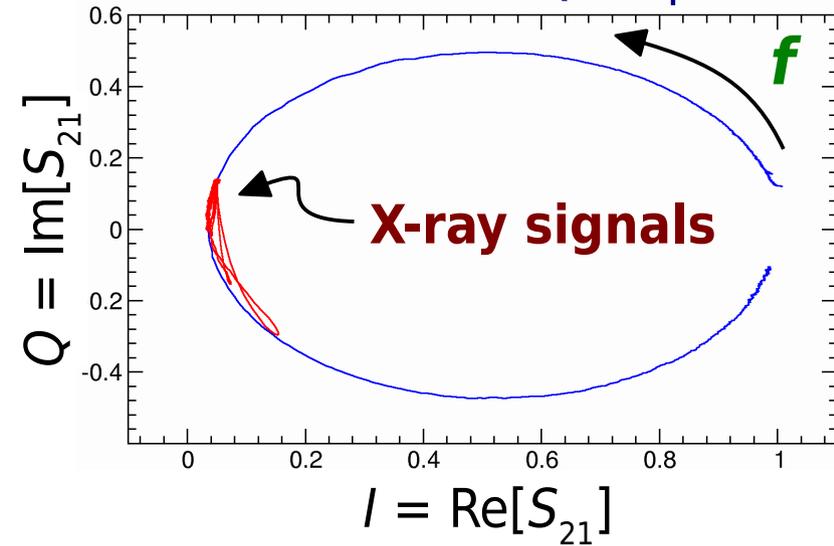
T = 300 K



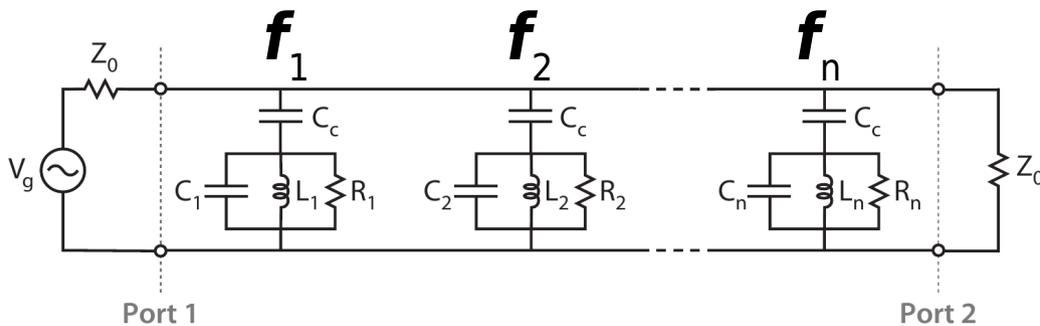
T = 4 K

T < 100 mK

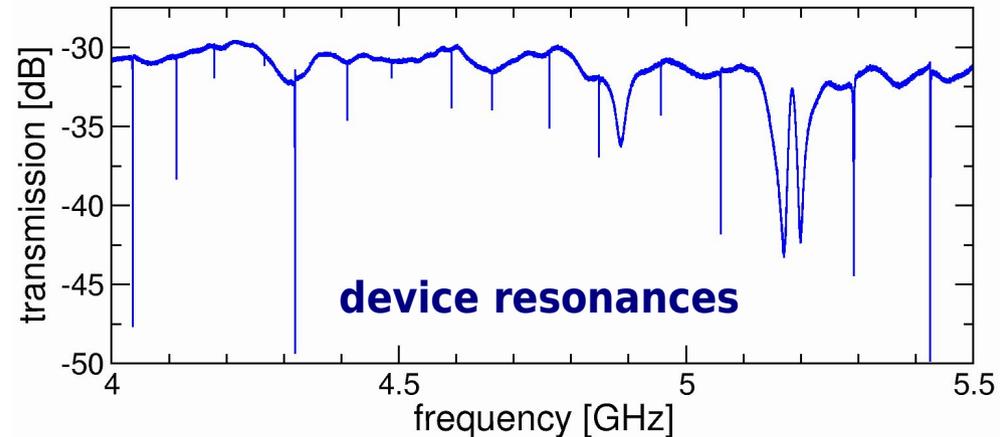
resonance IQ loop



microwave frequency multiplexing



device i is read out tuning the RF carrier to f_i



Advanced materials for KIDs at FBK

- **sub-stoichiometric TiN_x**

- ▶ tunable T_c (0 → 4.5K) by adjusting $x < 1$ → longer τ_{qp} for lower T_c
- ▶ low losses → high Q_i devices
- ▶ no surface oxide → low excess (TLS) noise
- ▶ high surface inductance fraction α → large signals
- ▶ hard to produce in a controlled way

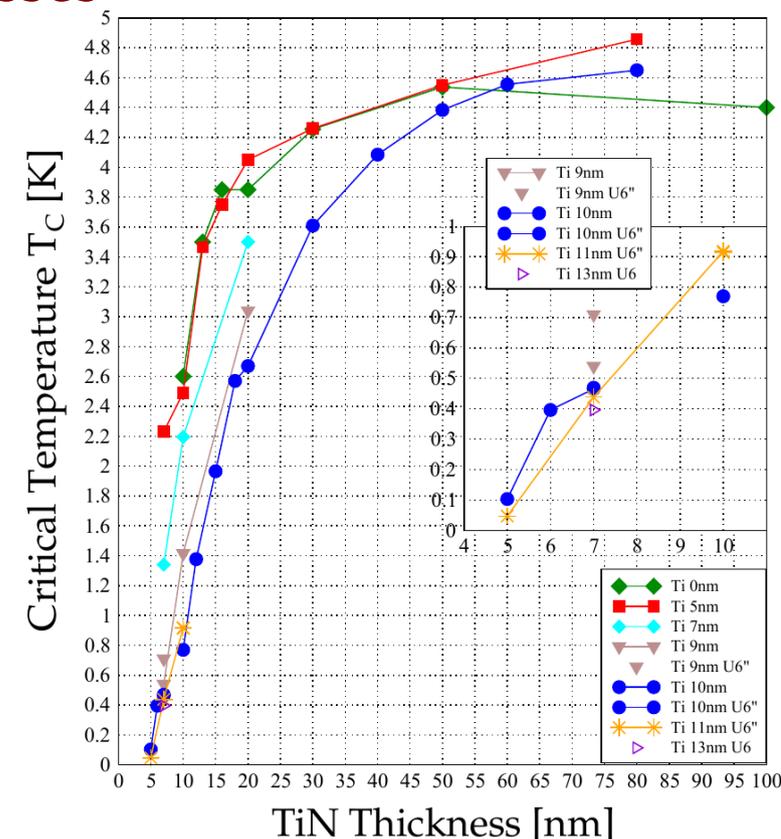
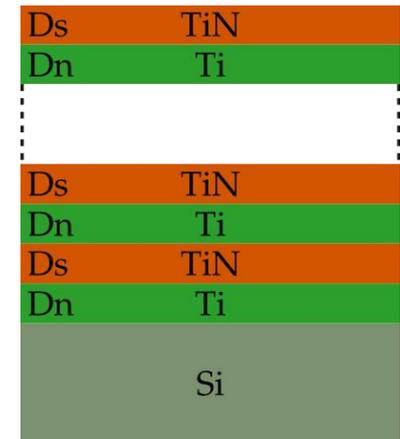
- **Ti/TiN sputtered multilayer (stoichiometric TiN)**

- **proximity effect** → T_c tuned by Ti and/or TiN thicknesses

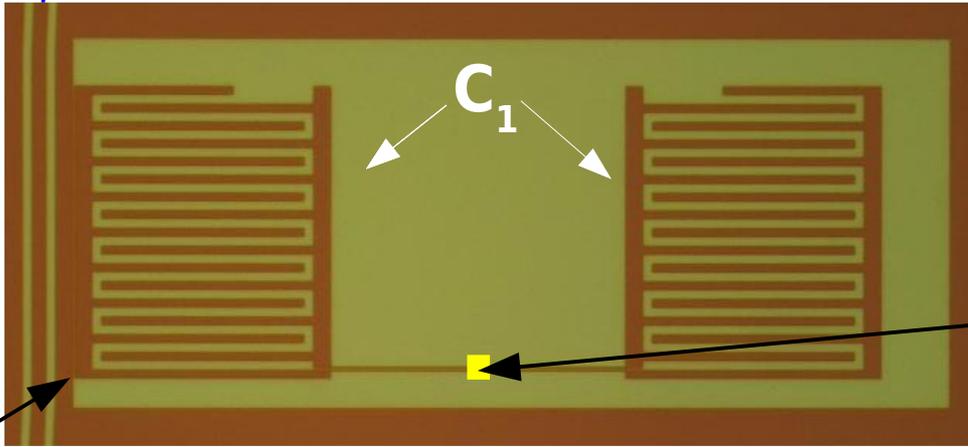
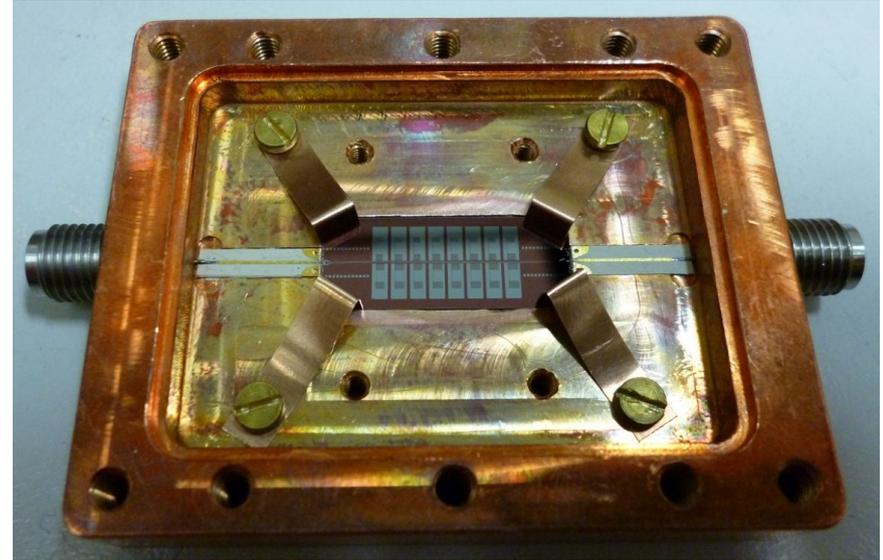
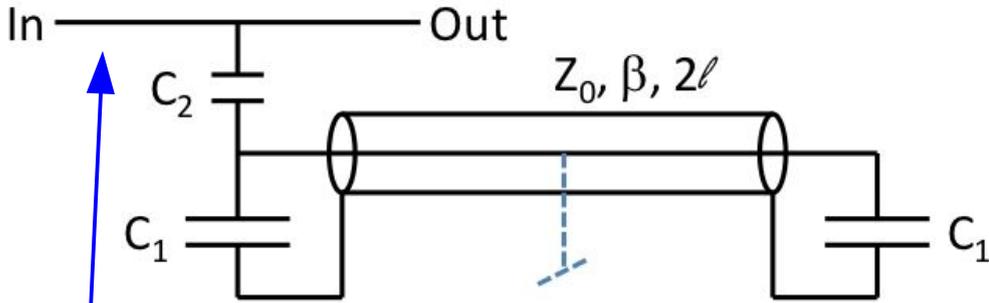
- ▶ good T_c control: tuning range **0.5 → 4.5K**
- ▶ **high Q_i** and **high $\alpha = L_s / L_{tot}$**
- ▶ good reproducibility and uniformity
- ▶ **equivalent to sub-stoichiometric TiN_x**

A. Giachero et al. J. Low Temp. Phys. 176 (2014) 155

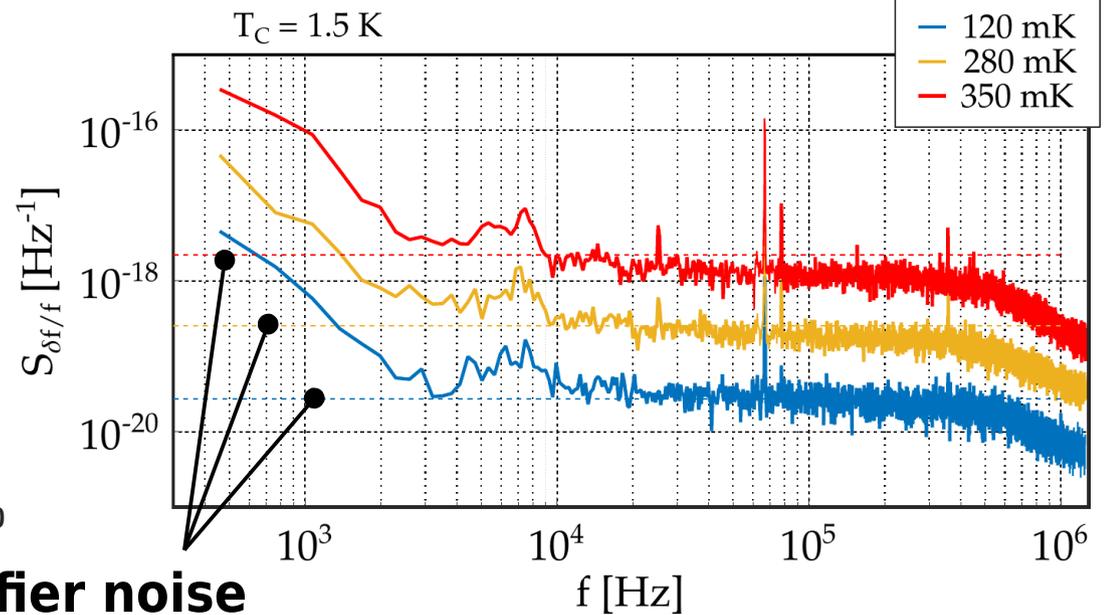
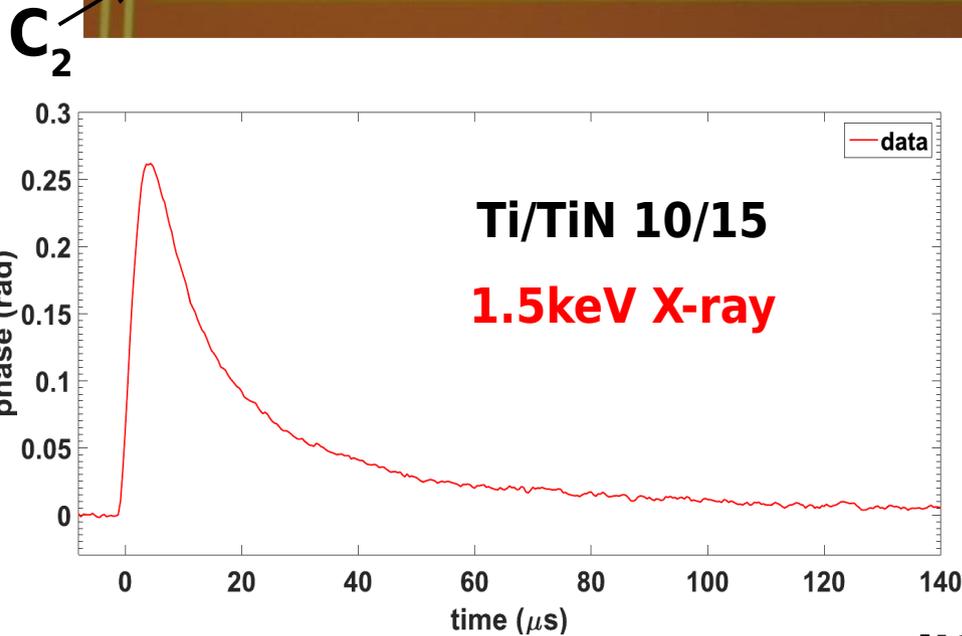
T_c [K]	Δ [meV]	α	L_s [pH/sq]	Q_i
1.5	0.200 ± 0.004	0.26 ± 0.01	13.1	$< 10^5$
0.640	0.091 ± 0.001	0.95 ± 0.01	36.9	$< 10^4$



Microresonator X-ray characterization

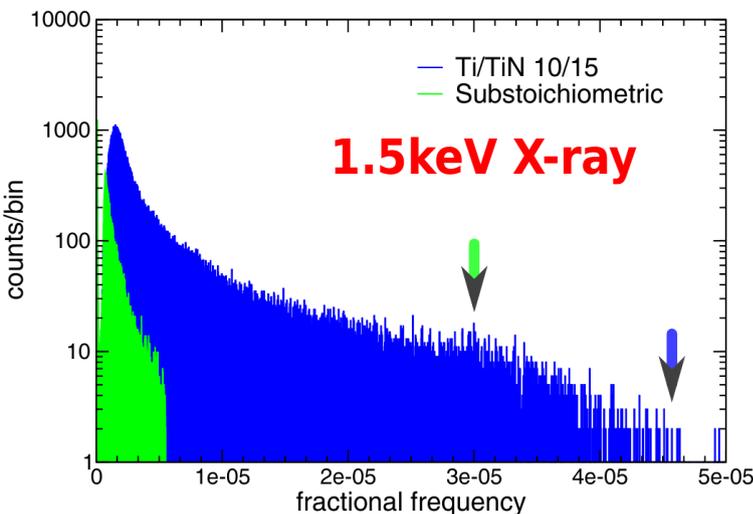


¹⁶³Ho implant
not yet implemented...
→ X-rays irradiation

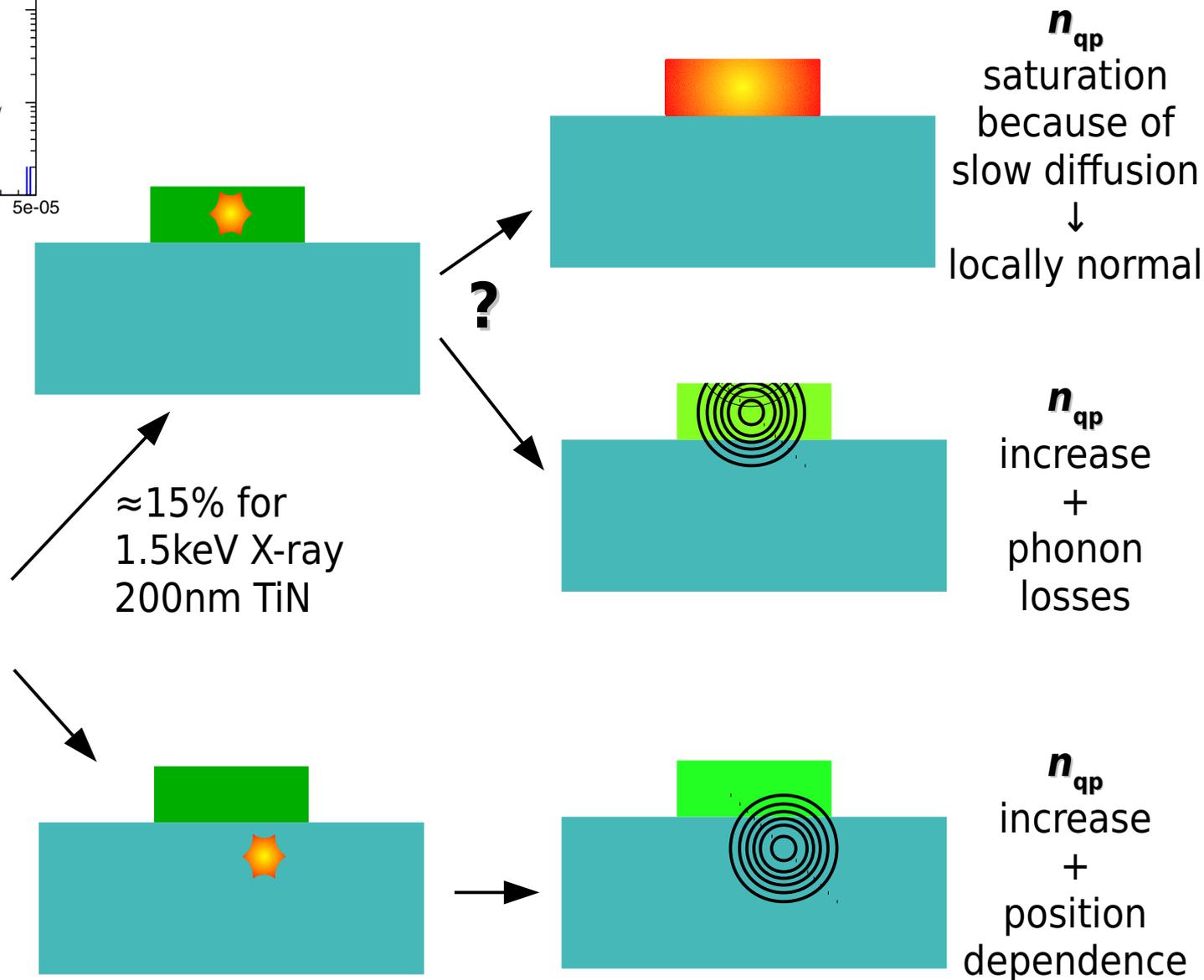
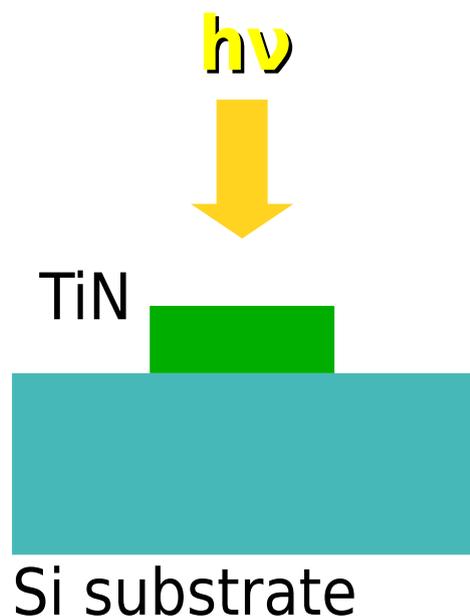


amplifier noise

Microresonator X-ray response

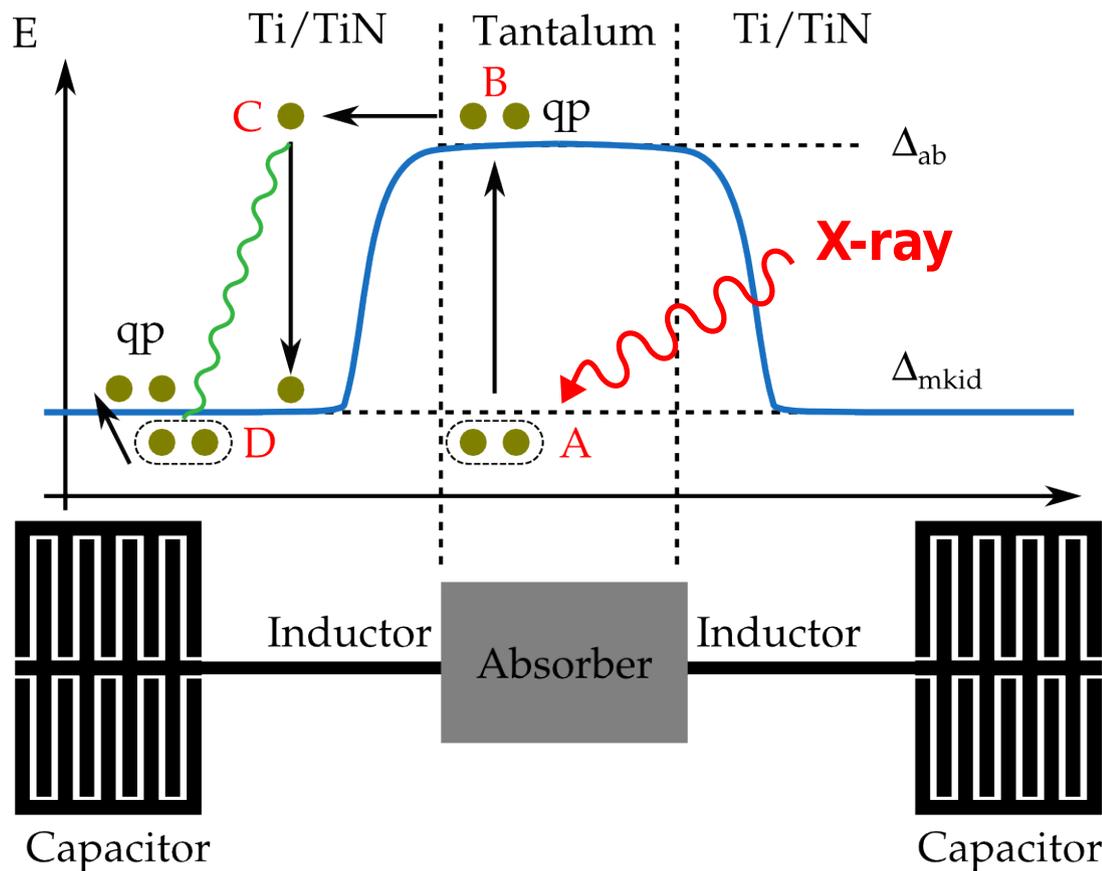


large (>1000) S/N but no resolving power!
HEMT noise dominated S/N

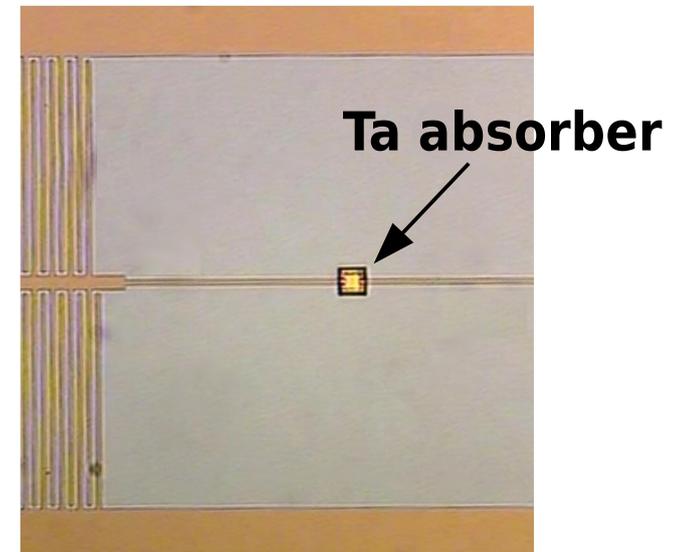


Microresonator detector improvements

- fast **qp diffusion** in Tantalum \rightarrow no saturation
- **qp trapping** in TiN/Ti: $\Delta(\text{Ta}) > \Delta(\text{TiN/Ti})$
- Tantalum: high $Z \rightarrow$ **high stopping power**
- Tantalum thickness: 200 / 500 nm

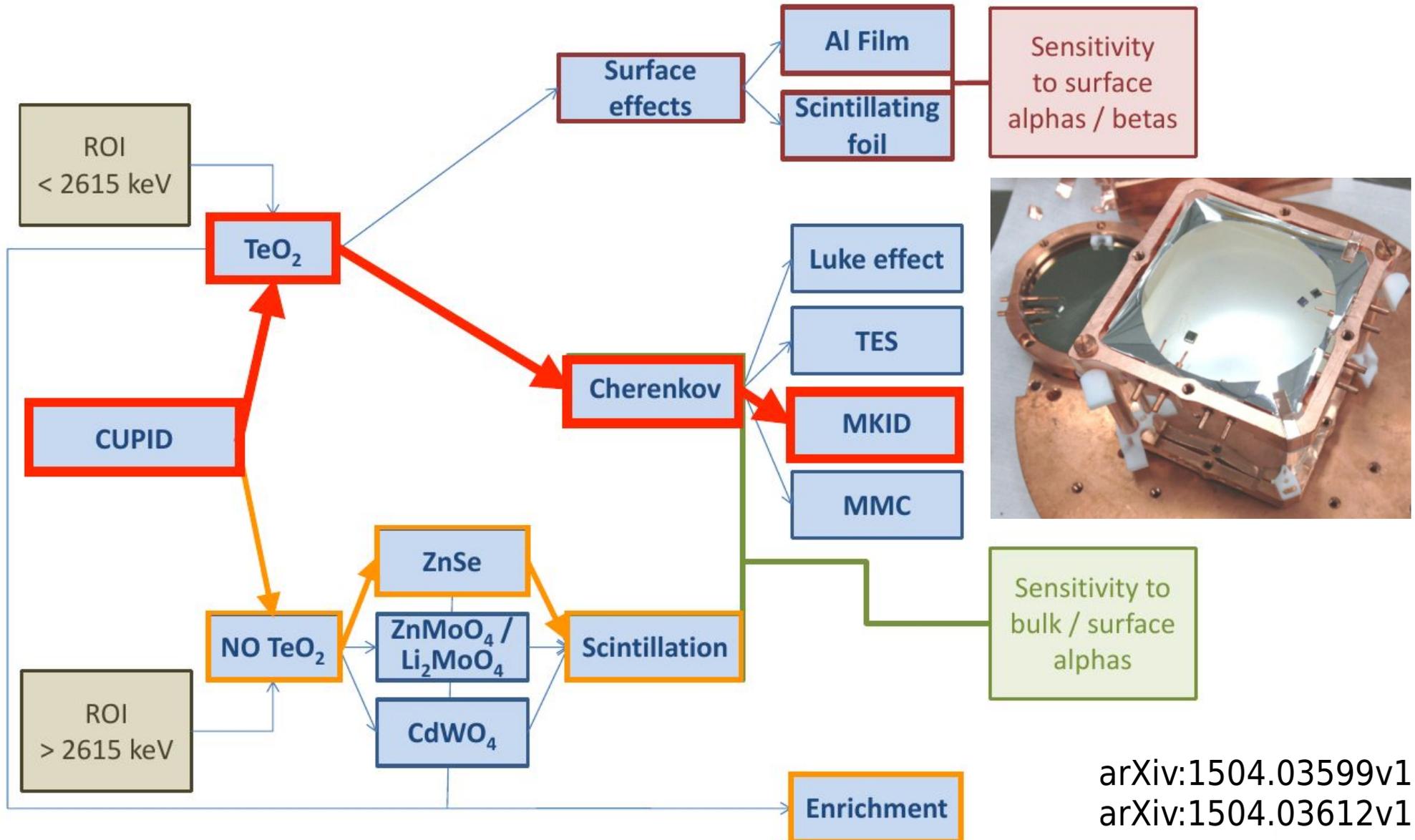


first prototype

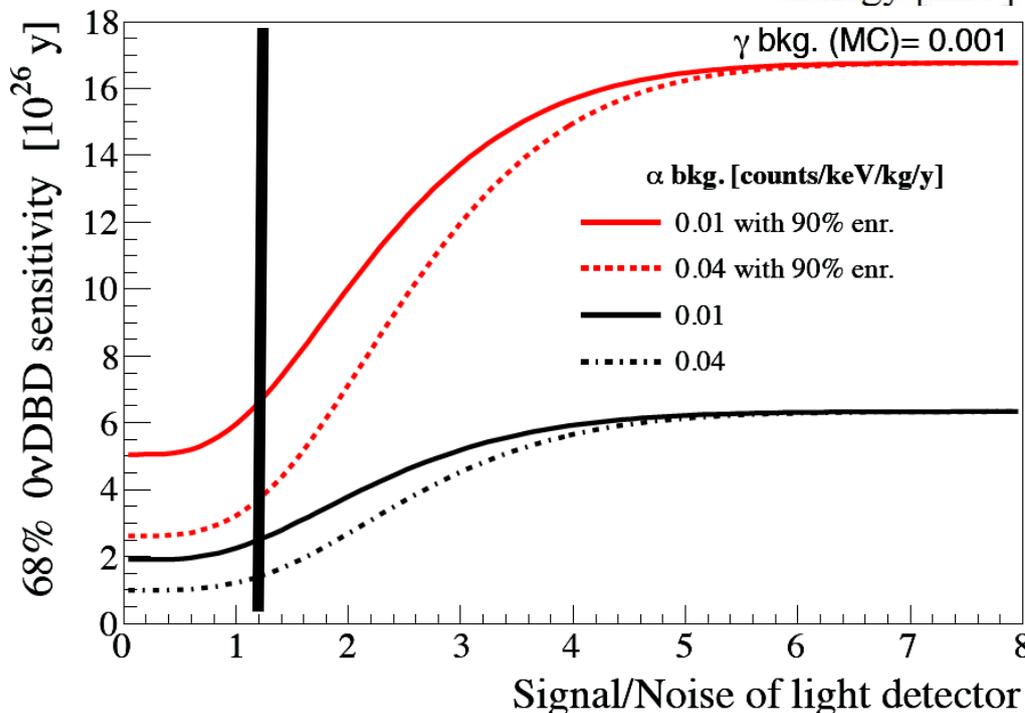
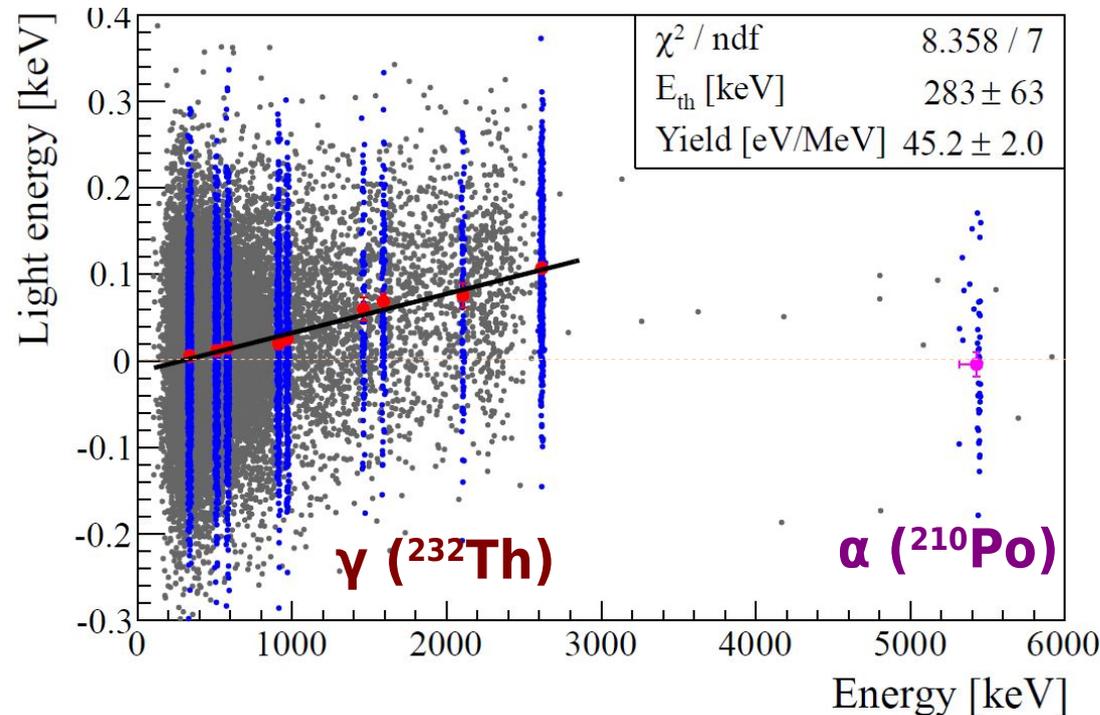


Beyond CUORE: CUPID

- Excluding **Inverted Hierarchy** → improve **CUORE** sensitivity by 5÷10
- CUORE upgrade with Particle ID (**CUPID**) → **reduce bkg from alphas**
 - ▷ **α/γ discrimination with Cherenkov** (T.Tabarelli de Fatis, EPJC65 (2010) 359)



CUPID: TeO₂ with Cherenkov



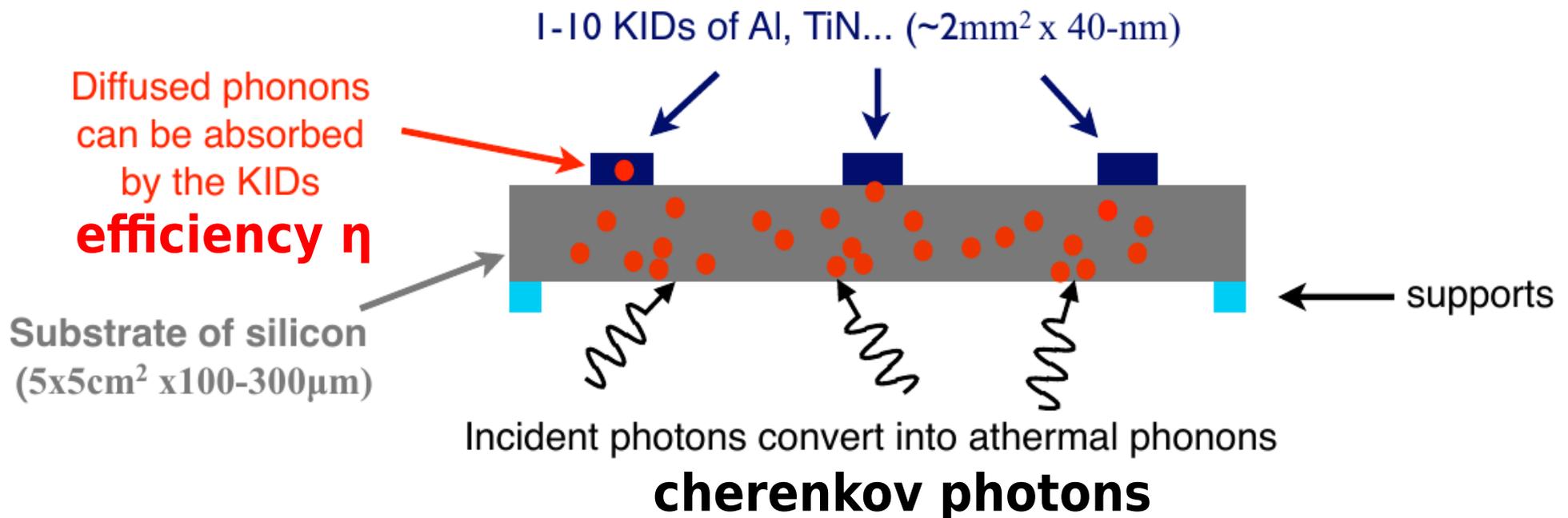
TeO₂ 770g (CUORE size)

- **Detected light not sufficient for event by event discrimination**
 - LY \approx 45 eV/MeV
 - \approx 100 eV for β/γ @ $Q_{\beta\beta}$ (S/N \approx 1.2)
- **10 times less than expected**
- **Montecarlo simulation**
 - Cherenkov light is self absorbed
 - more sensitive light detectors
- **5 year sensitivity to ¹³⁰Te $\beta\beta$ -0 ν**
 - CUORE + Cherenkov detection
 - with/without enrichment
 - **γ bkg (10^{-3} c/keV/kg/y) dominates for S/N \geq 5**

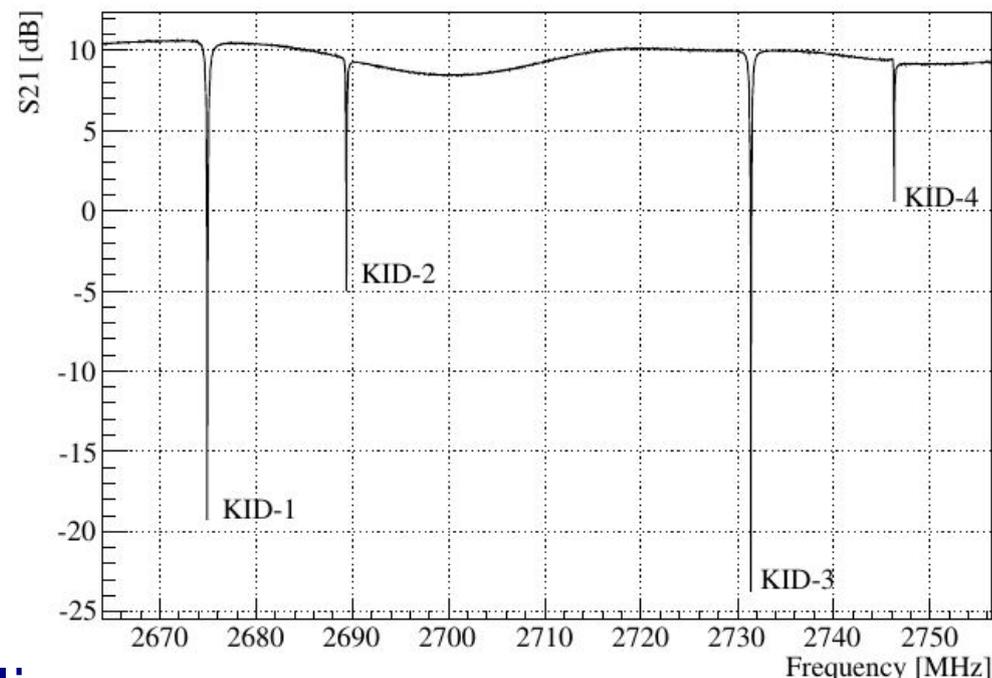
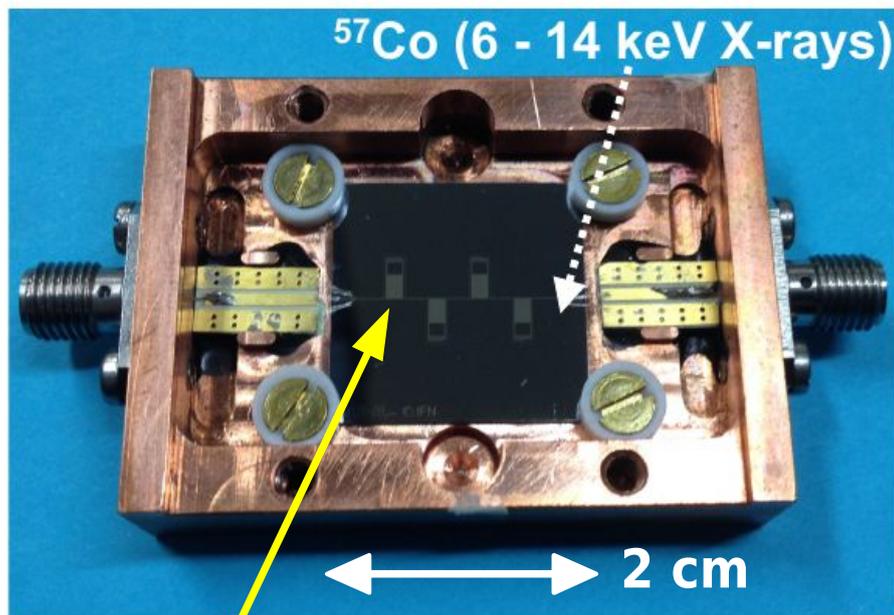
N. Casali et al., EPJ C75 (2015) 12
arXiv:1403.5528v1;



- **CALDER Target**
 - ▶ 50x50 mm² light detector
 - ▶ $E_{\text{thres}} = 20 \text{ eV}$
 - ▶ $T_{\text{op}} = 10 \text{ mK}$
- **Baseline detector design**
 - ▶ athermal phonon KID
 - ▶ Si substrate with up to 10 KID

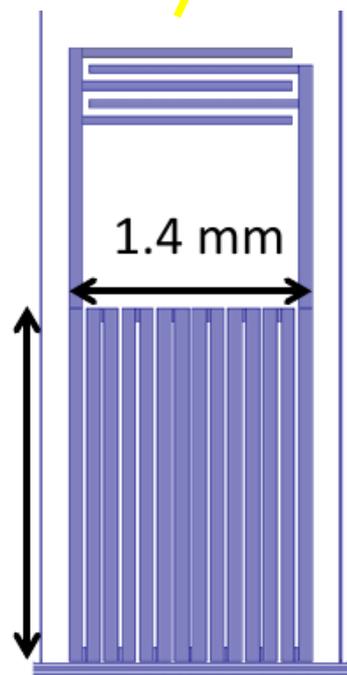
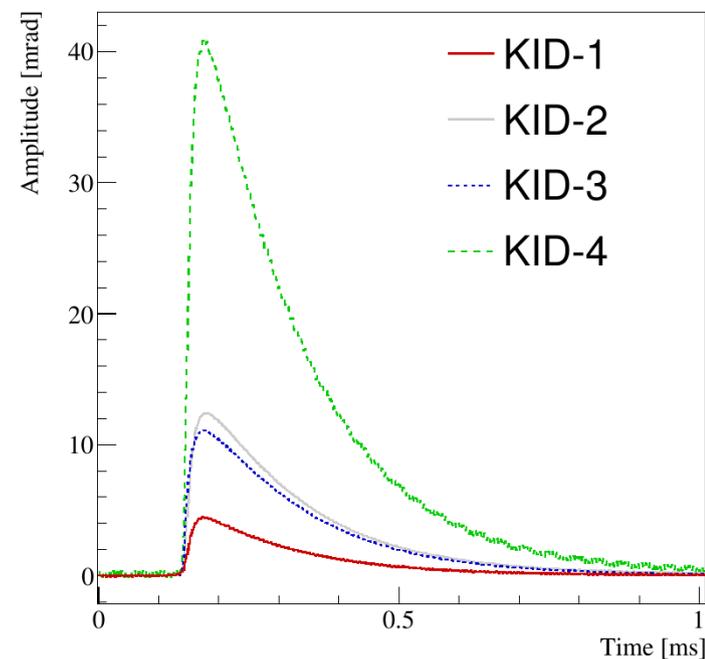


CALDER prototype detectors



- **4x AI KID on Silicon**
 - ▶ pair-breaking detectors
- **multiplexed read-out**
 - ▶ Software Defined Radio (NIXA)

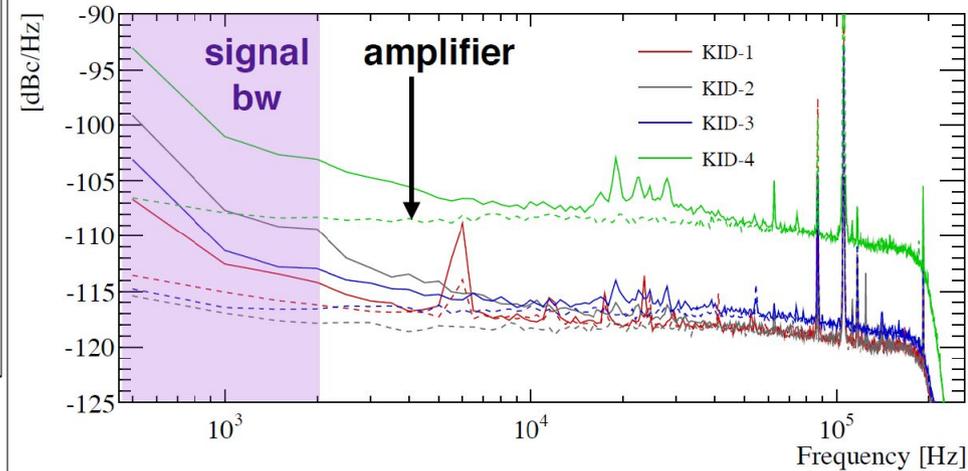
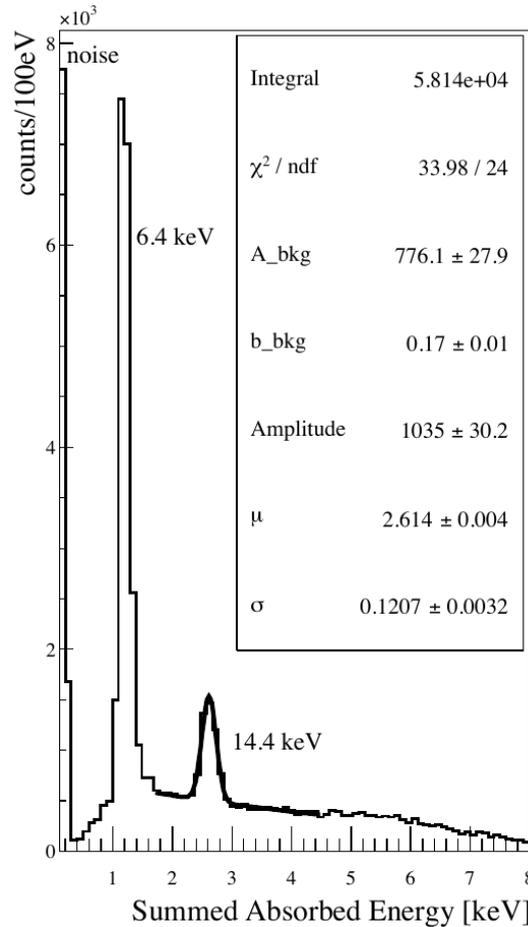
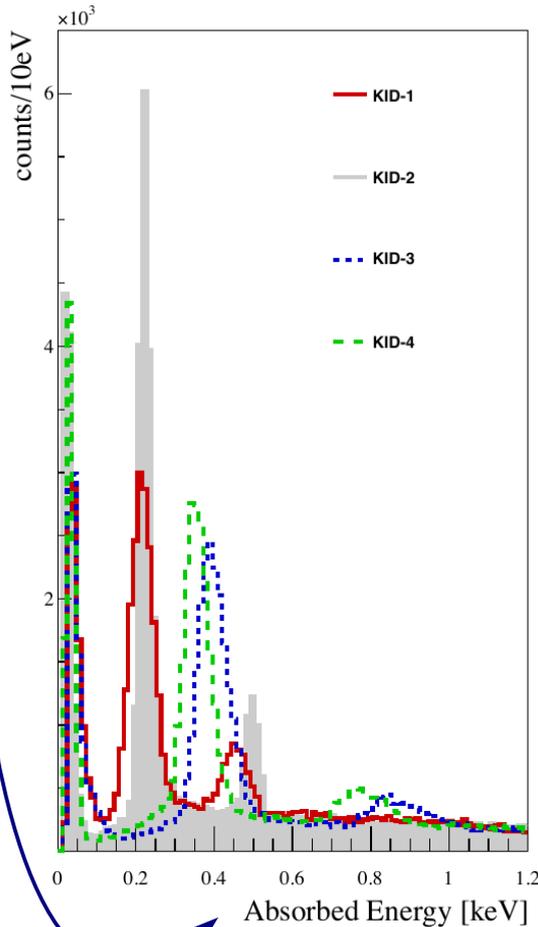
	f_0 [GHz]	Q [$\times 10^3$]
KID-1	2.675	6
KID-2	2.689	18
KID-3	2.731	8
KID-4	2.746	35



CALDER results

- KID response calibrated in energy
- summing up the energy absorbed by each KID
 - ▶ phonon collection efficiency $\eta = 18\%$
 - ▶ baseline width $\sigma_E = 154 \text{ eV}$
 - ▶ unknown **low f excess noise** dominates S/N

$$E = n_{qp} \Delta = \frac{\delta f_0 / f_0}{p_0} \Delta$$

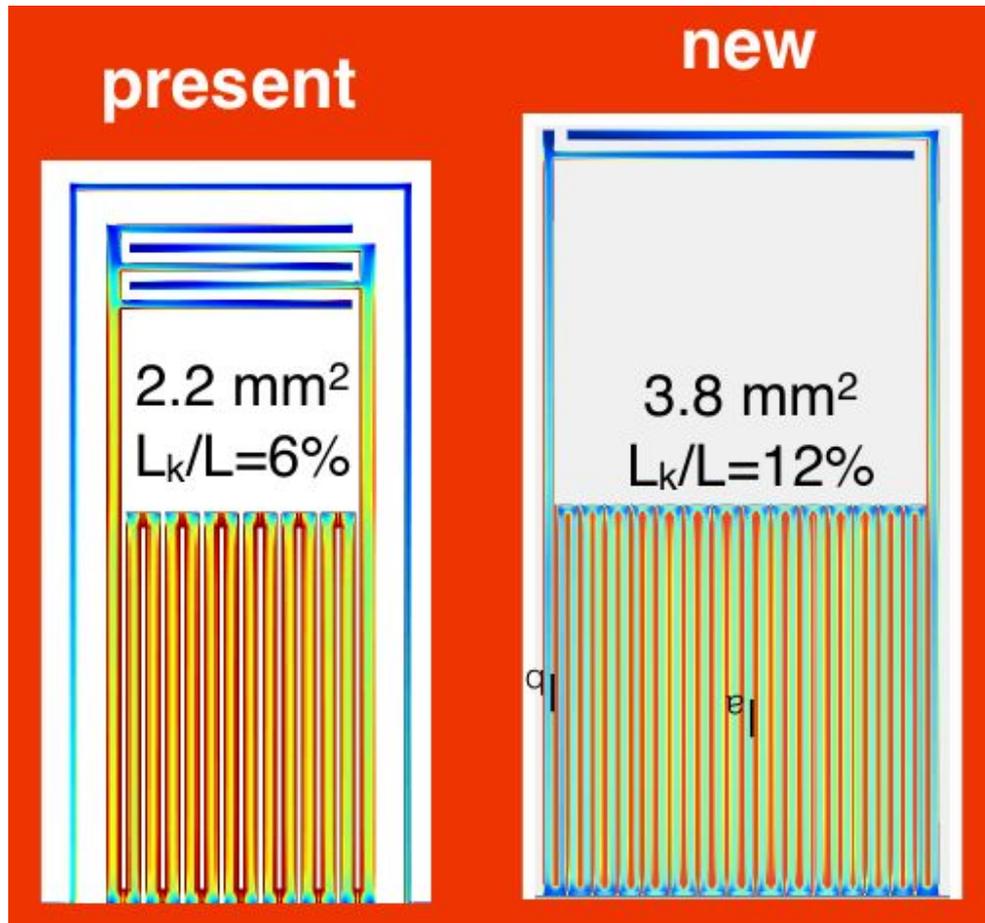


L. Cardani et al. Appl. Phys. Lett., 107 (2015) 093508

CALDER planned improvements

$$\Delta E \propto \frac{T_c}{\sqrt{Q \eta L_k / L_{tot}}}$$

- increase phonon collection efficiency η by design
- increase $\alpha = L_k / L_{tot}$ by design
- increase $\alpha = L_k / L_{tot}$ changing material
 - ▶ non stoichiometric **TiN_x**
 - ▶ **Ti/TiN** multilayers
- reduce **1/f excess noise**



	Al	TiN [non stoic.]	Ti+TiN [stoich]
T_c [K]	1.2	0.9	>0.4
L_k [pH/square]	0.5	3	30

Thermal-mode microresonators (KIDS_RD @CSN5)

Equivalence of temperature change and external pair breaking

- $h\nu$ absorption $\rightarrow \Delta T \approx h\nu/C$
- increase qp density n_{qp}

$$n_{qp}(T) = 2 N_0 \sqrt{2\pi k T \Delta} e^{-\frac{\Delta}{kT}}$$

- change in complex surface impedance $Z_s = R_s + i\omega L_s$

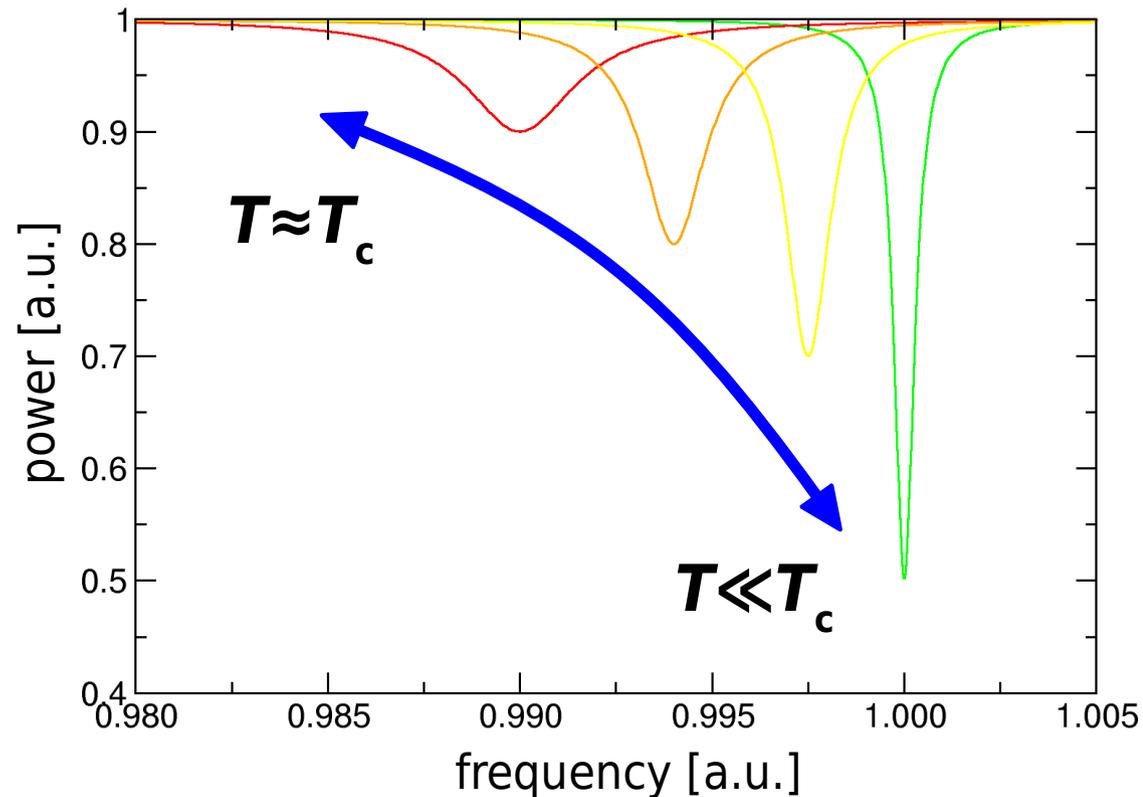


$$\frac{\delta f}{f_0} = -\frac{\alpha}{2} \frac{\delta L_s}{L_s} \quad \delta Q^{-1} = \alpha \frac{\delta R_s}{\omega L_s}$$

α surface inductance L_s fraction in circuit inductance

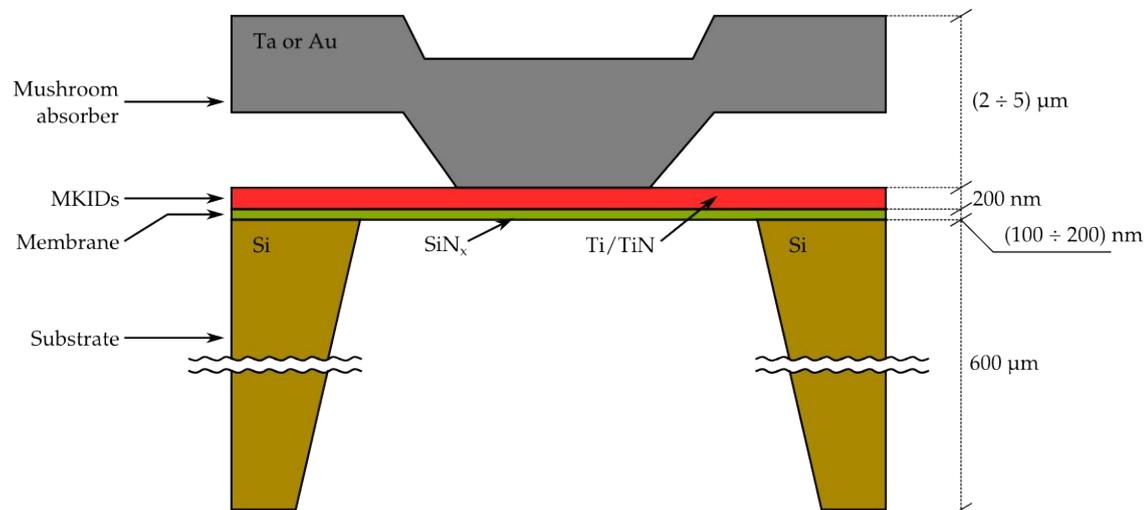
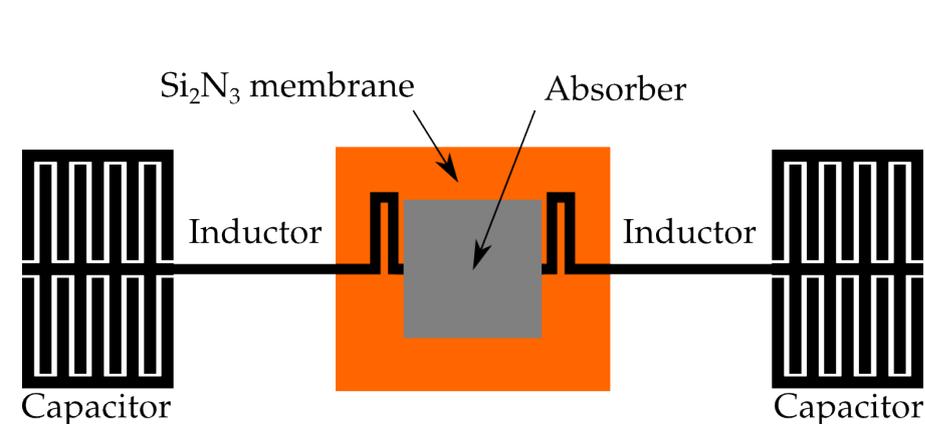
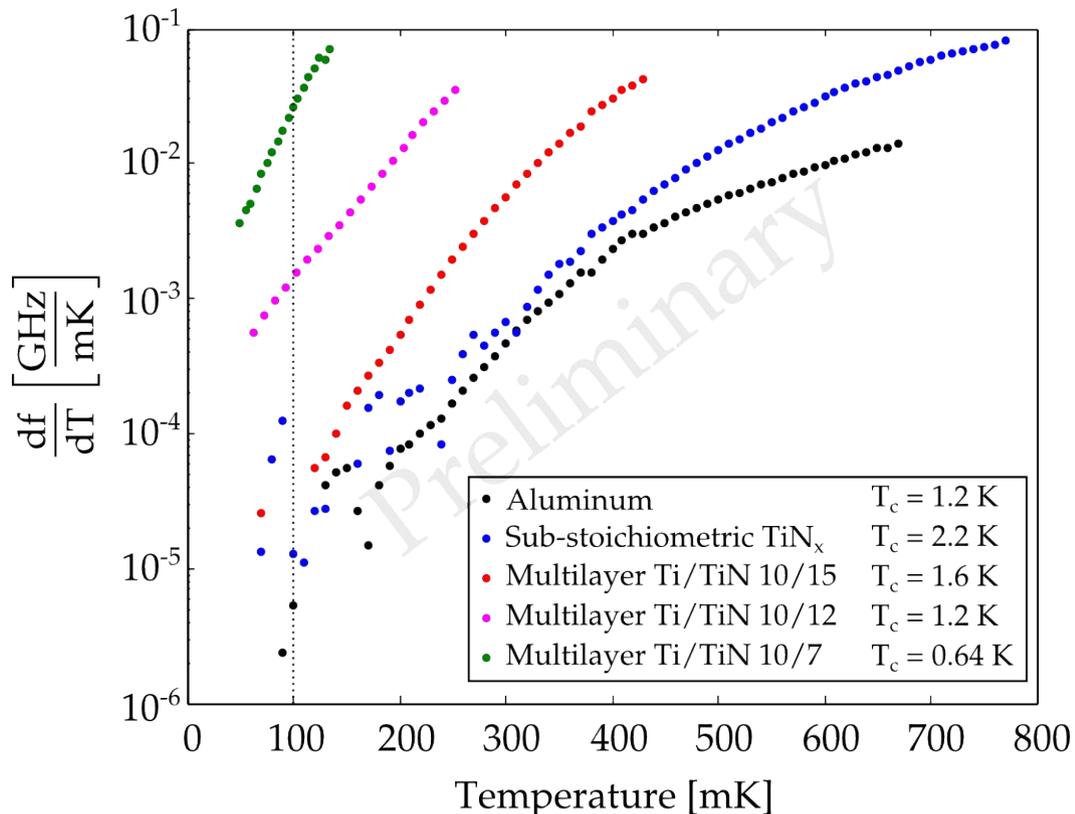
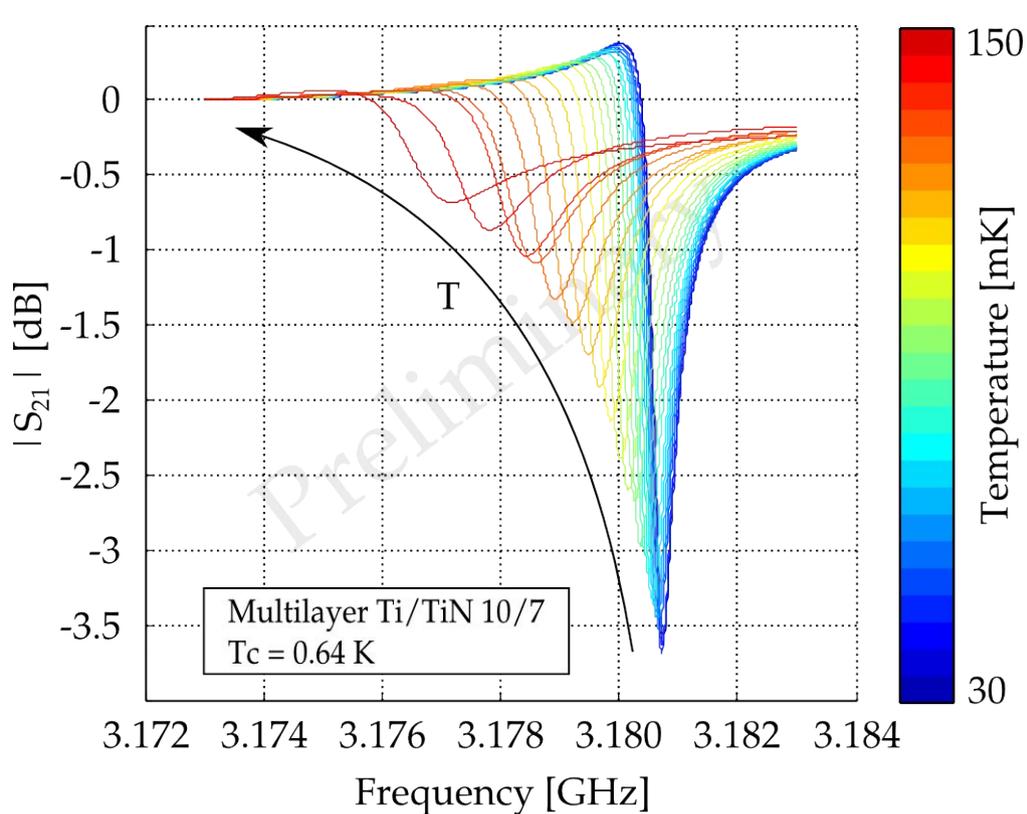


- relaxation time:
thermal time constant $\tau = C/G$

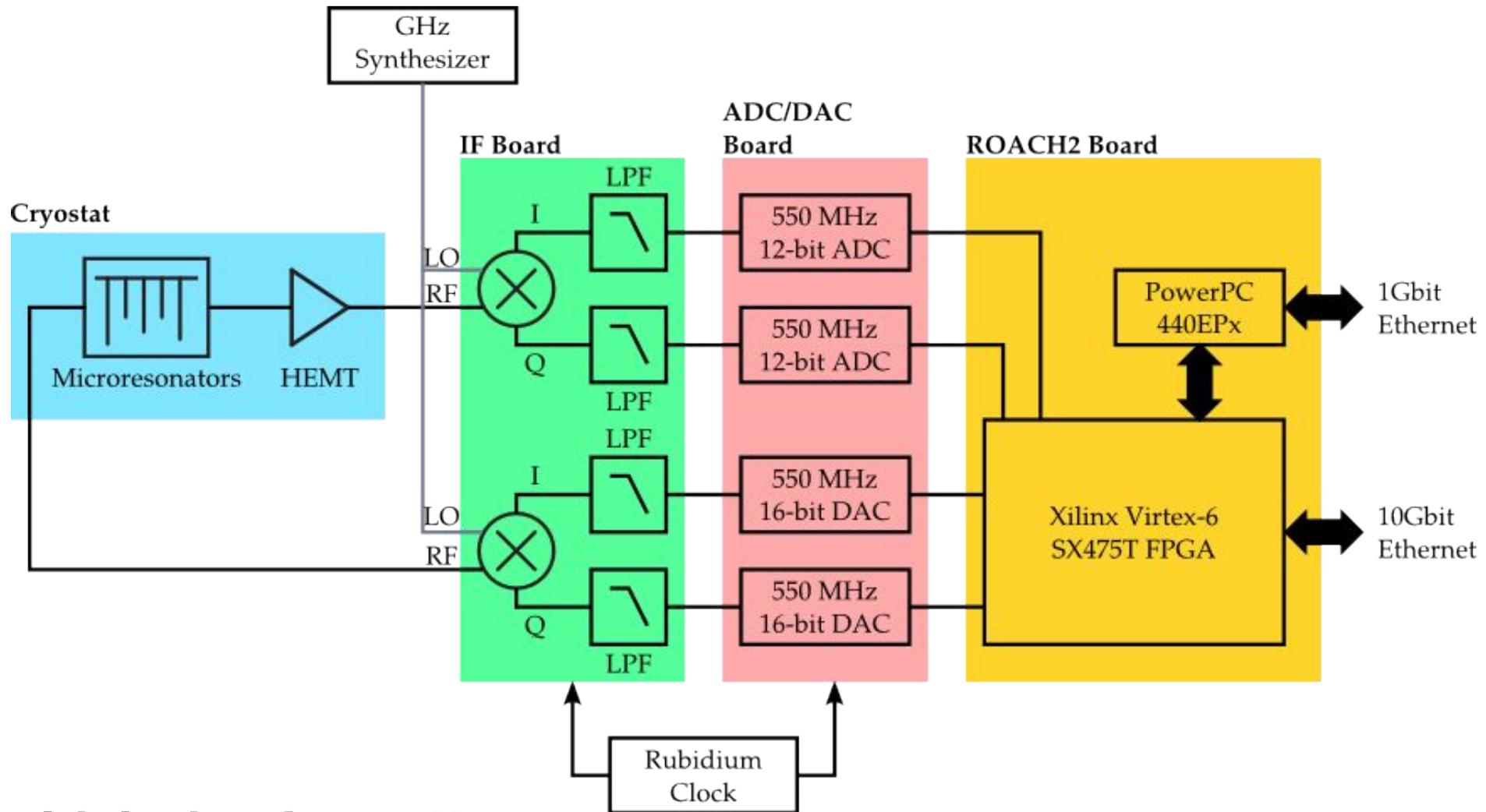


- **thermal KID** \leftrightarrow TES *replacement*
 - ▶ simpler micro-fabrication
 - ▶ simpler read-out (no SQUID)
 - ▶ easier high BW multiplexing
 - ▶ **$\Delta E \approx$ thermodynamic limit**

Thermal-mode KID microcalorimeters



KIDS multiplexing: Software Defined Radio



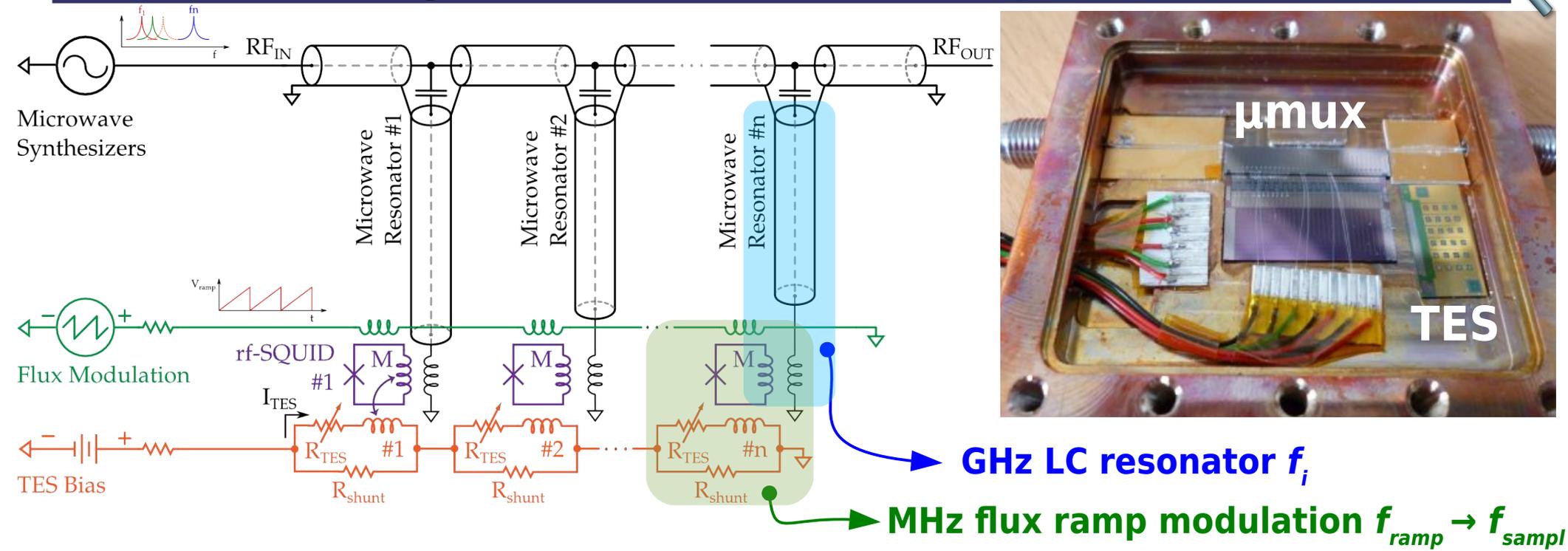
- multiplexing factor N_{mux}

- f_{BW} required bandwidth per channel $\approx 1/\tau_{rise}$

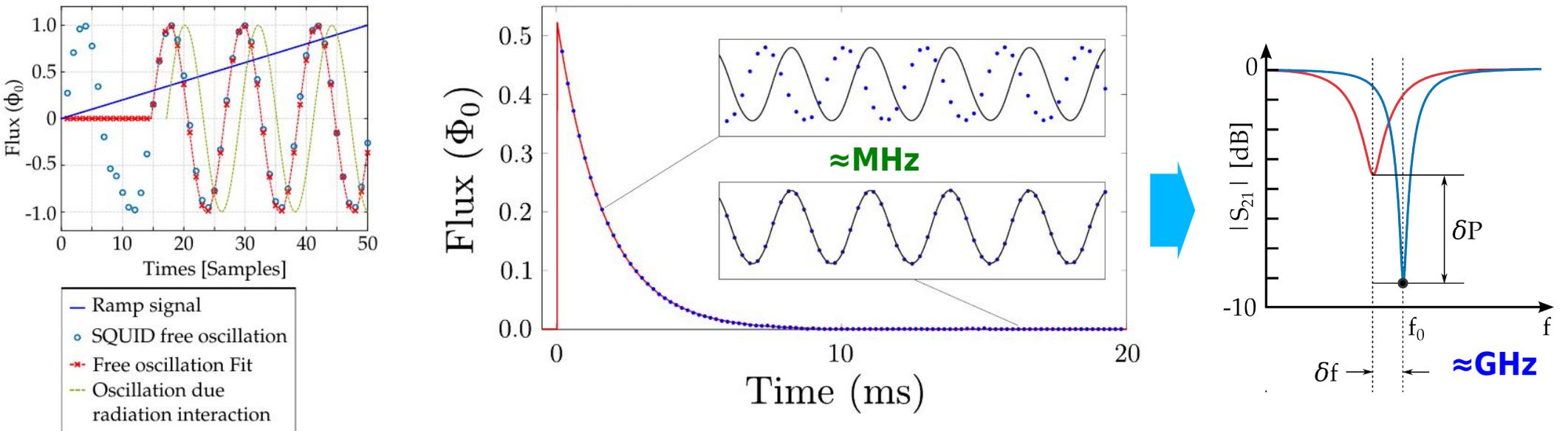
$$N_{mux} = \frac{f_{ADC}}{10f_{BW}}$$

$$f_{ADC} = 550 \text{ MHz}, f_{BW} = 100 \text{ kHz} \rightarrow N_{mux} \approx 500$$

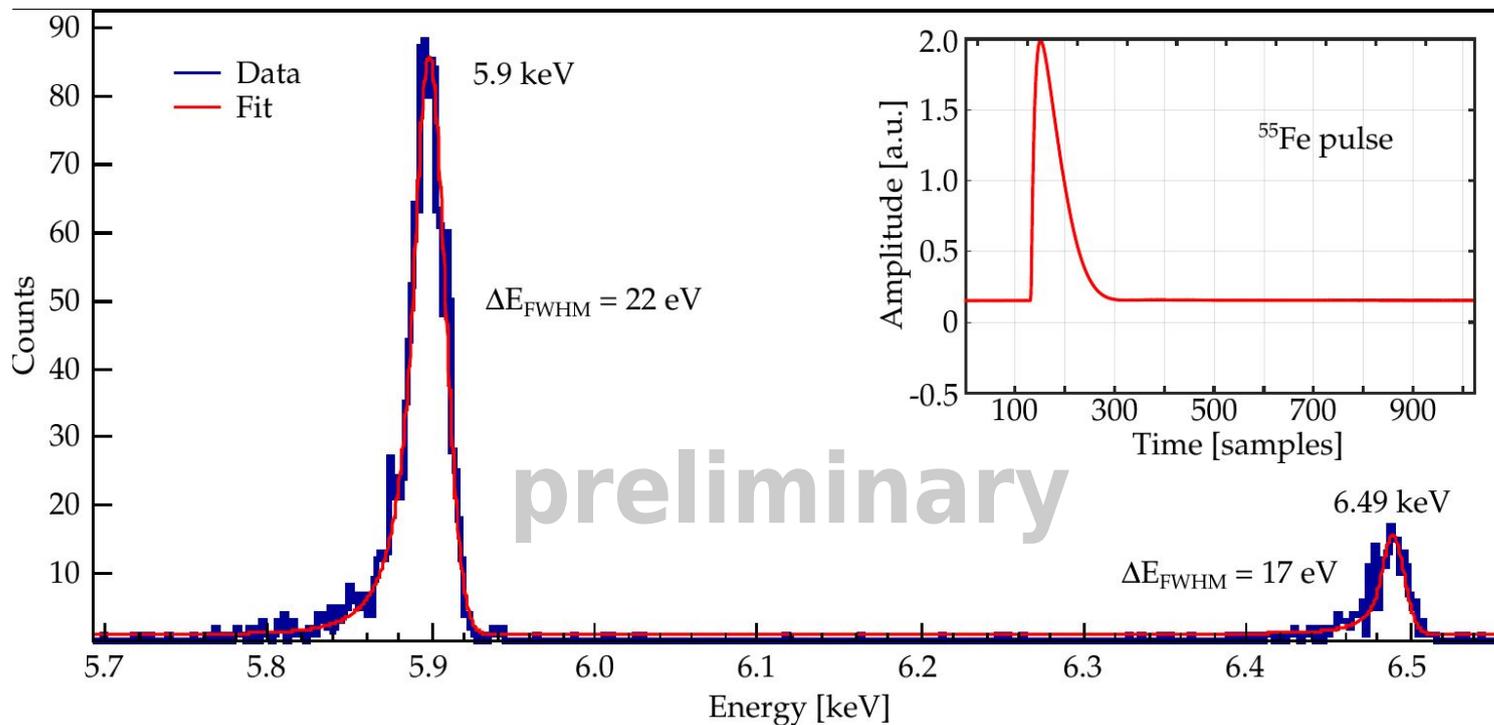
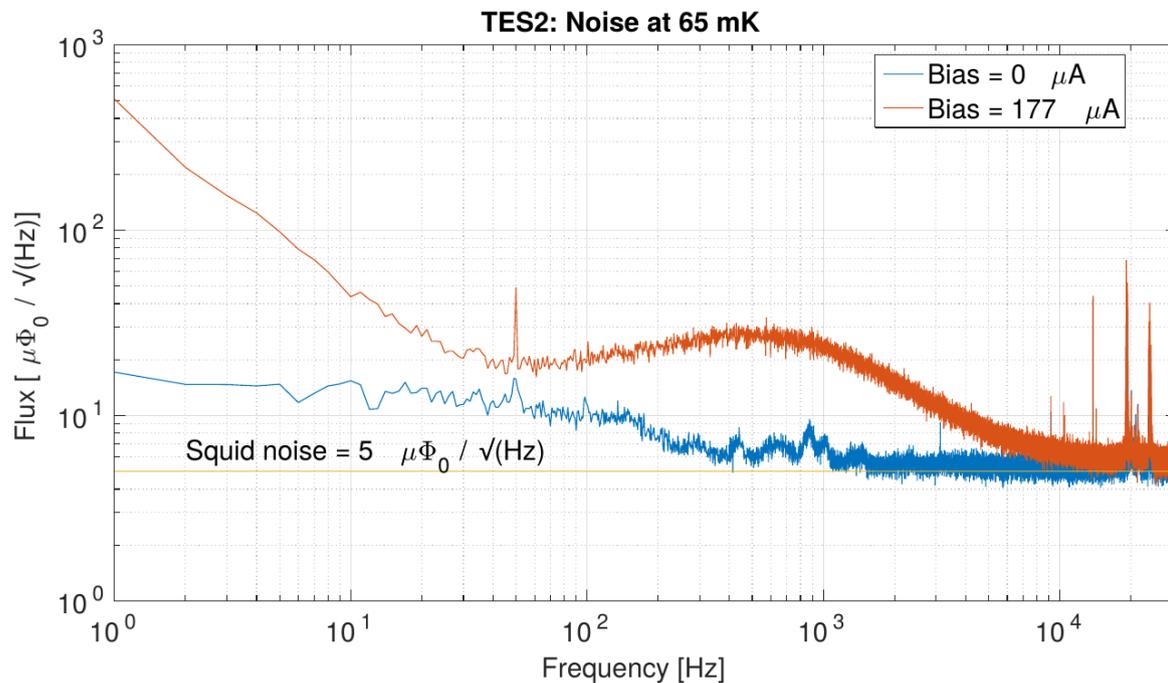
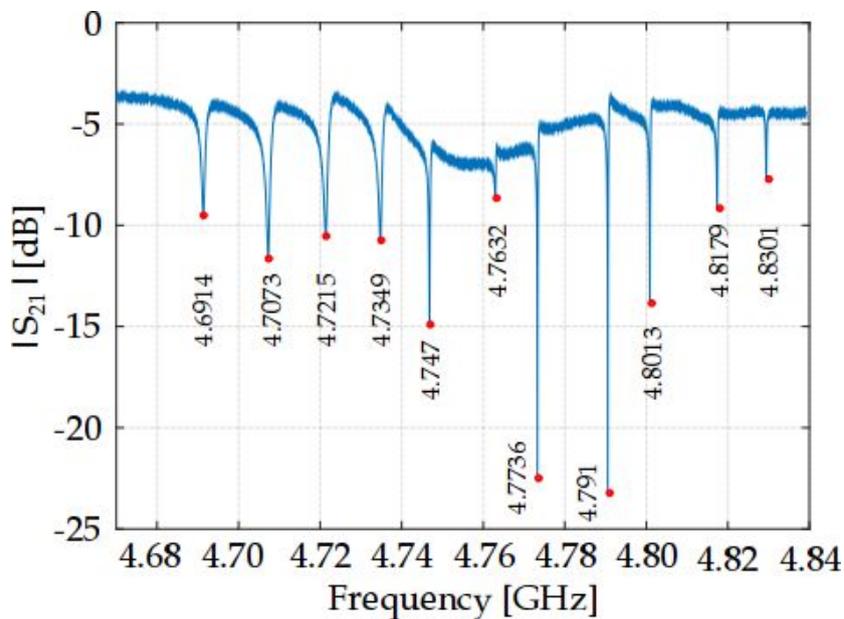
HOLMES array read-out: rf-SQUID μ wave mux



TES current signal frequency up-conversion + rf-SQUID response linearization



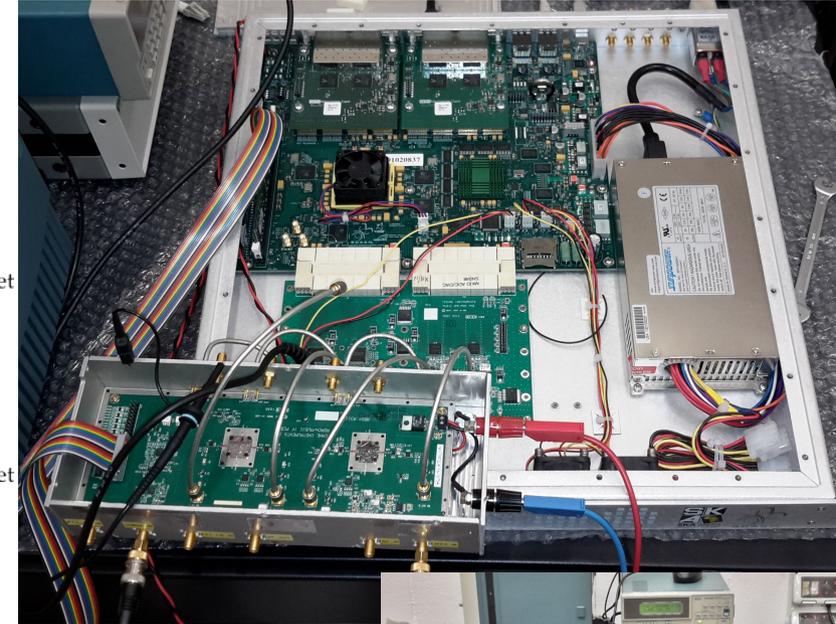
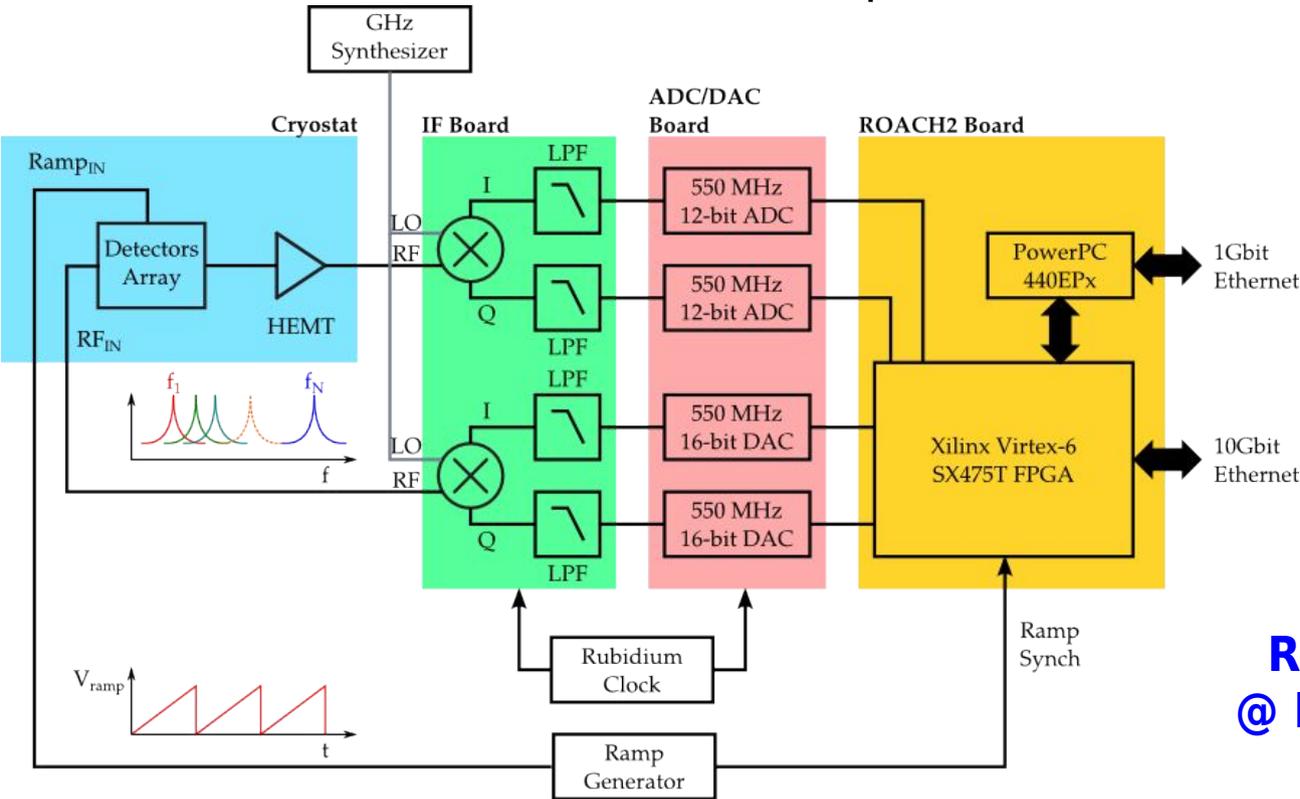
HOLMES TES μ wave read-out testing



TES microwave multiplexing (also for MMC)



Software Defined Radio + flux ramp modulation based on ROACH2



ROACH2 system @ Milano-Bicocca for HOLMES



- **multiplexing factor**

- f_{TES} required bandwidth per channel

$$f_{TES} = n_{\phi_0} f_{sampler} \quad (f_{sampler} = f_{ramp} \text{ from pile-up simulations})$$

$$N_{mux} = \frac{f_{ADC}}{10 f_{TES}} \quad f_{sampler} = 0.5 \text{ MHz}, n_{\phi_0} = 2 \rightarrow N_{mux} \approx 50$$

Conclusions

Today LTD sensor technologies
are powerful and flexible tools for
cutting edge science applications
(largely incomplete review...)

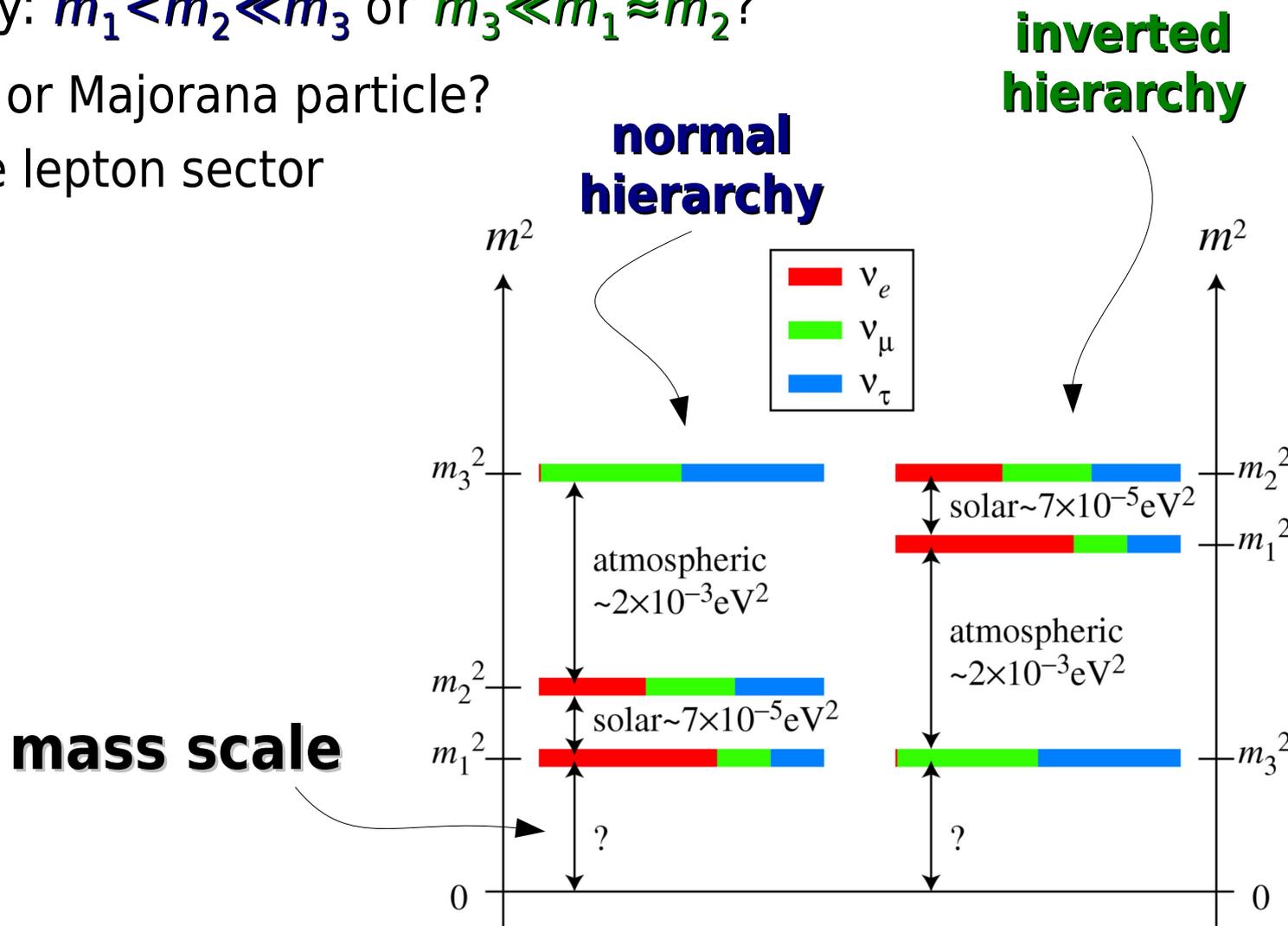
There is room for further improvements
to make devices more *friendly* and flexible

Many more applications are possible

Backup slides ...

Neutrino open questions

- mass scale: i.e. mass of the lightest ν
- degenerate ($m_1 \approx m_2 \approx m_3$) or hierarchical masses
 - ▶ mass hierarchy: $m_1 < m_2 \ll m_3$ or $m_3 \ll m_1 \approx m_2$?
- $\nu = \bar{\nu}$? i.e. Dirac or Majorana particle?
- CP violation in the lepton sector



CUORE sensitivity potential

- sensitivity depends strongly on **background** level
- **NME uncertainties** broaden the sensitivity expectations *

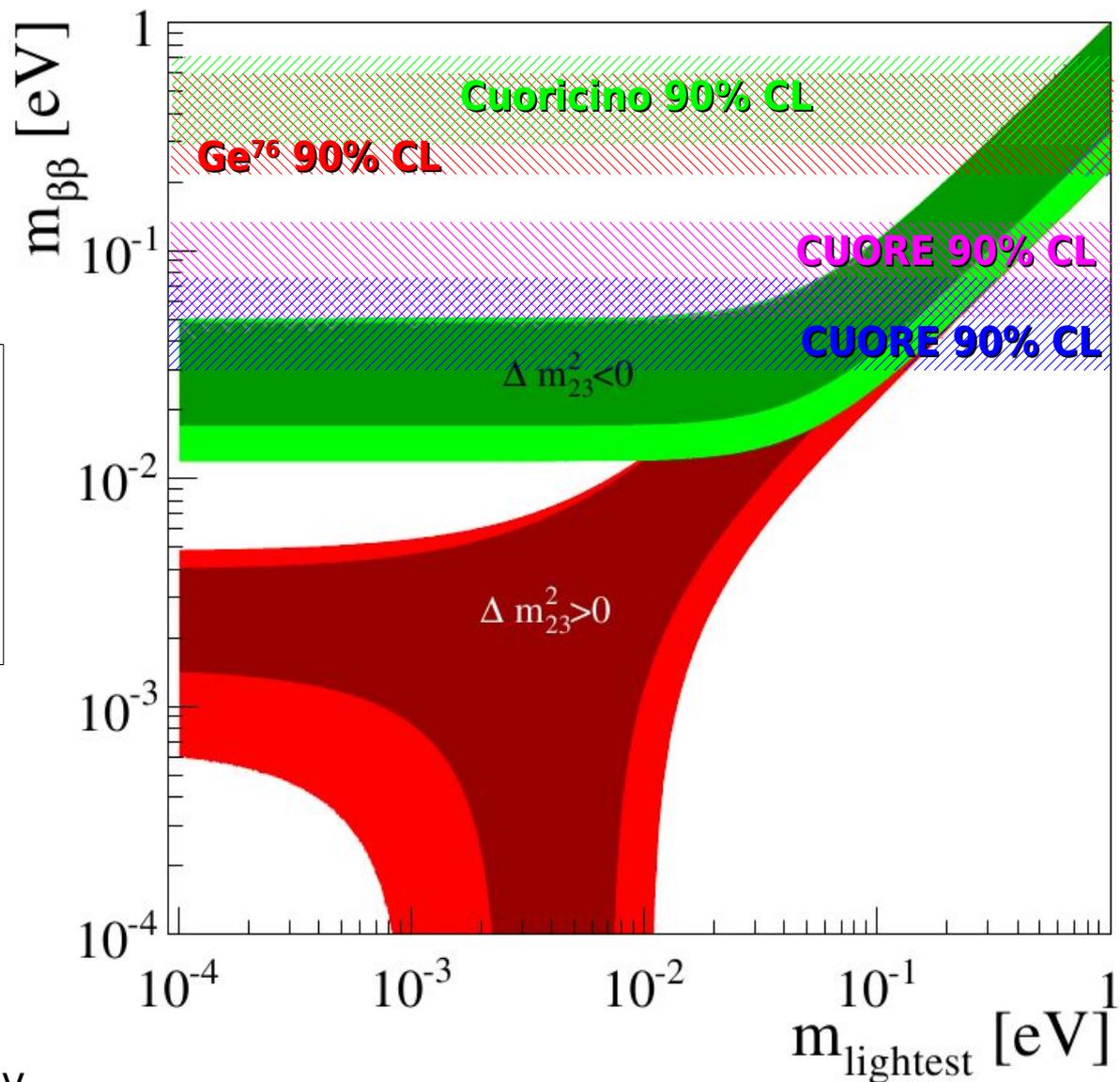
90% CL sensitivity in 5 years

bkg	ΔE	$\tau_{1/2}^{0\nu}$	$\langle m_\nu \rangle$
[c/keV/kg/y]	[keV]	[y]	[meV]
0.01	5	1.0×10^{26}	$50 \div 129$
0.001	5	2.8×10^{26}	$30 \div 77$

Nuclear Matrix Elements from:

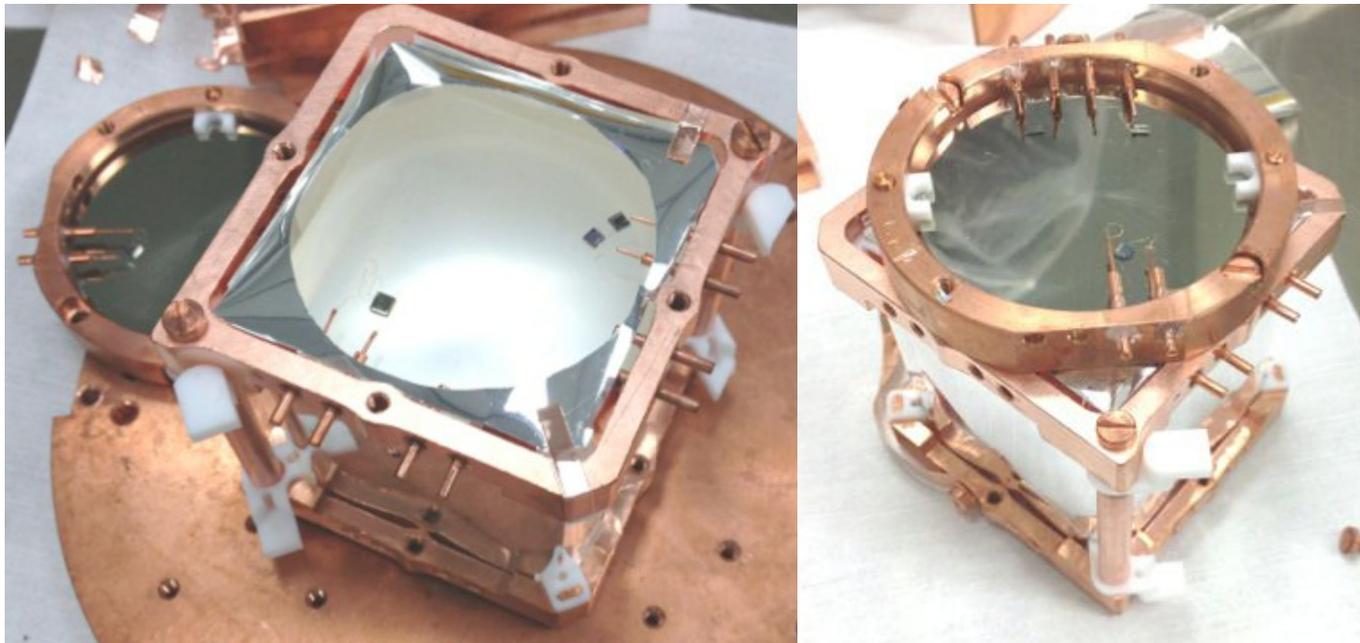
Poves et al., NPA 818 (2009) 139; Faessler et al., JoP G: Nucl. Part. Phys. 39 (2012) 124006; Fang et al., PRC 83 (2011) 034320; Suhonen et al., JoP G: Nucl. Part. Phys. 39 (2012) 124005; Iachello et al., PRC 87 (2013) 014315; P.K. Rath et al., Phys. Rev. C82 064310 (2010); T. R. Rodriguez et al., Phys. Rev. Lett 105 252503 (2010)

- * **NME uncertainties** can be reduced by observing $\beta\beta-0\nu$ in many different isotopes



TeO₂ with Cherenkov

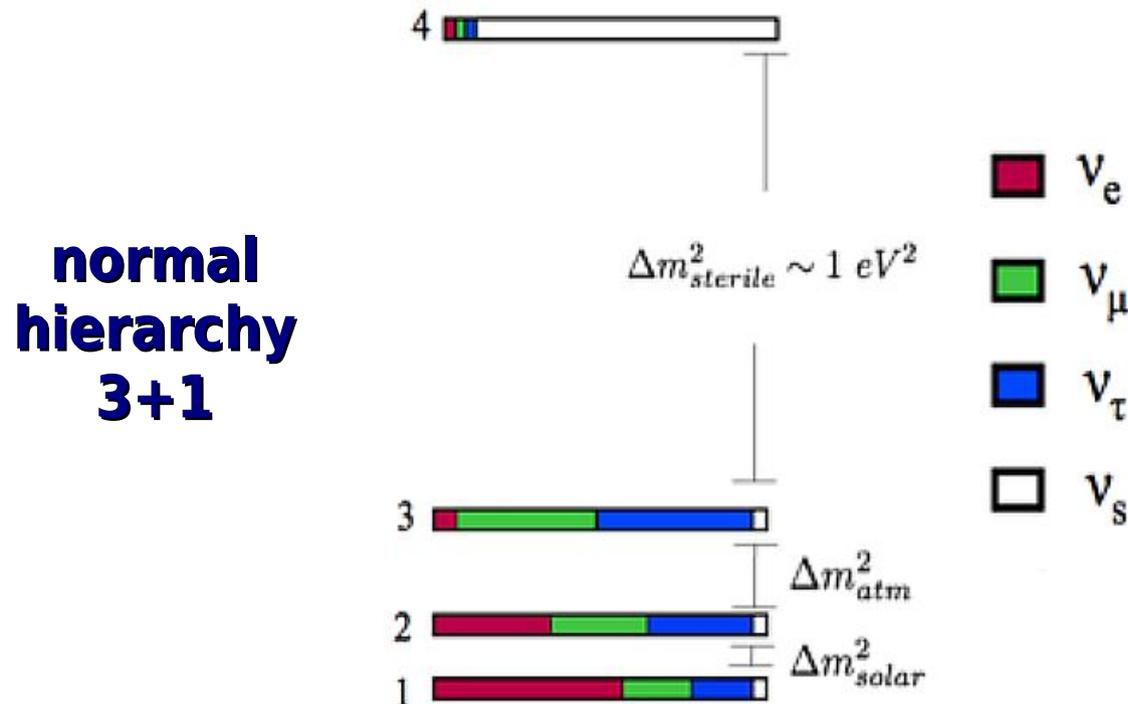
- TeO₂ does not scintillate, but MeV β s emit **Cherenkov radiation**
- threshold for Cherenkov light emission:
 - ▷ 50 keV for β/γ
 - ▷ 400 MeV for α
 - ▷ **α/γ discrimination is possible** (T.Tabarelli de Fatis, EPJC65 (2010) 359)
- Expected light emission: 140 eV/MeV
- Extremely challenging background discrimination



Anomalies and unobserved phenomenologies

- there are **anomalies** in data from
 - ▶ past reactor oscillation experiments (reanalyzed)
 - ▶ short baseline accelerator oscillation experiments (LSND, MiniBOONE)
 - ▶ solar experiment calibration with neutrino sources (GALLEX)
- call for 4th neutrino mass state **$\nu_4 \rightarrow$ sterile neutrino**

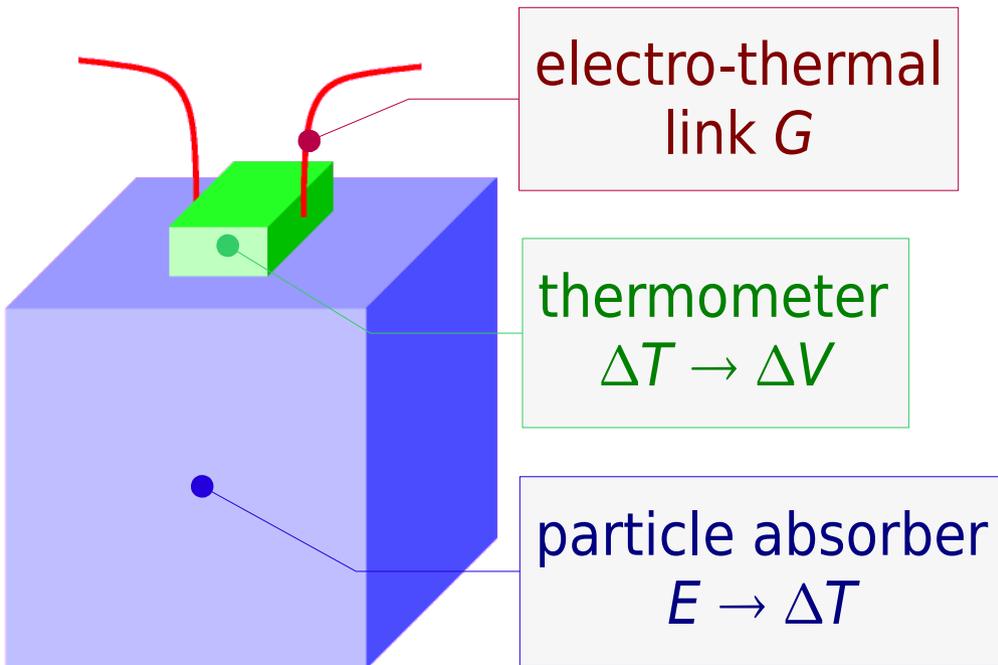
$$\Delta m^2 \approx 1 \text{ eV} \quad \text{and} \quad \sin^2 2\theta \gtrsim 0.1$$



Anomalies and unobserved phenomenologies

- **Sterile (Right Handed) neutrinos** are a natural extension to the Standard Model to include the mass of active neutrinos (ν MSM)
 - ▶ sterile neutrino in the **keV mass range** are perfect candidate as **Warm Dark Matter (WDM)** particles
 - **Standard Model also predicts:**
 - ▶ coherent neutrino scattering on nuclei → never observed!
 - ▶ cosmic neutrino background (CvB) → never observed!
- neutrinos probe physics beyond Standard Model**

Low temperature detectors as calorimeters

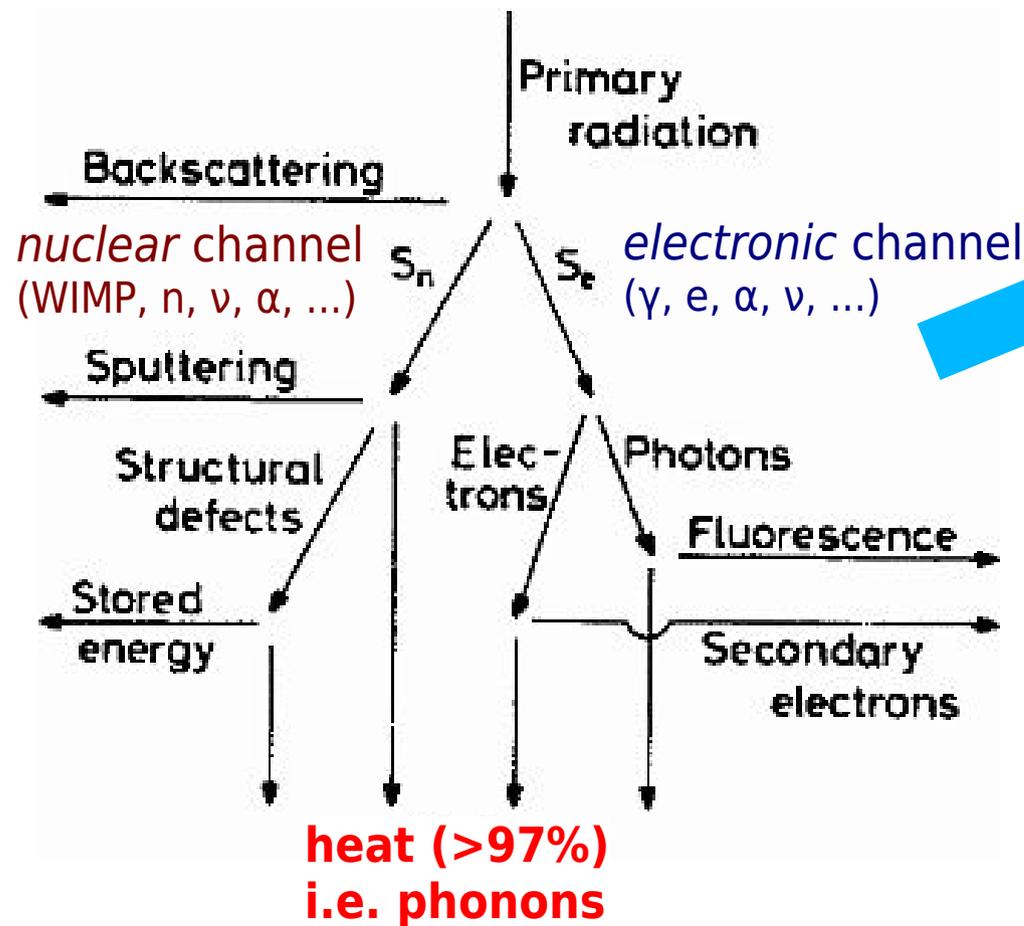


- (quasi-)equilibrium thermal detector
- complete energy *thermalization*
 - ▶ **calorimetry**
- **$\Delta T = E/C$** (C thermal capacity)
 - ▶ low C
 - ▷ low T (i.e. $T \ll 1\text{K}$)
 - ▷ dielectrics, superconductors
- **Pros and cons**
 - ▲ high energy resolution
 - ▲ large choice of absorber materials
 - ▲ true calorimeters
 - ▼ only energy and time informations
 - ▼ slow time response

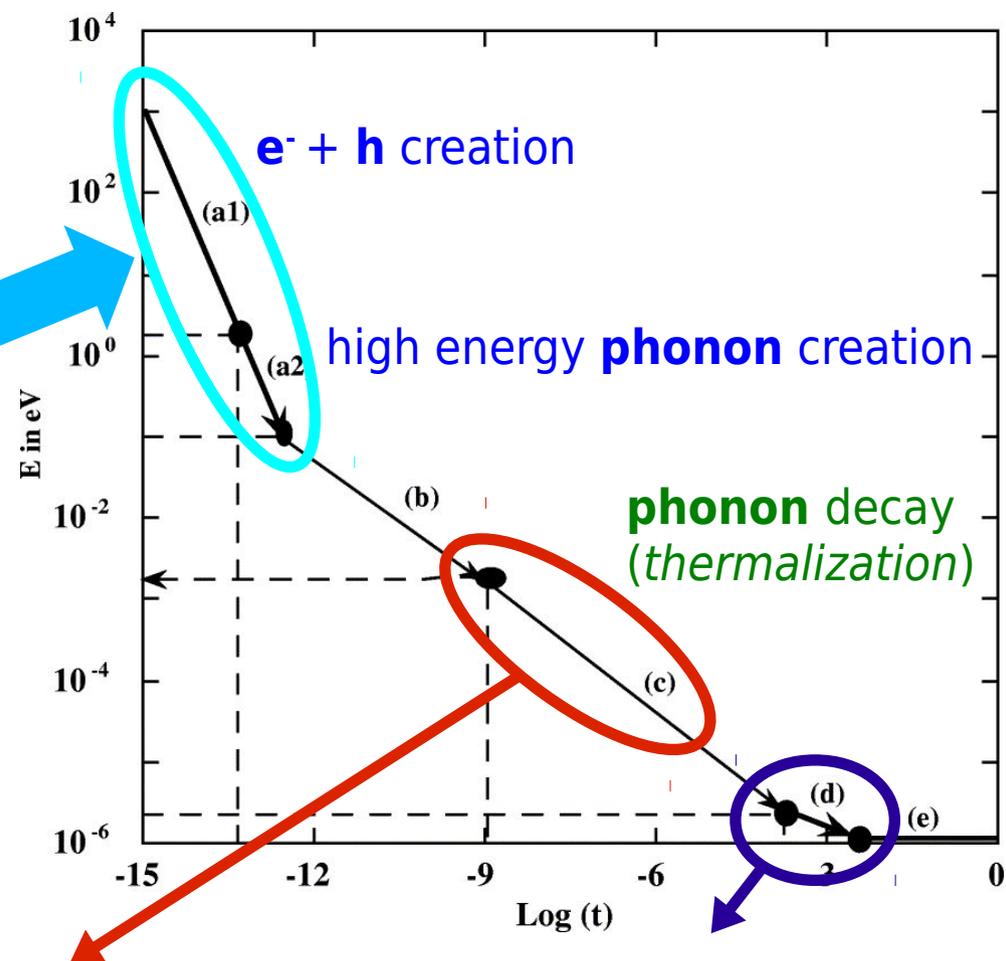
Particle interactions in LTD

directly ionizing: γ , e , α , ...

indirectly ionizing: n , WIMP, ν , ...



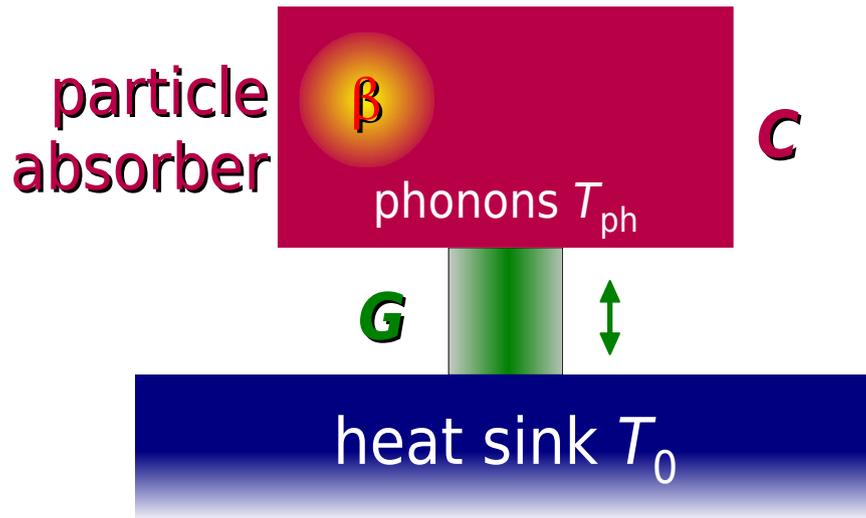
timescale



ballistic phonon propagation
→ **athermal** LTD signal

temperature rise/decay
→ **quasi-equilibrium** LTD signal

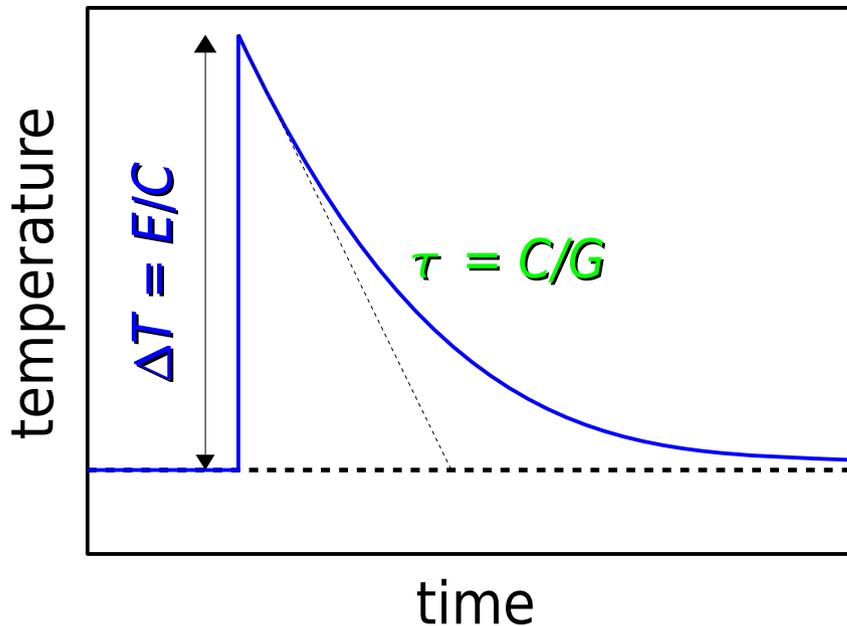
Low temperature detector principles



$$C(T_{ph}) \frac{dT_{ph}}{dt} + G(T_{ph}, T_0) = P(t)$$

$$P(t) = \Delta E \delta(t) \rightarrow T_{ph}(t) = T_0 + \frac{\Delta E}{C} e^{-t/\tau}$$

for $t > 0$ and with $\tau = C/G$



- **750 g of TeO_2 @ 10 mK**

$$C \sim T^3 \text{ (Debye)} \Rightarrow C \sim 2 \times 10^{-9} \text{ J/K}$$

$$1 \text{ MeV } \gamma\text{-ray} \Rightarrow \Delta T \sim 80 \mu\text{K}$$

$$G \sim 4 \times 10^{-9} \text{ W/K} \Rightarrow \tau = C/G \sim 0.5 \text{ s}$$

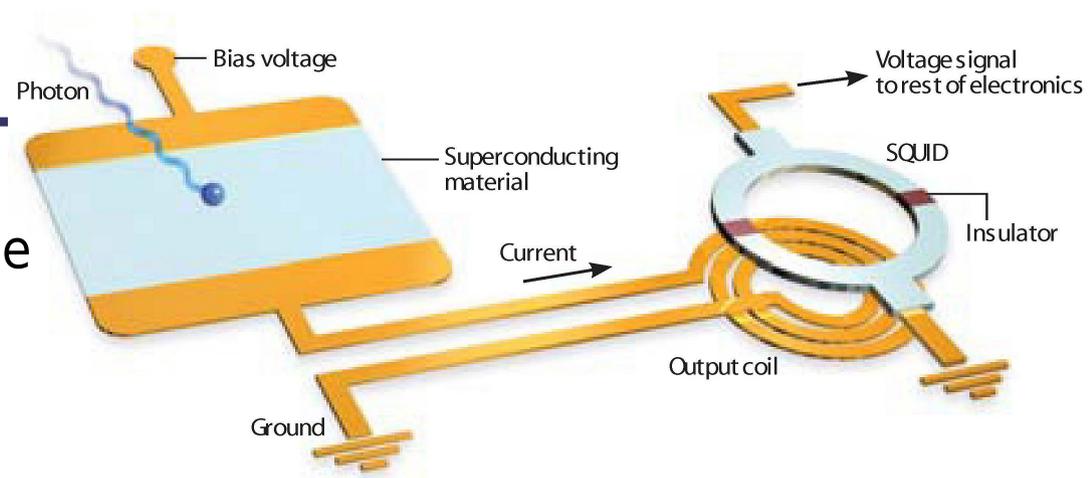
- **1 mg of Re @ 100 mK**

$$C \sim T^3 \text{ (Debye)} \Rightarrow C \sim 10^{-13} \text{ J/K}$$

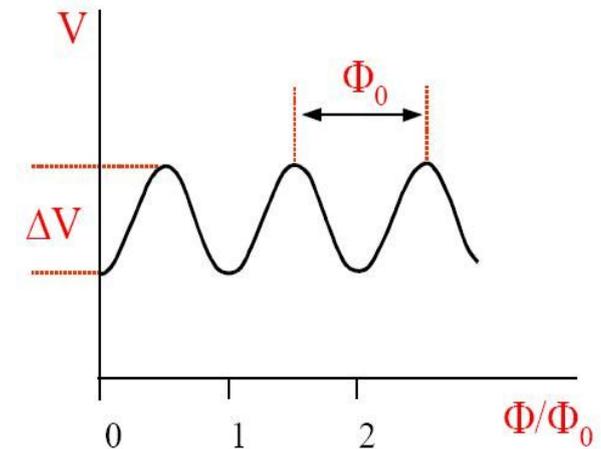
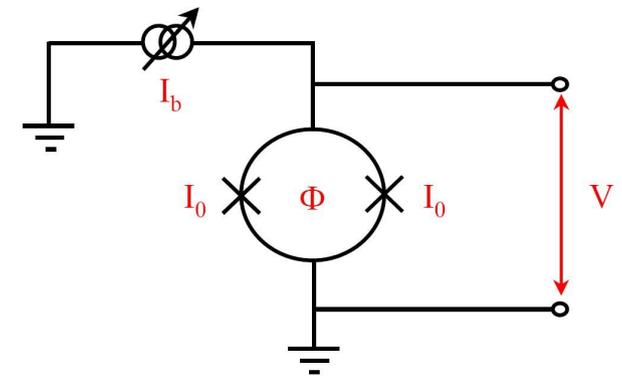
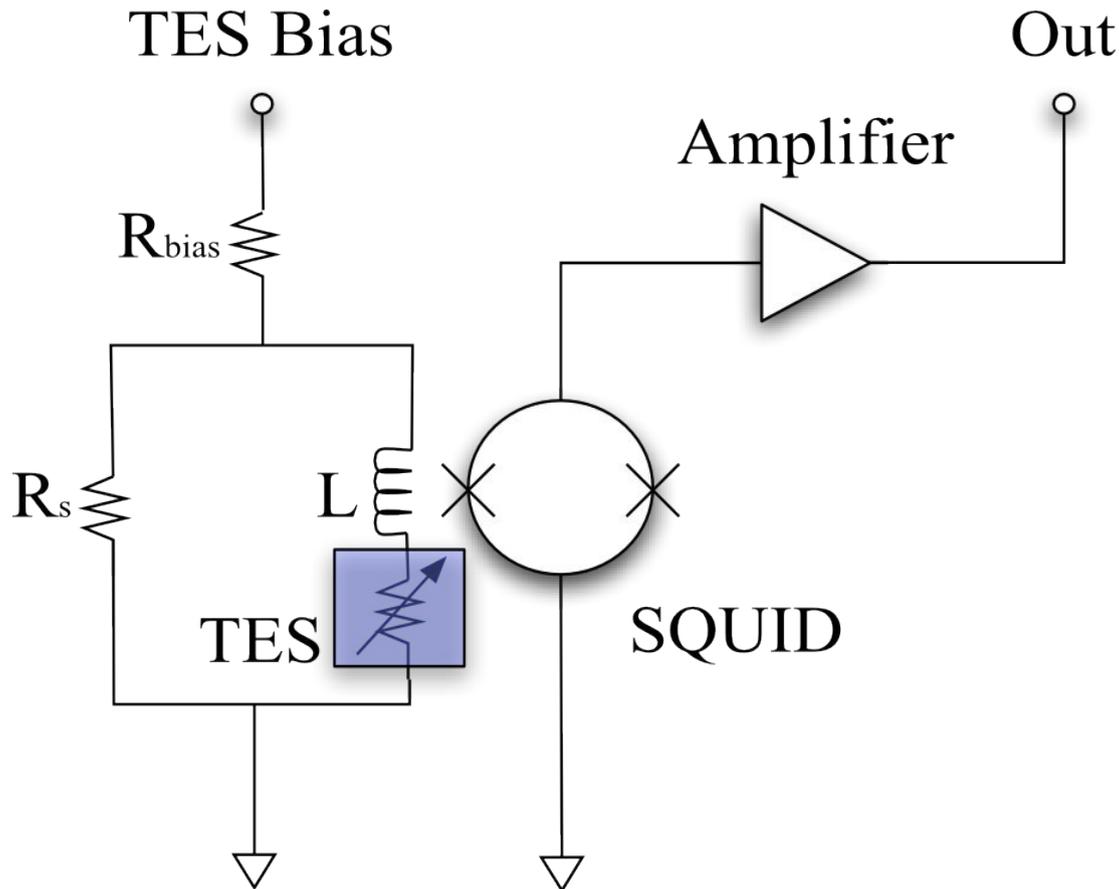
$$6 \text{ keV X-ray} \Rightarrow \Delta T \sim 10 \text{ mK}$$

$$G \sim 10^{-11} \text{ W/K} \Rightarrow \tau = C/G \sim 10 \text{ ms}$$

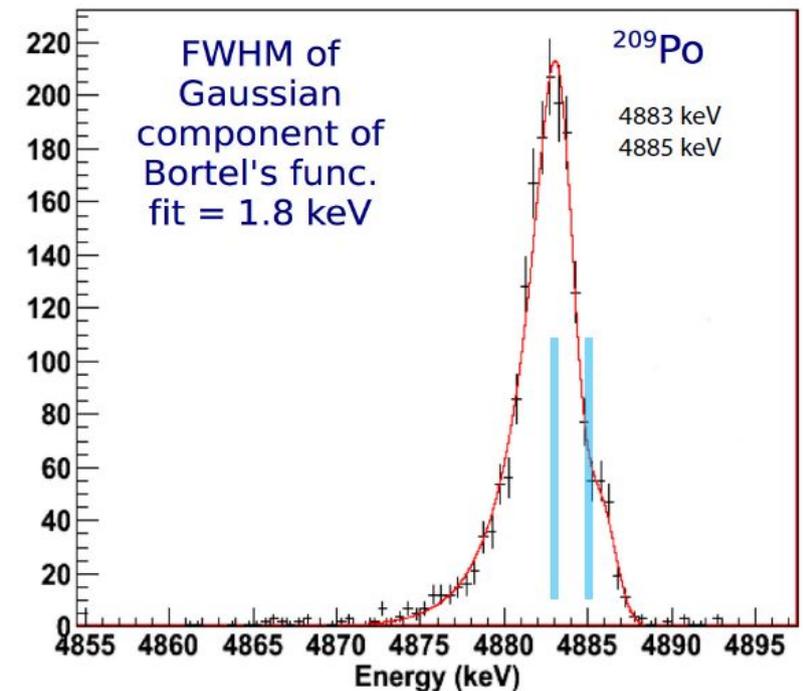
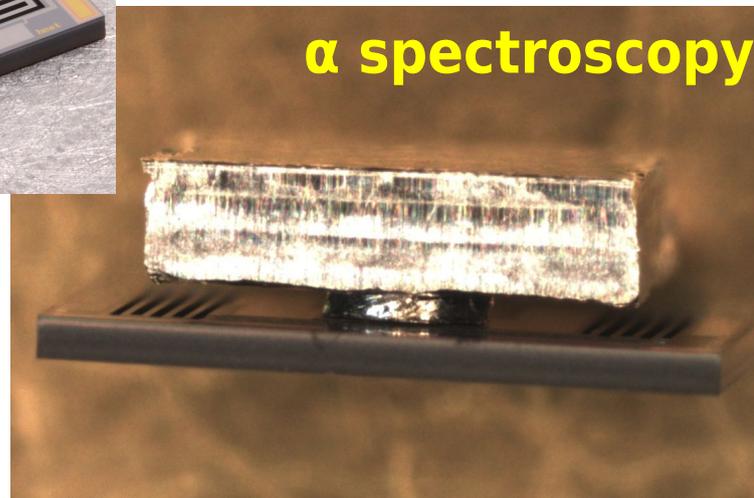
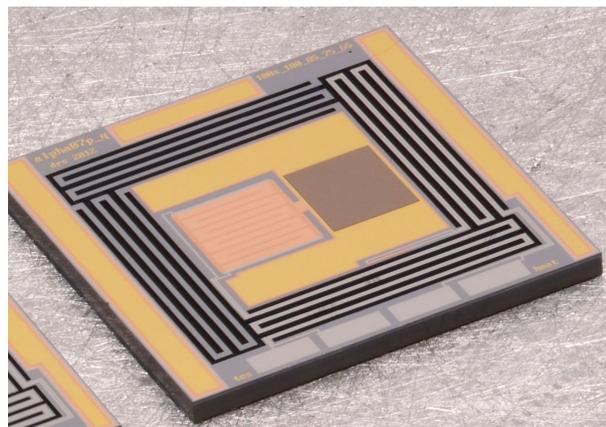
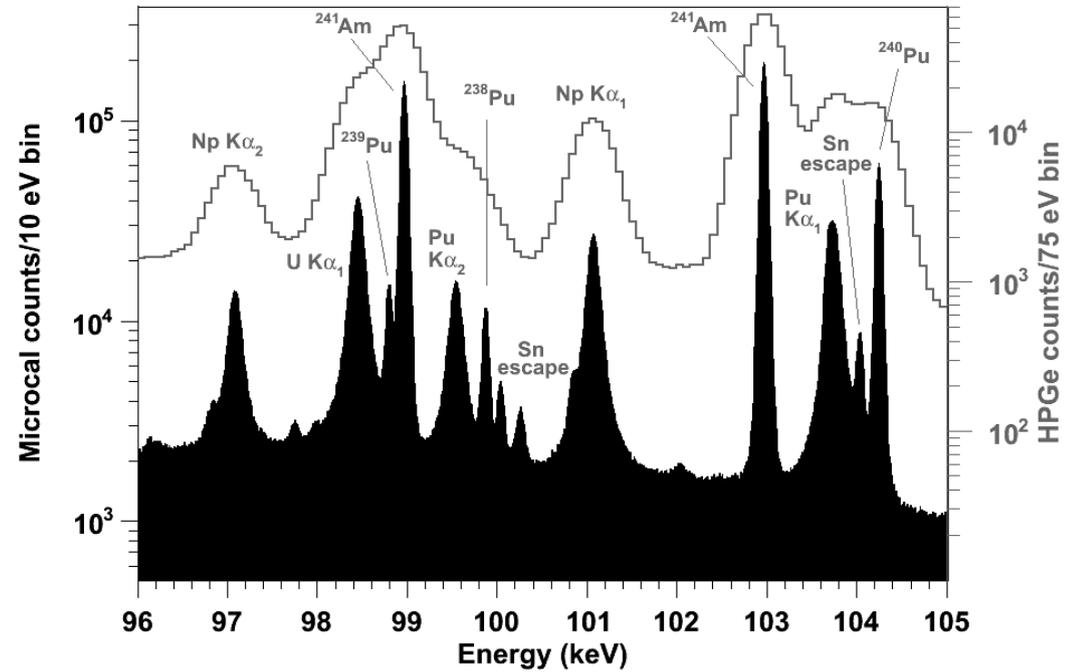
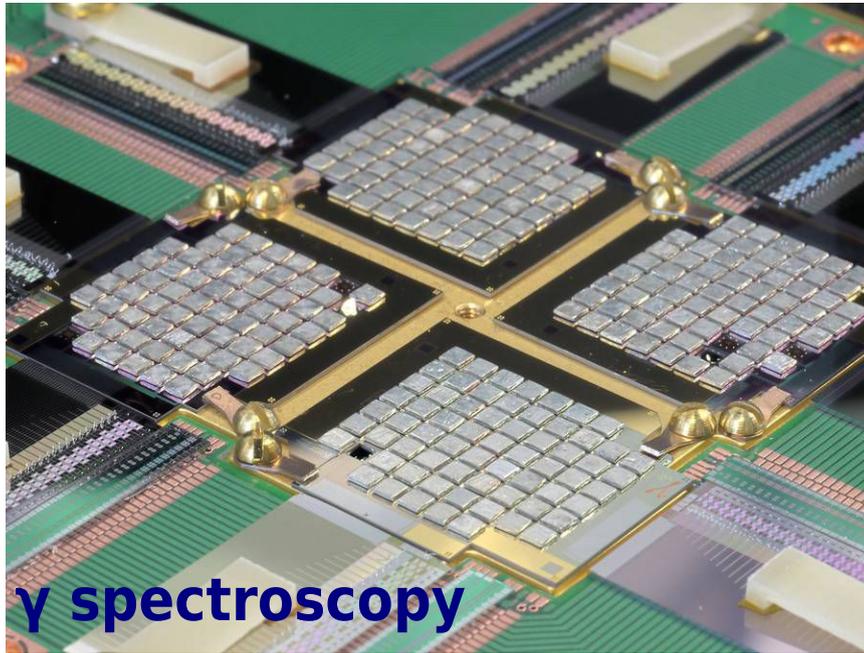
TES read-out: SQUID



- low impedance suitable for multiplexable dc-SQUID magnetometers
- current amplifier configuration
 - ▷ $\Delta I \rightarrow \Delta \Phi \rightarrow \Delta V$
- feedback linearized response

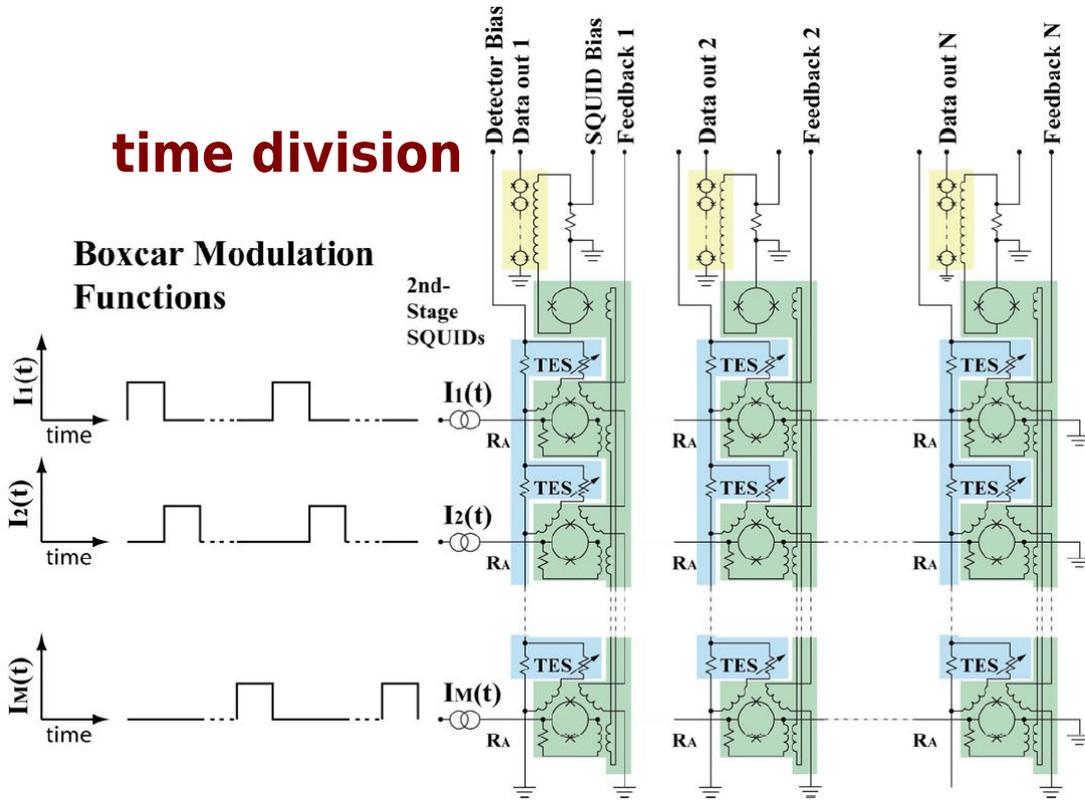


TES arrays for γ and α with Sn absorbers



TES array multiplexing

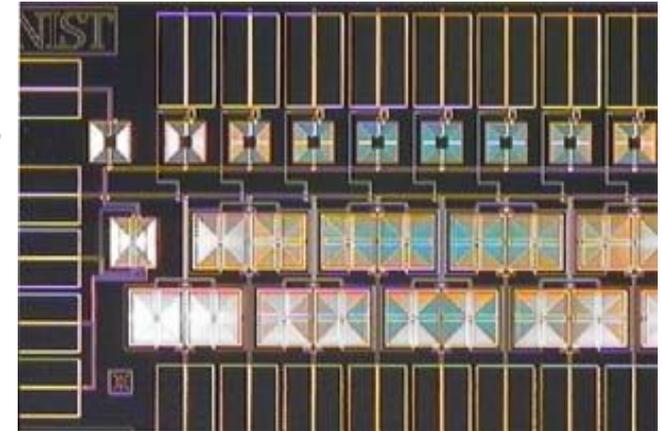
time division



32 channels in 1 column

32 addresses

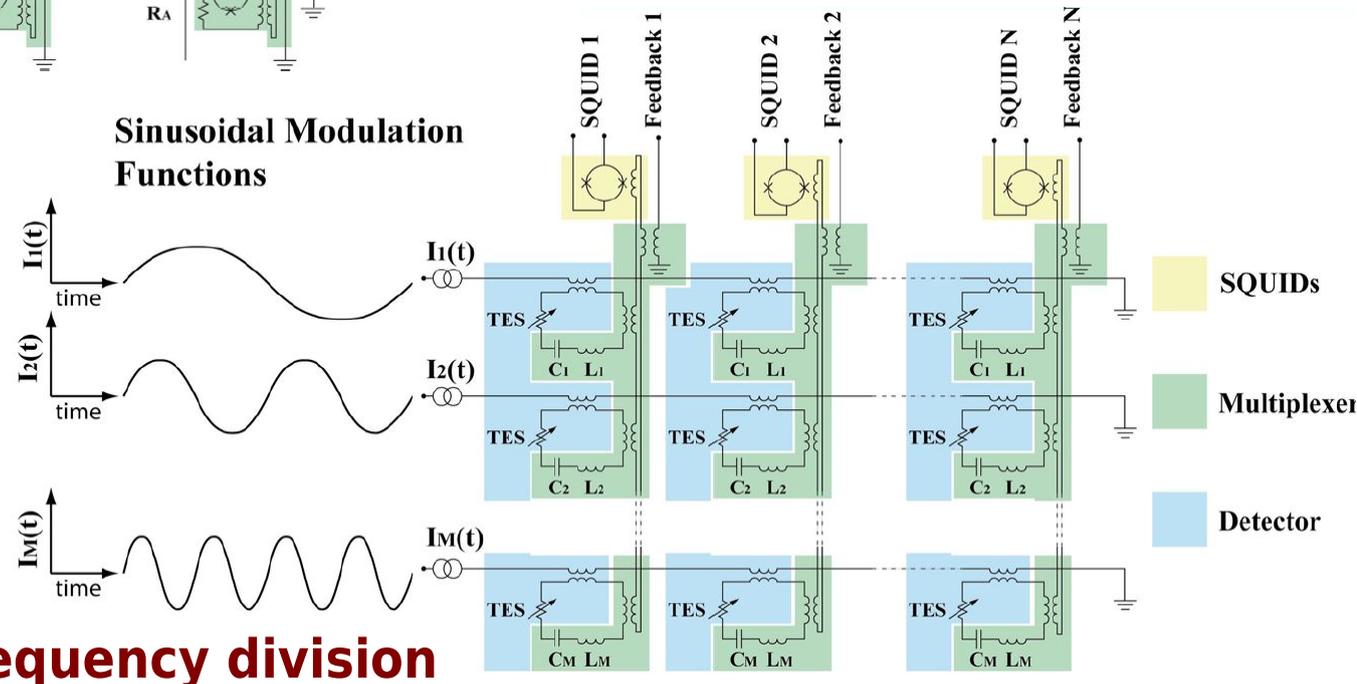
column outputs



32 inputs

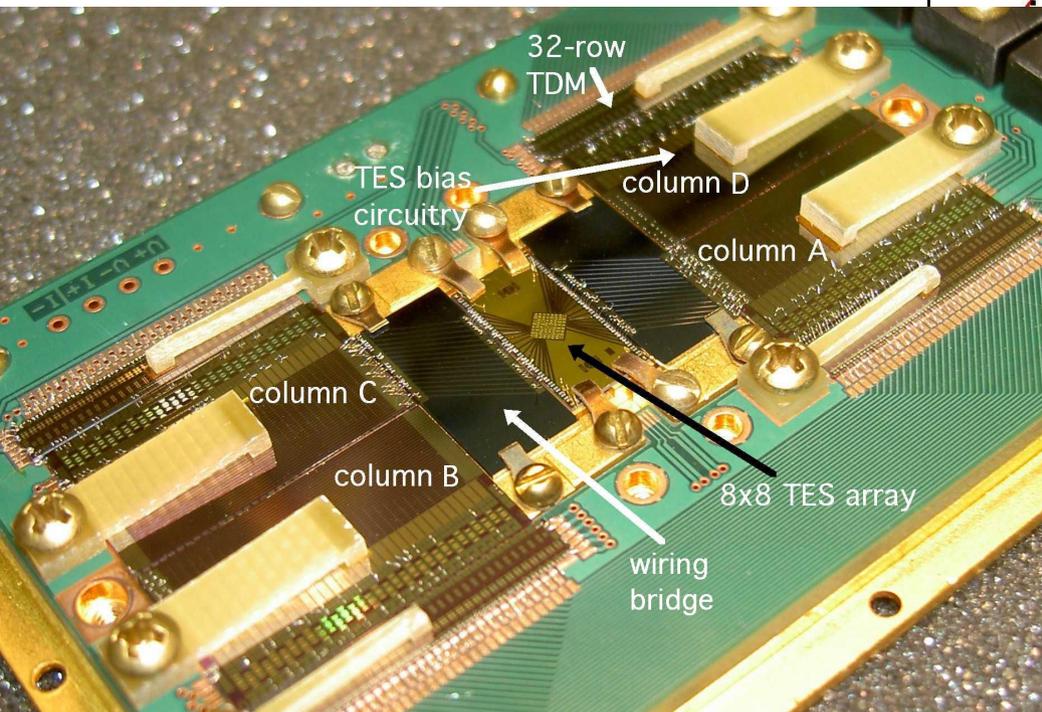
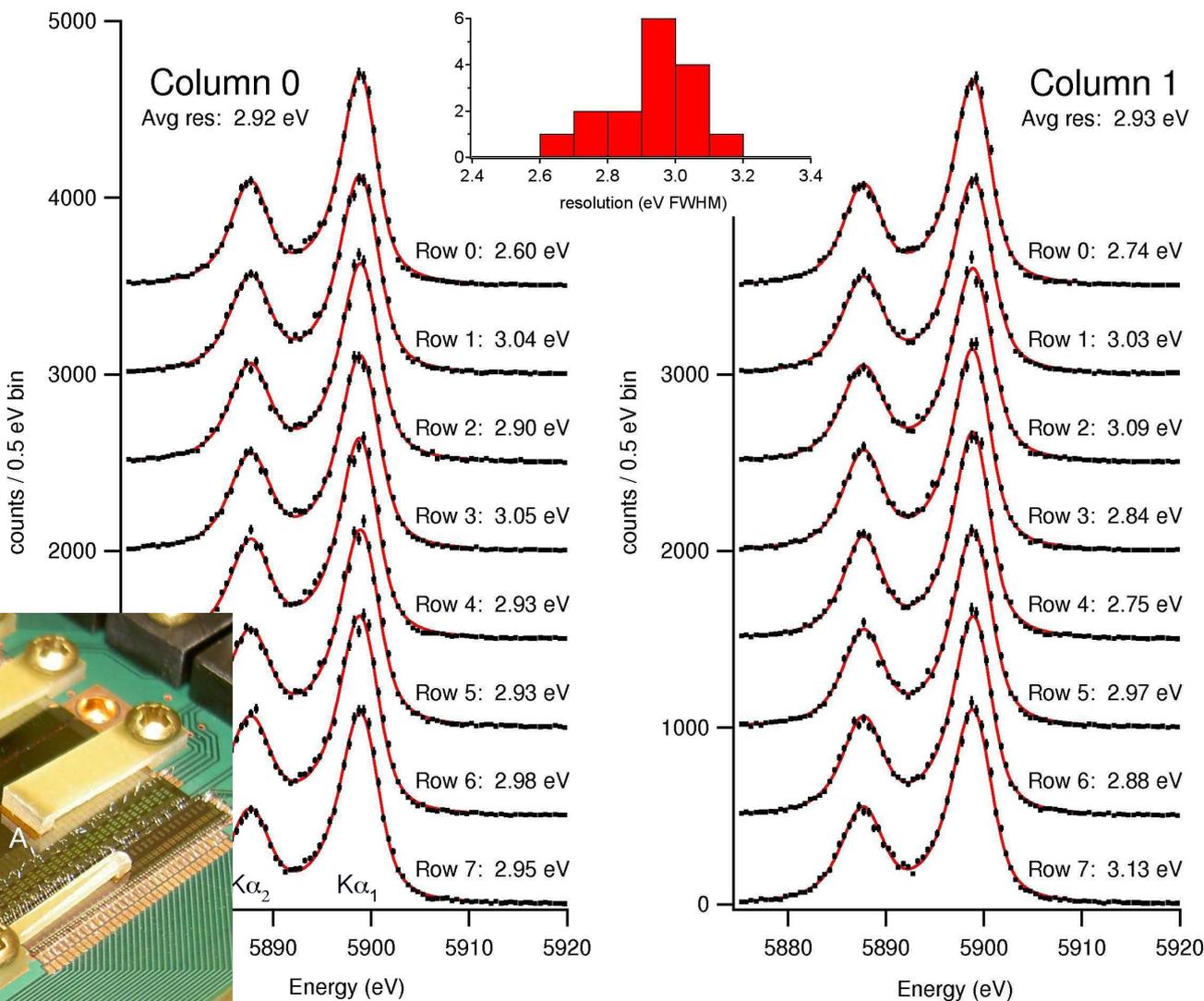
$N \times N$ array
without *mux*
 $N \times N$ readout channels
with *mux*
 N addresses + N outputs

Sinusoidal Modulation Functions

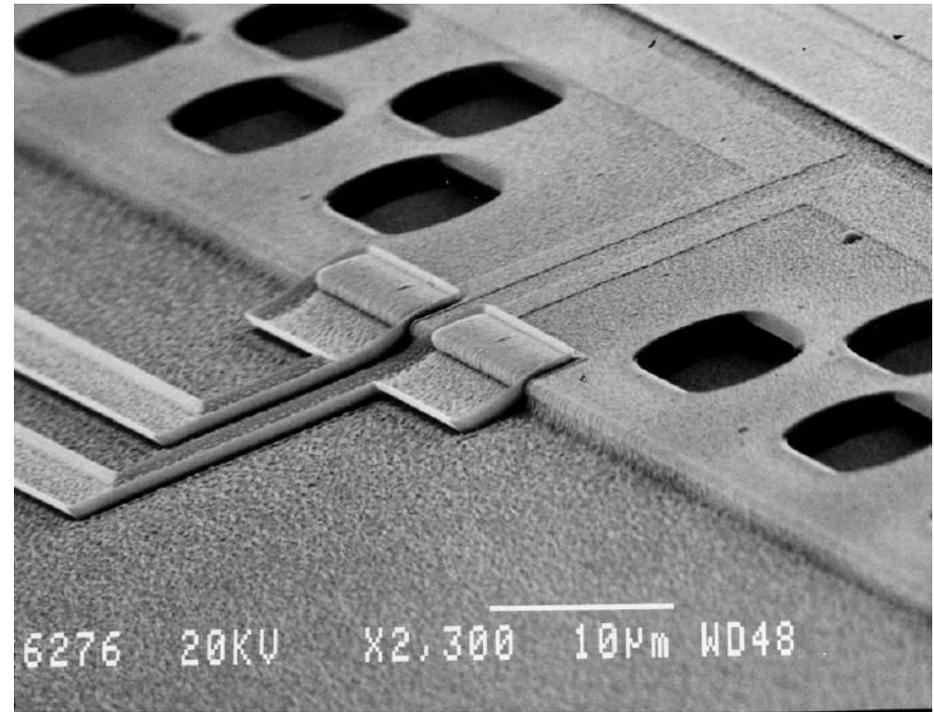
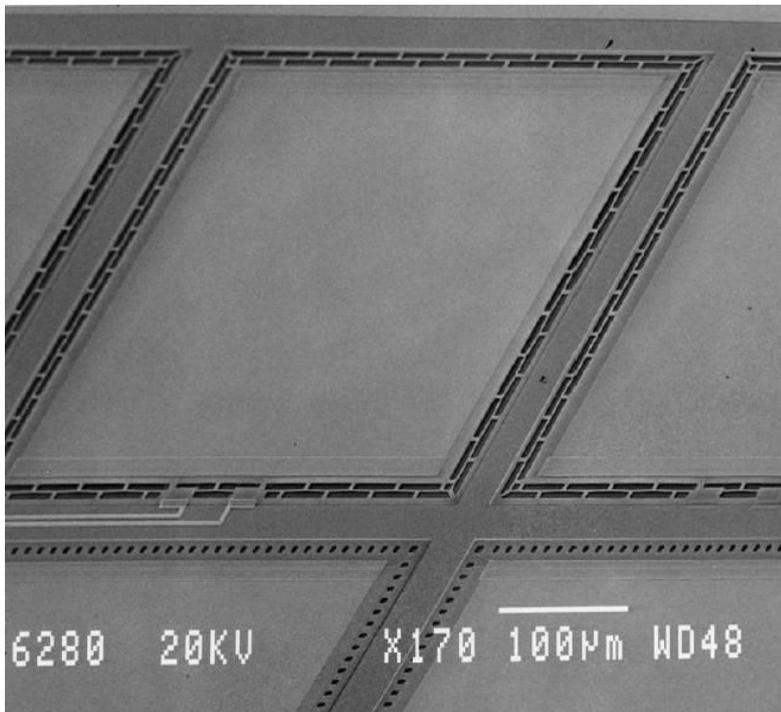
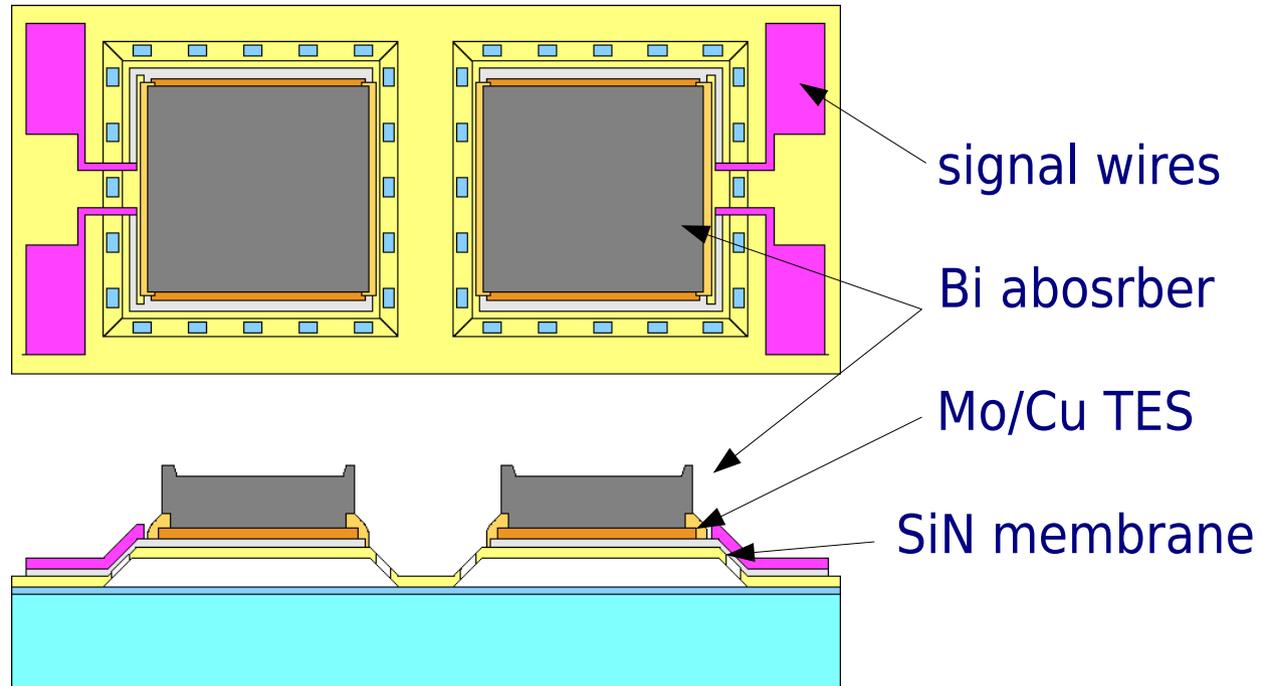


frequency division

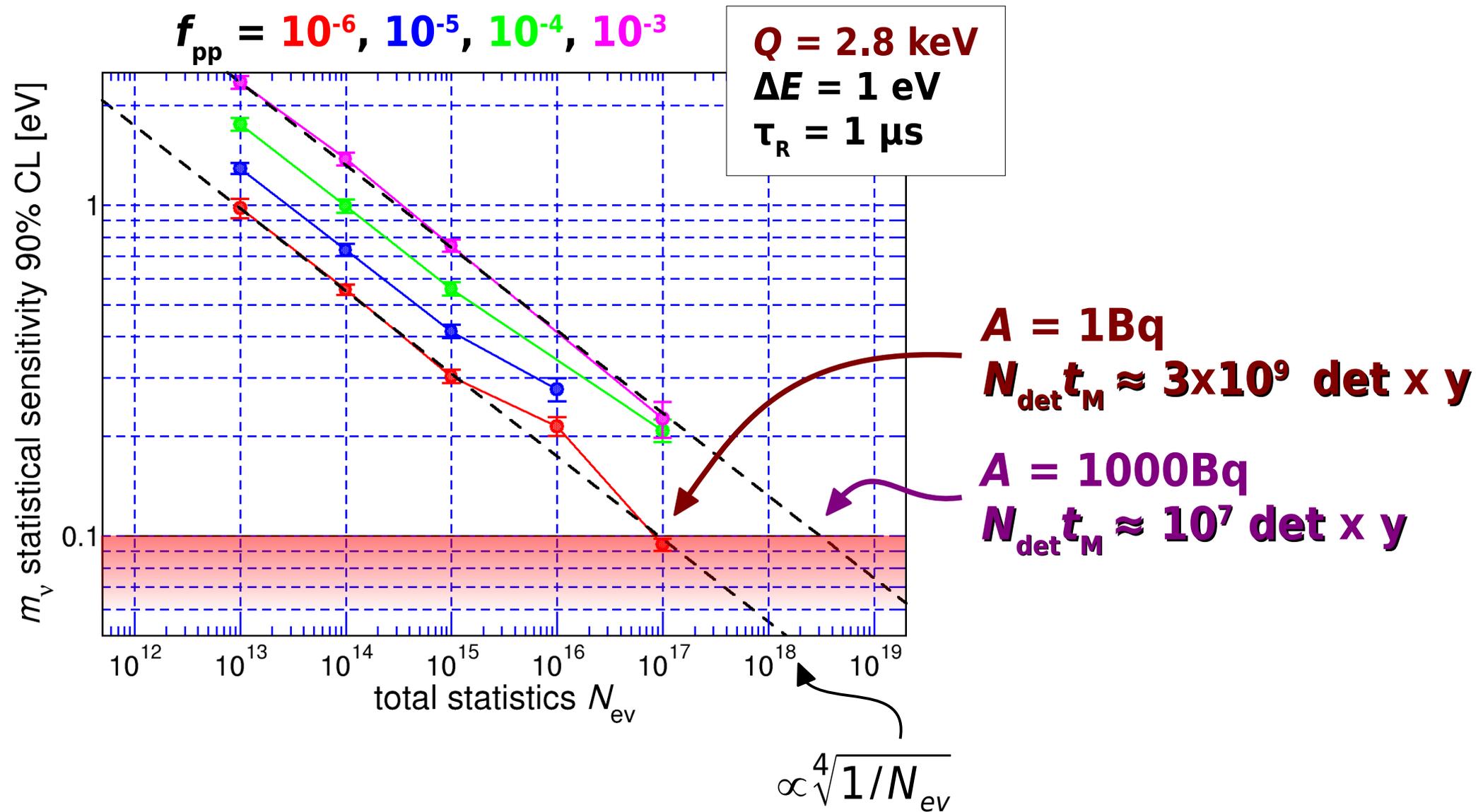
TDM array multiplexing



Micromachined TES arrays / 1



Statistical sensitivity: Montecarlo simulations / 2

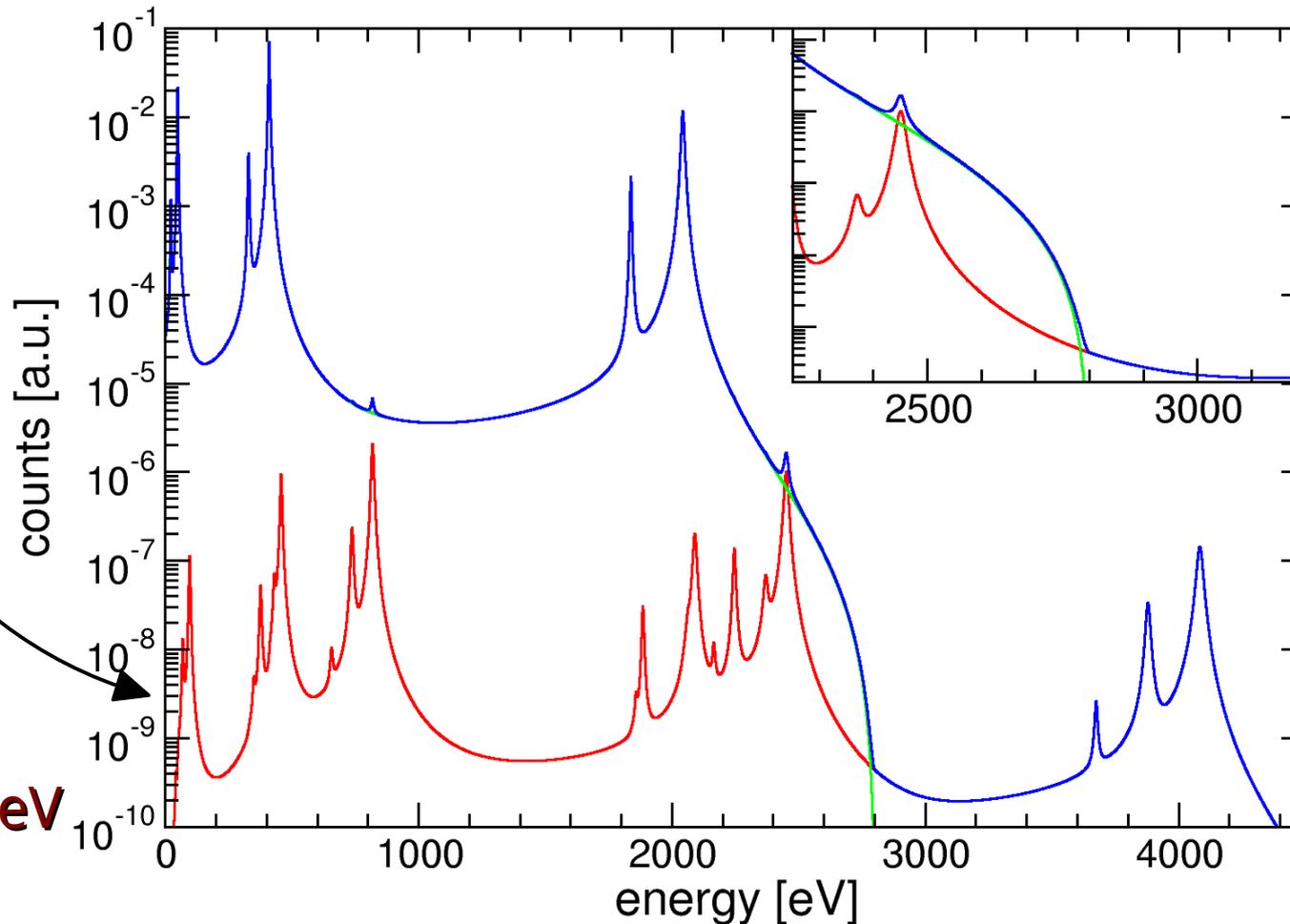


Electron capture end-point experiment / 2

- no direct calorimetric measurement of Q (end-point) so far
- complex pile-up spectrum

▶ $N_{pp}(E) = f_{pp} N_{EC}(E) \otimes N_{EC}(E)$ with $f_{pp} \approx A_{EC} \tau_R$

A_{EC} EC activity per detector
 τ_R time resolution (\approx rise time)



$Q = 2800 \text{ eV}$
 $f_{pp} = 10^{-4}$

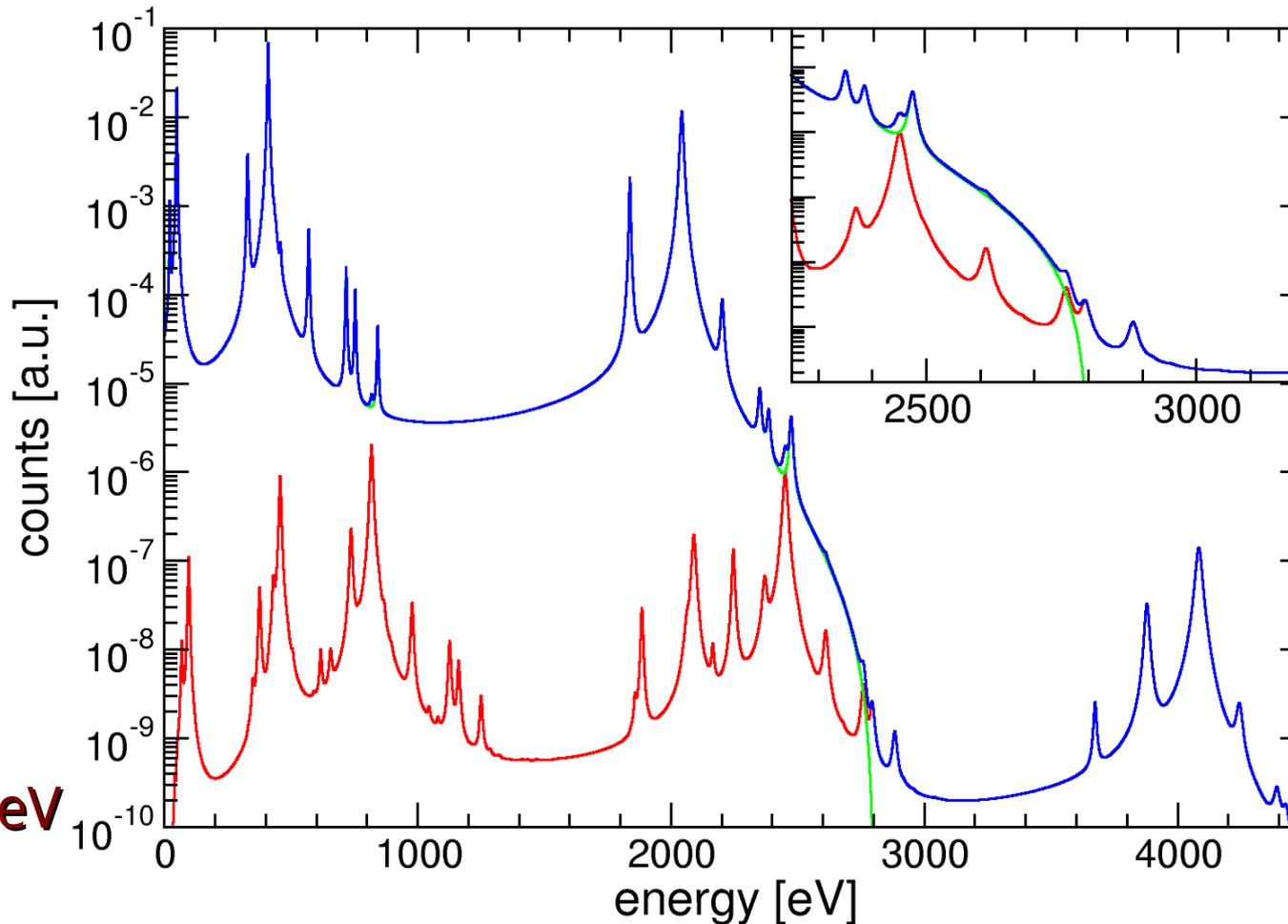
Electron capture end-point experiment / 3

- shake-up/shake-off → double hole excitations
 - ▶ n -hole excitations possible but less probable
 - ▶ authors do not fully agree on energies and probabilities
- even more complex pile-up spectrum
 - ▶ it may be worth keeping f_{pp} smaller than 10^{-4}

A.De Rujula, arXiv:1305.4857

R.G.H.Robertson, arXiv:1411.2906

A.Faessler et al., PRC 91 (2015) 45505



$Q = 2800$ eV

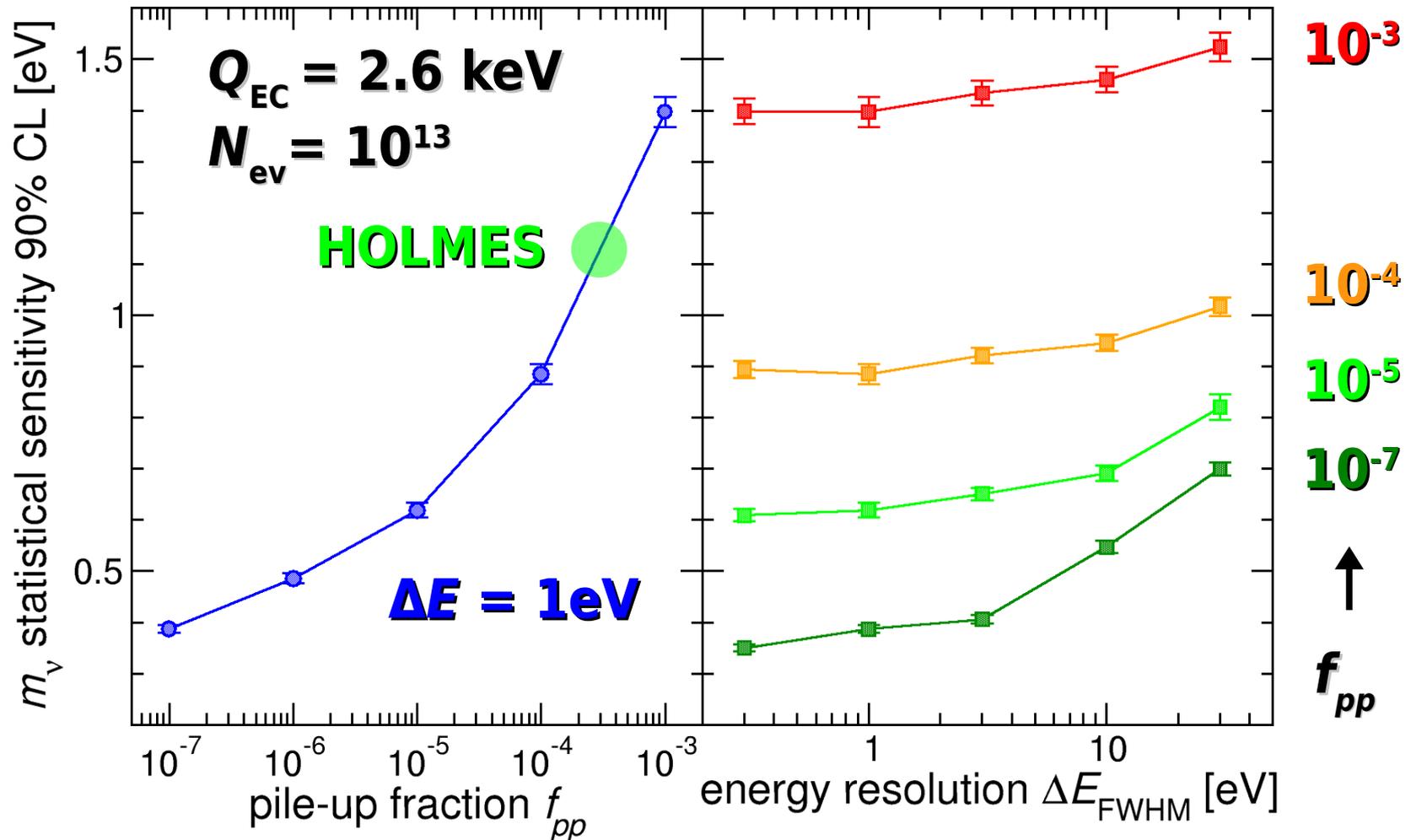
$f_{pp} = 10^{-4}$

HOLMES design: more MC simulations...

Statistical sensitivity $\Sigma(m_\nu)$ dependencies from MC simulations

- **strong** on statistics $N_{\text{ev}} = A_{\text{EC}} N_{\text{det}} t_M$: $\Sigma(m_\nu) \propto N_{\text{ev}}^{1/4}$
- **strong** on rise time pile-up (probability $f_{pp} \approx A_{\text{EC}} \tau_R$)
- **weak** on energy resolution ΔE

t_M measuring time
 N_{det} number of detectors
 A_{EC} EC activity per detector
 τ_R time resolution (\approx rise time)

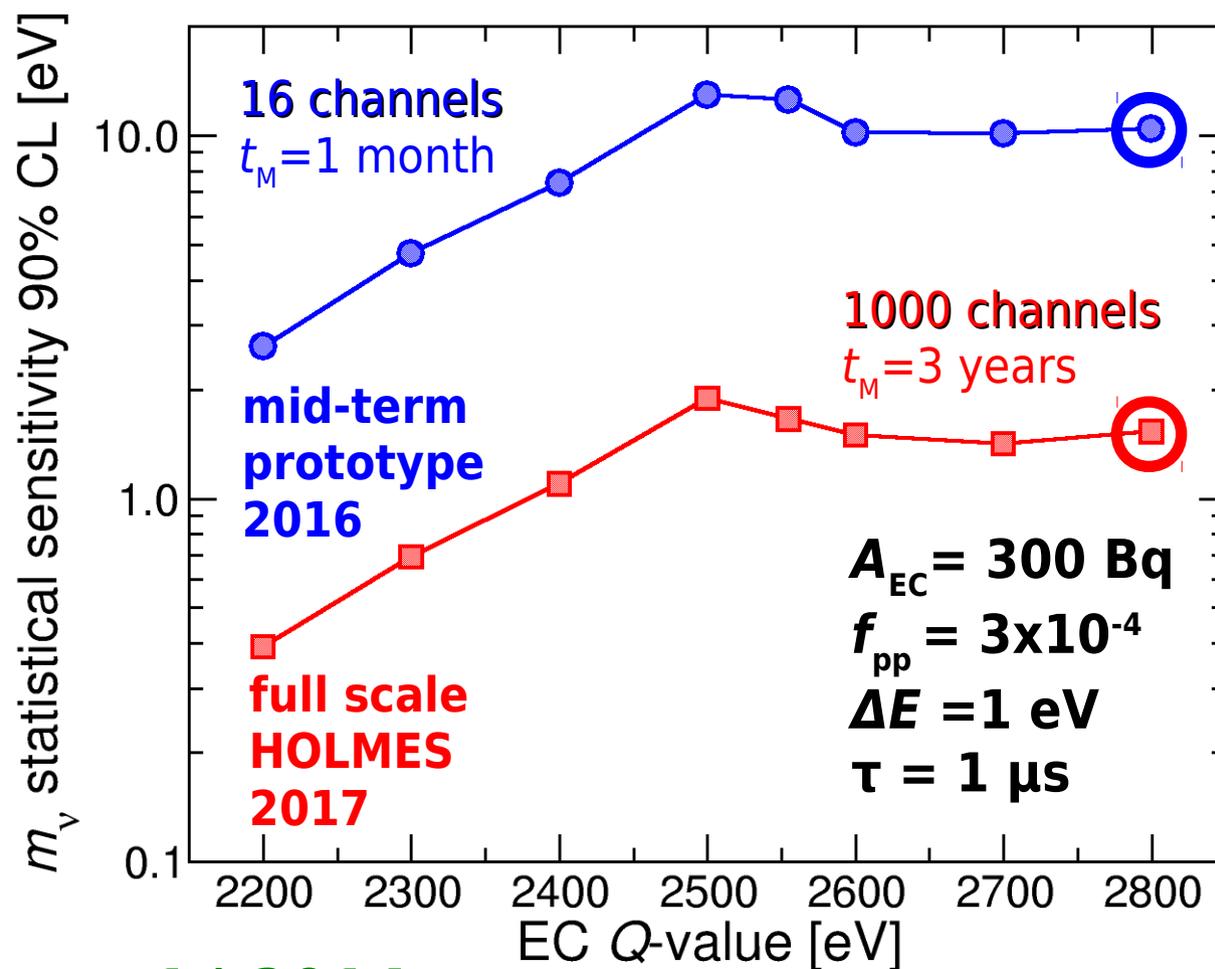


goal

- neutrino mass measurement: m_ν statistical sensitivity as low as 0.4 eV
- prove technique potential and scalability:
 - ▶ assess EC Q-value
 - ▶ assess systematic errors

baseline

- TES with implanted ^{163}Ho
 - ▶ 6.5×10^{13} nuclei per pixel
→ 300 dec/sec
 - ▶ $\Delta E \approx 1\text{eV}$ and $\tau_R \approx 1\mu\text{s}$
- 1000 channel array
 - ▶ 6.5×10^{16} ^{163}Ho nuclei
→ $\approx 18\mu\text{g}$
 - ▶ 3×10^{13} events in 3 years

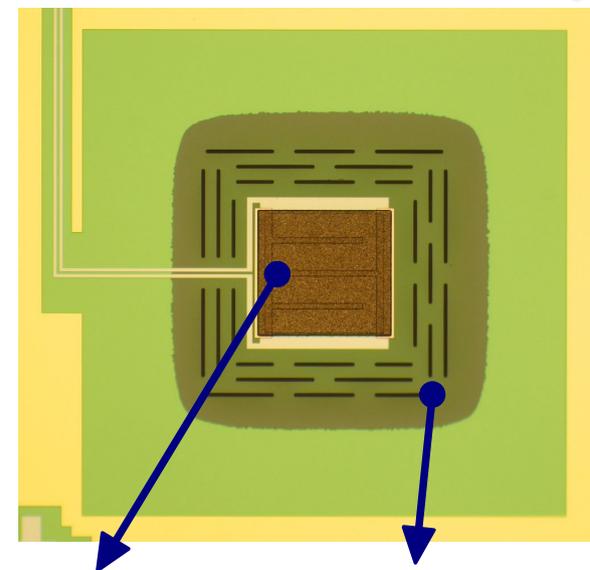


→ Project Started on February 1st 2014

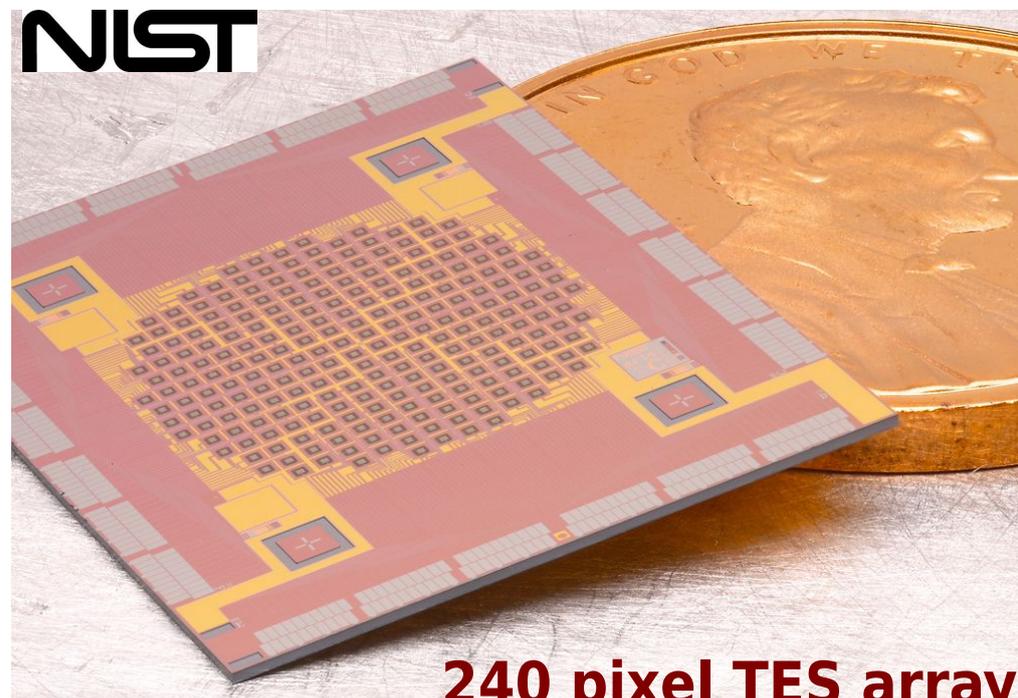
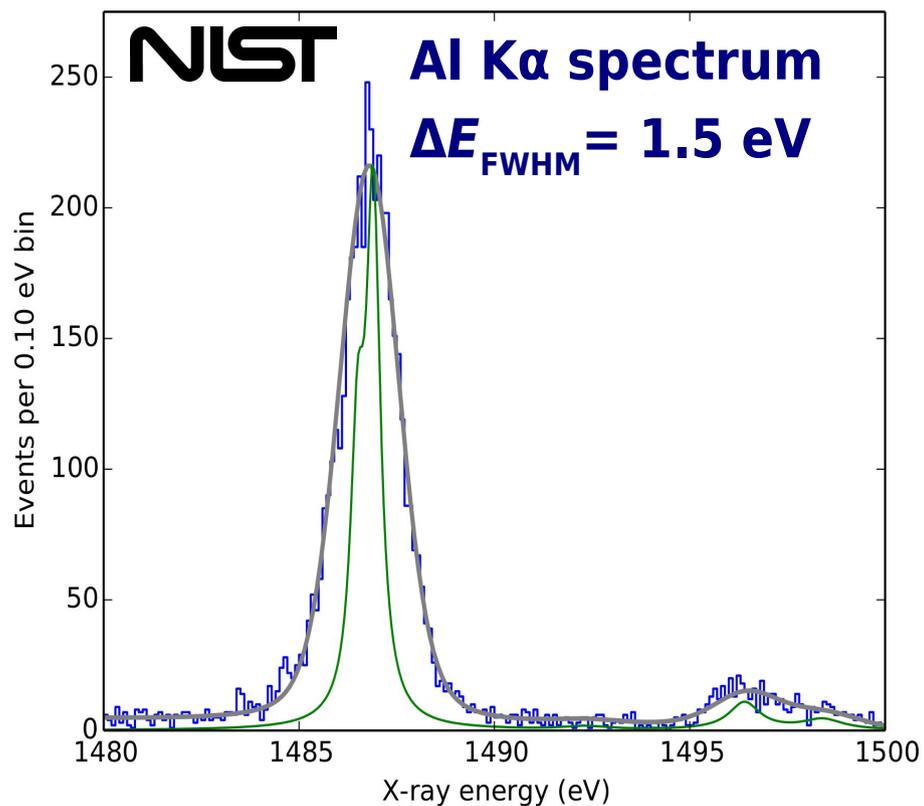
HOLMES detectors



- Transition Edge Sensors (TES) with Au absorber
 - ▷ hot electron microcalorimeters with electro-thermal feedback
 - ▷ electrodeposited Au for full absorption
- MoAu or MoCu proximity TES → $T_c \approx 100\text{mK}$
- on Si_2N_3 membrane



TES with Au absorber membrane Si_2N_3



240 pixel TES array

KIDS multiplexing

